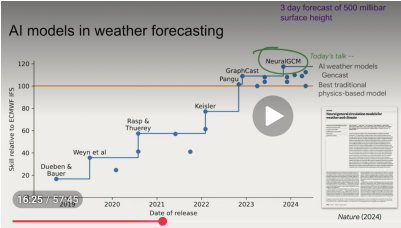
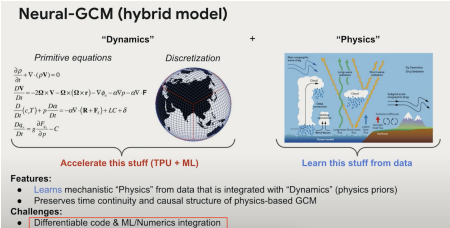

A nebulous presentation on
hybrid models, transport, AI and more

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Hybrid models

New generation of solvers



(Michael Brenner)

Being optimistic, hybrid modeling is also part of a possible future of AI for plasma physics.

Here I show some internal contradictions for this scenario (hybrid modeling for plasma physics).

What is hybrid modeling?

Take your favorite **PDE-based model** and perform **data-enhancement**

$$\begin{cases} \partial_t \mathbf{U} + \mathcal{L}(\mathbf{U}, \partial_x \mathbf{U}, \partial_{xx} \mathbf{U}, \dots) = \mathcal{S}_\theta^{\text{NN}}(\alpha), & \theta \text{ contains all NN weights,} \\ \mathbf{U}(\mathbf{x}, 0) = \beta, & \text{initial data at time } t = 0, \end{cases}$$

where one has a large **dataset of observations**

$$\mathcal{D} = \{(\alpha_i, \beta_i, \gamma_i) \mid i = 1, 2, \dots, N\} \subset \mathbb{R}^{m+n+n}.$$

The cost function evaluated at given final time $T > 0$ is something like

$$J(\theta) = \frac{1}{N} \sum_i \int_{x \in \Omega} |U(x, T : \alpha_i, \beta_i, \theta) - \gamma_i|^2 dx \quad (+ \text{ many other terms}).$$

The cost function is minimized with gradient descent (training)

$$\theta'(t) = -\nabla J(\theta(t)).$$

Take a fully-connected **feed-forward neural network** function

$$f : \mathbb{R}^{a_0} \rightarrow \mathbb{R}^{a_{\ell+1}}$$

$$f = f_{\ell} \circ S_{\ell} \circ f_{\ell-1} \circ S_{\ell-1} \circ \cdots \circ f_1 \circ S_1 \circ f_0$$

where $f_i(x_i) = W_i x_i + b_i \in \mathbb{R}^{a_{i+1}}$ ($x_i \in \mathbb{R}^{a_i}$) is linear

and $S_i \in C^0(\mathbb{R})$ is a non linear activation function $0 \leq S'_i(x) \leq 1$.

- Sigmoidal functions $S_i = \sigma$ are smooth activation functions.
- ReLU function $S_i = R$ with $R(x) = \max(0, x)$ is an extremely popular activation function used in production codes.
- *Max-pooling* used in **convolutional neural networks** is based on the maximum function $(a, b, \dots) \mapsto \max(a, b, \dots) \in \mathbb{R}$.

Training is based on explicit evaluation of the gradient of the cost

$$\theta'(t) = -\nabla J(\theta(t)), \quad J = \text{Loss}.$$

-
- $f \in \text{Lip}(\mathbb{R}^{a_0} : \mathbb{R}^{a_{\ell+1}})$: Rademacher Th. $\implies \nabla f \in L^\infty(\mathbb{R}^{a_0} : \mathcal{M}_{a_{\ell+1}, a_0}(\mathbb{R}))$.
 - ML is universally based on the **chain rule/backpropagation/autodiff**

$$\nabla f(x) = A_\ell(x)B_\ell(x)A_{\ell-1}(x)B_{\ell-1}(x)\dots A_1(x)B_1(x)A_0(x).$$

Non ambiguous in C^1 . **Highly ambiguous** for Lipschitz functions.

Example taken from Karpathy's microgpt.py

```
The most atomic way to train and run inference for GPT in Python. @karpathy
def log(self): return Value(math.log(self.data), (self,), (1/self.data,))
def exp(self): return Value(math.exp(self.data), (self,), (math.exp(self.data),))
def relu(self): return Value(max(0, self.data), (self,), (float(self.data > 0),));
```

Last line means $R'(0) = 0$, where $R(x) = \max(0, x)$.

- Do as if ReLU, defined by $R(x) = \max(0, x)$, is smooth at $x = 0$.
- From $R(R(x)) = R(x)$, the chain rule $R'(x)R'(R(x)) = R'(x)$ yields

$$R'(0)^2 = R'(0) \implies R'(0) = 0 \text{ or } 1.$$

- From $R(x) - R(-x) = x$, we deduce

$$2R'(0) = 1 \implies R'(0) = 1/2.$$

- **Of course, no solution is satisfactory.**

Same phenomenon for max (Boustany paradox)

- For $x = (1, 2, 3, 4)$ encode the function $t \mapsto f(t) = \max_1(tx) - \max_2(tx)$.

```
def max1(x):
    res = x[0]
    for i in range(1, a):
        if x[i] > res:
            res = x[i]
    return res
```

```
def max2(x):
    return torch.max(x)
```

- Pytorch-autodiff** is (would be similar with TensorFlow, Scikilearn, ...),

t	-1	-0.5	-0.01	0	0.01	0.5	1
$f(t)$	0	0	0	0	0	0	0
$f'(t)$	0	0	0	-1.5	0	0	0

- The paradox can be solved in the good algebra.
- Since such pathologies concerns only the **blue part**, let's us decide we will never use ReLU for modeling data in hybrid models for plasma physics.

Red part : Discrete transport schemes

- Take your favorite advection equation $\mathbf{t} \geq 0, \mathbf{x} \in \mathbb{R} : \partial_t \mathbf{u} + \mathbf{a} \partial_x \mathbf{u} = 0$.
- The upwind scheme is

$$\frac{u_j^{n+1} - u_j^n}{\Delta t} + \mathbf{a} \frac{u_{j+\frac{1}{2}} - u_{j-\frac{1}{2}}}{\Delta x} = 0,$$

where $u_{j+\frac{1}{2}} = \mathbf{G}_a(u_j^n, u_{j+1}^n)$ with

$$\mathbf{G}_a(\alpha, \beta) = \alpha \text{ if } \mathbf{a} > 0, \quad \mathbf{G}_a(\alpha, \beta) = \beta \text{ if } \mathbf{a} \leq 0.$$

- The upwind scheme is the basis of many schemes (all stable under CFL).
- The time steps are composed one after the other. That is any numerical scheme can be seen as the composition of basic functions.

Principle

$\{\text{Numerical schemes}\} \subset \{\text{Neural Nets}\}.$

Autodiff and discrete transport schemes

What matters is the regularity of the flux function F

$$au_{j+\frac{1}{2}} = F(a, u_j, u_{j+1}) = aG_a(u_j, u_{j+1}) = R(a)u_j - R(-a)u_{j+1}.$$

- Unfortunately the function $a \mapsto R(a)$ is not differentiable at $a = 0$ (in ML codes, usual convention is $R'(0) = 0$).

Principle

The non-differentiable ReLU function $R(x) = \max(0, x)$ is (almost) everywhere in advection dominated numerical codes.

Even if the code does not blow up, we just don't know what's going on.

- Thiry-Li-Mémin-Roullet, G. A unified formulation of quasi-geostrophic and shallow water equations 2024.

Even worse for high order schemes

Consider the MinMod TVD scheme

$$\frac{u_j^{n+1} - u_j^n}{\Delta t} + a \frac{u_{j+\frac{1}{2}}^n - u_{j-\frac{1}{2}}^n}{\Delta x} = 0,$$

where

$$\begin{cases} a \geq 0 & u_{j+\frac{1}{2}} = u_j + \frac{1}{2}(1 - \nu) \min\text{mod}(u_{j+1} - u_j, u_j - u_{j-1}), \\ a < 0 & u_{j+\frac{1}{2}} = u_{j+1} + \frac{1}{2}(1 - \nu) \min\text{mod}(u_{j+1} - u_j, u_j - u_{j-1}). \end{cases}$$

Check that $\min\text{mod}(x, y) = R(-\max(-x, -y)) - R(-\max(x, y))$.

Lemma

The 2nd order Lax Wendroff flux belongs to the regularity class C_{pw}^1 .

Presumably, all high order conservative fluxes
(TVD, TVB, BV, WENO, MUSCL, ...)
belong to the same class.

Conclusion : The mathematics of autodiff for hybrid transport models generates new interesting questions. (we will soon investigate with Petar Samardzic (new PhD) and Louis Thiry (co-supervisor)).

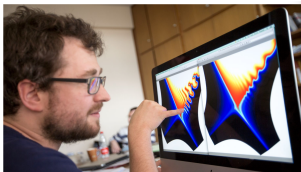
To develop hybrid models in plasmas physics, what is needed is

- Recode to insure compatibility with automatic differentiation (use JAX, ...).
In particular for the transport part.
- At the same time, ensure numerical integration is stable.

This was what Michael Brenner focused on in his presentation on meteorology modeling.

Presumably this holds true for plasma physics as well.

- This part is entirely based on the work of Emmanuel Franck, Victor Michel-Dansac and the INRIA/MACARON team



Numerical schemes for hyperbolic equations
enhanced by Scientific Machine Learning
Victor Michel-Dansac (Inria)

- Scimba is an entirely new library based on PYTORCH for solving PDES, for education and research

<https://www.scimba.org>

They do Claude coding, so development runs at super fast pace.

- It uses some ideas coming from PINNS, but enhanced with a modeler, JAX, natural gradients and other optimization techniques.

- Take your favorite model=PDE + **Boundary Conditions** (or IC), written in symbolic form as

$$\begin{cases} \mathcal{L}u(x) = f(x) & x \in \Omega, \\ u(x) = g(x) & x \in \Gamma = \partial\Omega \end{cases}$$

- Rewrite as $u = \operatorname{argmin} J(\theta)$ where

$$J(\theta) = \int_{x \in \Omega} |\mathcal{L}v_{\theta}^{\text{NN}}(x) - f(x)|^2 dx + \lambda \int_{x \in \Gamma} |v_{\theta}^{\text{NN}}(x) - g(x)|^2 d\sigma \quad (\lambda > 0).$$

- What they did is :
 - implement a modeler to design domains/shapes (Ω)
 - define functions with NN (layers and neurons)
 - use only smooth activation function (no ReLU stuff)
 - use Monte-Carlo to approximate integrals (on any Ω and Γ)
 - use autodiff to assemble $\mathcal{L}v_{\theta}^{\text{NN}}(x_i)$ at MC points x_i
 - use autodiff to minimize J and find quasi-optimal values θ_*

Weinan+Bing, The Deep Ritz method : A deep learning-based numerical algorithm for solving variational problems, 2017.

Solve elliptic problems like the model equation

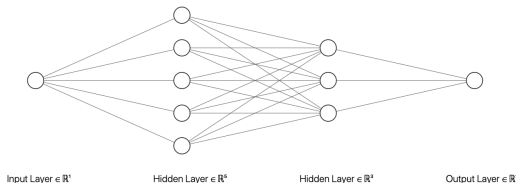
$$\begin{cases} -\Delta u + u = f, & x \in \Omega, \\ \partial_n u = 0, & x \in \partial\Omega, \end{cases}$$

using

$$u = \operatorname{argmin}_{v \in V} \int_{\Omega} \left(\frac{1}{2} |\nabla v(x)|^2 + \frac{1}{2} |v(x)|^2 - f(x)v(x) \right) dx, \quad V = H^1(\Omega).$$

-
- Lagaris+Likas+Fotiadis, Artificial Neural Networks for Solving Ordinary and PDEs, 1997.
 - Raissi+Perdikaris+Karniadakis. Physics-informed neural networks (PINNs), 2019.
 - Kharazmi+Zhang+Karniadakis, VPINNs : Variational Physics-Informed Neural Networks For Solving Partial Differential Equations, 2019,
 - Berrone+Canuto+Pintore. VPINNs quadratures and test functions, 2021.

- Step 1 : choose a NN architecture (hyper-parameters)



from which constructs functions $x \mapsto v_{\theta}^{\text{NN}}(x)$ where we renamed the weights $\theta = W$. It defines the space

$$V^{\text{NN}} = \{v_{\theta}^{\text{NN}} \text{ for all weights } \theta\}.$$

- Step 2 : minimize $u^{\text{NN}} = \underset{v_{\theta}^{\text{NN}} \in V^{\text{NN}}}{\operatorname{argmin}} J(\theta)$

$$J(\theta) = \int_{\Omega} \left(\frac{1}{2} |\nabla v_{\theta}^{\text{NN}}(x)|^2 + \frac{1}{2} |v_{\theta}^{\text{NN}}(x)|^2 - f(x) v_{\theta}^{\text{NN}}(x) \right) dx.$$

- Step 3 : replace $\int h(x) dx$ with $\sum h(x_i) w_i$: **Riemann, Monte Carlo, ...**

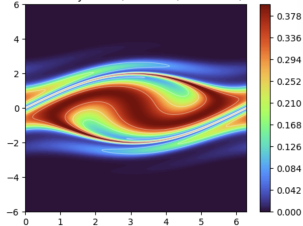
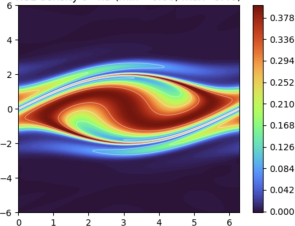
Extremely easy. Install :

```
pip3 install --upgrade pip
pip3 install --user virtualenv
virtualenv myenv2
source myenv2/bin/activate
pip install torch
pip install matplotlib
pip install jupyter
pip install scimba
pip install jax
```

Then run with either :

```
jupyter notebook
or :
python file.py
```

Let us run Jupyter Notebooks

`geometries.ipynb``gradshaf.ypnb``transport.ipynb`SL density $t=4.5$ (min=0.00, max=0.40)NSL density $t=4.5$ (min=-0.00, max=0.40)

Efficiency in high dimensions

dimen- sion d	classical SL				neural SL			
	n_x	Time	relative error		Time		relative error	
			e_{L^2}	e_{L^∞}	init.	iter.	e_{L^2}	e_{L^∞}
2	64	0.86 s	9.39×10^{-4}	1.15×10^{-2}	21.27 s	10.84 s	8.75×10^{-4}	1.14×10^{-2}
3	64	1.27 s	3.08×10^{-4}	1.09×10^{-2}	23.97 s	14.23 s	3.44×10^{-4}	1.10×10^{-2}
4	64	27.16 s	1.15×10^{-4}	9.42×10^{-3}	29.76 s	17.86 s	1.47×10^{-4}	9.26×10^{-3}
5	32	62.91 s	1.03×10^{-4}	1.58×10^{-2}	41.45 s	25.12 s	7.35×10^{-5}	8.02×10^{-3}
6	23	85.59 s	1.12×10^{-4}	1.75×10^{-2}	55.86 s	33.86 s	3.93×10^{-5}	4.09×10^{-3}
7	15	106.87 s	1.19×10^{-4}	3.82×10^{-2}	79.70 s	47.69 s	3.62×10^{-5}	5.34×10^{-3}
8	11	139.95 s	2.47×10^{-4}	1.08×10^{-1}	109.83 s	64.34 s	3.69×10^{-5}	3.70×10^{-3}

From :

Neural semi-Lagrangian method for high-dimensional advection-diffusion problems, 2026

<https://arxiv.org/pdf/2504.20715>

Extension to full Vlasov-Poisson at the end of the paper.

Comparison with approximation theory

Theorem (Cybenko 1989)

Let S be any sigmoid. Then finite sums of the form

$$G(\mathbf{x}) = \sum_{j=1}^N \alpha_j S(\langle \mathbf{y}_j, \mathbf{x} \rangle + \theta_j)$$

are dense in $C^0([0, 1]^d)$ equipped with the norm of the maximum.

Theorem (Yarostsky)

There exists a ReLU-NN approximating all bounded functions belonging to $W^{n, \infty}([0, 1]^d)$ with uniform accuracy ε and at most $O(\varepsilon^{-d/n} \log 1/\varepsilon)$ computational units.

Numerical results obtained with real algorithms are not well explained by such theoretical results.

Summary of that part

- SCIMBA makes Physically-Informed-Neural-Nets (just least square by the way) accessible.
- Clearly good potential and good accuracy for problems in very high dimension (thanks to MC integrator) and parametric studies.
- NN enhanced δf simulations in another work (Non-linear control variate in δf particle-in-cell methods using symplectic neural networks, Fournet, Campos-Pinto, Franck, Michel-Dansac, 2026).