A new quasi-linear transport model for electromagnetic turbulence in high-beta spherical tokamaks and its application to STEP

<u>M. Giacomin^{1,2,*}</u>, D. Dickinson², W. Dorland³, N. Mandel⁴, A. Bokshi², F. J. Casson⁵, H. G. Dudding⁵, D. Kennedy⁵, B. S. Patel⁵, and C. M. Roach⁵

¹Department of Physics, University of Padua, Padua, Italy.

²York Plasma Institute, University of York, York, YO10 5DD, United Kingdom.

³Department of Physics, University of Maryland, College Park, MD 20740, USA.

⁴Princeton Plasma Physics Laboratory, Princeton, 08543, NJ, USA.

⁵UK Atomic Energy Authority, Abington, OX14 3bD, United Kingdom.

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*maurizio.giacomin@unipd.it

Outline

- Why building a reduced transport model for STEP?
- The quasi-linear model
 - The quasi-linear metric $\boldsymbol{\Lambda}$
 - Quasi-linear fluxes
- Transport simulations with T3D
 - Implementation of the quasi-linear model in T3D
 - First flux-driven simulations of a STEP flat-top operating point
- Summary

For details see Giacomin et al arXiv:2404.17453v1, accepted for publication in JPP.





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D. Kennedy's talk showed that the STEP flat-top operating point obtained with JINTRAC modelling is <u>not</u> in kinetic equilibrium: gyrokinetic simulations predict large electromagnetic turbulent fluxes driven by kinetic ballooning modes.



What are our options?

1. Use other existing saturation rule implemented in available transport models such as TGLF [1] or QuaLiKiz [2, 3].

Fast solution but inaccurate for electromagnetic turbulence such as hybrid-KBMs in STEP.

2. Use high-fidelity transport models based on nonlinear gyrokinetic simulations, such as Tango-GENE [4], PORTALS-CGYRO [5], T3D-GX [6].

Very accurate solution but too expensive computationally.

3. Build an electromagnetic reduced transport model to specifically addressed STEP-like regimes.

- [2] C. Bourdelle et al. Plasma Phys. Control. Fusion 58 (2015).
- [3] J. Citrin et al. Plasma Phys. Control. Fusion 58 (2017).
- [4] A. Di Siena et al. Nucl. Fusion 62 (2022).
- [5] P. Rodriguez-Fernandez *et al.* Nucl. Fusion 62 (2022).
- [6] M. Barnes et al. Phys. Plasmas 17 (2010); T. Qian et al. Bulletin of APS (2022).

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^[1] G.M. Stabler et al. Phys. Plasmas 14 (2007).

A reduced transport model for high- β ST



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The quasi-linear reduced model is composed of:

- The quantity Λ General function
- The coefficients Q_0 and α Specific to STEP

Building the function Λ



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• Perpendicular wavenumber averaged along θ

$$k_{\perp}^{2}\rangle_{k_{y},\theta_{0},i} = \frac{\int \mathrm{d}\theta \ k_{\perp}^{2}(\theta,k_{y},\theta_{0}) \ J(\theta) \ |\chi_{i}(\theta)|^{2}}{\int \mathrm{d}\theta \ J(\theta) \ |\chi_{i}(\theta)|^{2}}$$

where $\chi_i \in \{e\delta\phi/T_e, \delta A_{\parallel}/(\rho_s B_0), \delta B_{\parallel}/B_0\}$

• Summing $\gamma/\langle k_{\perp}^2 \rangle$ over all the three fields

$$\hat{\Lambda}(k_y,\theta_0) = \sum_{i=1}^{3} \frac{\max_{\theta} |\chi_i(k_y,\theta_0,\theta)|}{\max_{\theta} |\chi_1(k_y,\theta_0,\theta)|} \frac{\gamma(k_y,\theta_0)}{\langle k_{\perp}^2(k_y,\theta_0,\theta) \rangle_{\theta,i}}$$

• Integrating over θ_0

$$\bar{\Lambda}(k_y) = \frac{1}{\theta_{0,\max}(k_y,\gamma_E)} \int_0^{\theta_{0,\max}(k_y,\gamma_E)} \hat{\Lambda}(k_y,\theta_0) \,\mathrm{d}\theta_0$$

The reduced model retains the effect of the equilibrium flow shear



- A linearly growing mode in presence of flow shear is "tilted" in time, $k_x = \gamma_E k_y \Delta t$.
- The value of k_x increases to $k_x = \gamma_E k_y / \gamma$ after a growth time.
- k_x is related to θ_0 as $k_x = \hat{s}\theta_0 k_y$ (in ballooning space)
- The value of θ_0 reached after a growth time (and limited to π) is therefore

$$heta_{0,\max} = \min\left(rac{\gamma_E}{\hat{s}\gamma}, \ \pi
ight)$$

- Small values of γ_E are sufficient to reduce the value of Λ if $\gamma(\theta_0)$ decreases with θ_0 .
- The average over $\theta_0 \in [0, \theta_{0,max}]$ reduces to the value at $\theta_0 = 0$ at $\gamma_E = 0$.

Building the function Λ



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• Perpendicular wavenumber averaged along θ

$$\langle k_{\perp}^2 \rangle_{k_y,\theta_0,i} = \frac{\int \mathrm{d}\theta \ k_{\perp}^2(\theta,k_y,\theta_0) \ J(\theta) \ |\chi_i(\theta)|^2}{\int \mathrm{d}\theta \ J(\theta) |\chi_i(\theta)|^2}$$

where $\chi_i \in \{e\delta\phi/T_e, \delta A_{\parallel}/(\rho_s B_0), \delta B_{\parallel}/B_0\}$

• Summing $\gamma/\langle k_{\perp}^2 \rangle$ over all the three fields

$$\hat{\Lambda}(k_y,\theta_0) = \sum_{i=1}^{3} \frac{\max_{\theta} |\chi_i(k_y,\theta_0,\theta)|}{\max_{\theta} |\chi_1(k_y,\theta_0,\theta)|} \frac{\gamma(k_y,\theta_0)}{\langle k_{\perp}^2(k_y,\theta_0,\theta) \rangle_{\theta,i}}$$

• Integrating over θ_0

$$\bar{\Lambda}(k_y) = \frac{1}{\theta_{0,\max}(k_y,\gamma_E)} \int_0^{\theta_{0,\max}(k_y,\gamma_E)} \hat{\Lambda}(k_y,\theta_0) \,\mathrm{d}\theta_0$$

• Finally integrating over k_y

$$\Lambda = \int \mathrm{d}k_y \,\bar{\Lambda}(k_y) \qquad \qquad Q_{\mathrm{ql}} = Q_0 \Lambda^\alpha$$

The quasi-linear model



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Linear flux ratio (species over total) as quasi-linear weight (similar for particle flux)

$$Q_{ql,s} = \frac{Q_0 \Lambda^{\alpha-1}}{\rho_* c_s} \int dk_y \frac{1}{\theta_{0,\max}} \int_0^{\theta_{0,\max}} d\theta_0 \frac{Q_{l,s}(k_y,\theta_0)}{Q_{l}(k_y,\theta_0)} \sum_{i=1}^3 \frac{\max_{\theta} |\chi_i(k_y,\theta_0,\theta)|}{\max_{\theta} |e\delta\phi(k_y,\theta_0,\theta)/T_e|} \frac{\gamma(k_y,\theta_0)}{\langle k_{\perp}^2(k_y,\theta_0,\theta) \rangle_{\theta,i}}$$
Normalizing factor Equilibrium flow shear suppression Weighted sum over each field contribution

$$\Lambda = \frac{1}{\rho_* c_s} \int \mathrm{d}k_y \, \frac{1}{\theta_{0,\max}} \int_0^{\theta_{0,\max}} \mathrm{d}\theta_0 \, \sum_{i=1}^3 \frac{\max_{\theta} |\chi_i(k_y,\theta_0,\theta)|}{\max_{\theta} |e\delta\phi(k_y,\theta_0,\theta)/T_e|} \frac{\gamma(k_y,\theta_0)}{\langle k_{\perp}^2(k_y,\theta_0,\theta) \rangle_{\theta,i}} \quad \Longrightarrow \quad \sum_s Q_{\mathrm{ql},\mathrm{s}} = Q_0 \Lambda^{\alpha}$$

The coefficients Q_0 and lpha



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- The function $\Lambda(k_y, \theta_0)$ is computed directly from linear GK simulations.
- The coefficients Q_0 and α link Λ to the quasi-linear total heat flux.
- How are these coefficients computed? From a database of nonlinear GK STEP simulations!

The database includes heat flux values from about 20 nonlinear simulations with and without the equilibrium flow shear, different pressure gradients, β (and β'), \hat{s} and q.



Implementation in T3D



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- T3D [1] is an improved version of Trinity [2], written in Python with a highly modular structure.
- Transport equations solved by T3D:

$$\frac{\partial n_s}{\partial t} + \frac{1}{V'} \frac{\partial}{\partial \Psi} (V' \overline{\Gamma_s}) = \overline{S_{n,s}}$$
$$\frac{3}{2} \frac{\partial p_s}{\partial t} + \frac{1}{V'} \frac{\partial}{\partial \Psi} (V' \overline{Q_s}) = \frac{3}{2} n_s \sum_u \nu_{su} (T_u - T_s) + \overline{S_{p,s}}$$

- Transport equations require the value of the fluxes Γ_s and Q_s , which are computed from the reduced transport model.
- T3D is coupled to GS2 to compute the function Λ .
- Different neoclassical models are available in T3D. Here we use NEO [3].

[1] T. Qian *et al.* Bulletin of APS (2022).

[2] M. Barnes et al. Phys. Plasmas 17 (2010).

[3] E.A. Belli & J. Candy et al. Plasma Phys. Control. Fusion 50 (2008).

Numerical setup and initial state





- The initial state is the STEP-EC-HD from JINTRAC-JETTO [1].
- 150 MW ECHR, 340 MW radiated power, 360 MW initial alpha heating.
- Particle source from pellet injection.
- Shearing rate from neoclassical radial electric field (JINTRAC-JETTO).
- $(N_p+1)(N_r-1)n_{k_v}n_{\theta_0} = 1440$ single mode linear GS2 simulations per T3D iteration.
- Finite Dirichlet boundary condition at r/a = 0.9.

T3D-GS2 numerical parameters	
N _r	6
r _{edge} /a	0.9
N _p	3
n _{ky}	12
n _{θo}	6
n _{min}	2
n _θ	33
n _λ	24
n _e	10

There are three main approximations



- 1. Only density and temperature profiles of electrons and a single thermal ion species are evolved. The quasi-neutrality is enforced on the ion species, such that $n_i = n_D + n_T = n_e$, where $n_D = n_T = n_i/2$. Impurities and fast α particles are neglected.
- 2. The geometrical shaping parameters, q and \hat{s} are considered constant while evolving the pressure profile. The parameter β' is evolved with the pressure profile.
- 3. The **equilibrium flow** shear is computed from the initial radial electric field evaluated from the ion pressure gradient and neoclassical flows, and it **is kept constant** as the kinetic profiles evolve. Sensitivity to γ_E is evaluated.

T3D steady state solution





Initial strong power and particle imbalance

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Initial power and particle balance Power and particle losses largely exceed available sources.

Final power and particle balance

Excellent agreement between power and particle losses and available sources.







r/a

Time evolution of power and particle loss





- Evolution divided in three phases:
 - Short initial phase (t < 0.1 s, 50 iterations): power and particle losses decrease suddenly.
 - Long second phase (t < 9 s, 100 iterations): particle and power losses converge to available sources.
 - Final phase (t > 9 s, 40 iterations): particle and power losses are constant.
- Presence of some ripple from low k_y modes being stabilized/destabilized.
- Particle flux vanishes at $r/a \le 0.33$.
- Long time is required to relax to a steady state solution (particle confinement time is 12.5 s).

Fusion power reduced by 10%



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- Fusion power is 10% smaller in the T3D solution than in the initial STEP operating point. The fusion power drops is just an indicative value. Approximations may be important.
- Energy and confinement time are similar to STEP-EC-HD.

Are the T3D-GS2 final profiles a steady state solution for GK turbulent transport?

Linear GK analysis is reassuring



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The T3D steady state is characterized by GK instabilities similar to STEP-EC-HD.





Growth rate values are in general smaller, especially as low k_y .

The T3D solution is a steady state for GK



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Nonlinear GK simulations predict heat and particle fluxes that are equal or below the available sources in the T3D steady state solution



Summary



- Confinement prediction in STEP requires modelling turbulent transport from hybrid-KBMs, but GK turbulent transport calculations are prohibitively expensive.
- A new quasi-linear transport model is built for hybrid-KBM transport in STEP-like regimes and implemented in the T3D transport code.
- The first STEP T3D simulations return plasma profiles similar to the initial condition with fusion power reduced by 10%.
- A number of approximations are considered: no fast α particles, no impurities and constant equilibrium.
- GK nonlinear simulations show that the T3D solution is compatible with the available STEP sources.

Future work

- Effect of impurities and fast α particles.
- Stochastic transport from MTMs.
- Sensitivity on the safety factor or evolution of the equilibrium with kinetic profiles.
- Sensitivity to boundary condition.
- Experimental validation in high- β spherical tokamaks.
- Surrogate model for Λ .