Full Flux Surface (FFS) δf -gyrokinetic code; stella

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Vienna Workshop

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- 2. Tokamak Formalism
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- 4. Full Flux Surface Code
- 5. Results
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Motivation

To get fusion working

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Tokamak formalism



Tokamak formalism



- ▶ Magnetic field geometry described by Clebsch representation: $B = \nabla \alpha \times \nabla \psi$
- Cooridinate system: $(\psi, \theta, \zeta) \leftrightarrow (x, y, z)$
- $\blacktriangleright \ \psi \leftrightarrow \text{flux surface label} \leftrightarrow x$
- $\bullet \ \theta \leftrightarrow \text{poloidal angle} \leftrightarrow z \text{ (for tokamak)}$
- $\zeta \leftrightarrow$ toroidal angle $\leftrightarrow z$ (for stellarators)

$$\blacktriangleright \ \alpha = \theta - \iota \zeta \leftrightarrow y$$

Tokamak formalism



- ▶ Despite fusion being only 10 years away, δf -gyrokinetic simulations for full devices are out of numerical reach
- ▶ Consider just a single field line and simulate turbulence along it

Current state of affairs



▶ Simulation coordinates: $(x, y, z) \to (\psi, \alpha, \zeta)$

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- ▶ Initialise some δf and ϕ at t = 0 on simulation domain
- Evolve gyrokinetic equation pseudo-spectrally

Current state of affairs



- Simulation coordinates: $(x, y, z) \to (\psi, \alpha, \zeta)$
- ▶ Initialise some δf and ϕ at t = 0 on simulation domain
- Evolve gyrokinetic equation pseudo-spectrally
- ▶ Decay in v_{\parallel} ; $g(t, \boldsymbol{x}, v_{\parallel} \to \pm \infty, \mu) \to 0$
- ▶ Turbulence is taken as periodic in perpendicular directions, $k_x, k_y \gg 1/L$
- ▶ Use twist and shift boundary conditions in z to capture extended modes





▶ If $\hat{s} \propto dq/d\psi \neq 0$ then domain gets sheared as it travels around device



- ▶ If $\hat{s} \propto dq/d\psi \neq 0$ then domain gets sheared as it travels around device
- Eddies get sheared
- Pushed to higher perpendicular wavenumbers



▶ Use "twist-and-shift" boundary conditions to map sheared domain back onto original one



▶ Demand that any function, $A(t, x, y, z) = \sum_k \hat{A}_{k_x, k_y}(t, z) e^{ik_y(y-y_0) + ik_x(x-x_0)}$, be periodic at the same poloidal location, θ



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$$But \ y = y(\psi(x), \theta, \zeta) = \zeta - q(\psi)\theta A(t, x, y(x, \theta, z), z) = A(t, x, y(x, \theta, z + 2p\pi), z + 2p\pi)$$
(1)



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• Let \mathbf{k}' be the wavenumber that satisfies the above: $\hat{A}_{\mathbf{k}}(z) = \hat{A}_{\mathbf{k}'}(z + 2p\pi)$ (phase factor)

- $\{k'_x, k'_y\} = \{k_x + \delta k_x, k_y\}$ $\delta k_x \propto \frac{\mathrm{d}y}{\mathrm{d}x}, \ \delta k_x = 2\pi \hat{s} k_y$
- phase factor = $-2\pi pik_y \frac{\mathrm{d}y}{\mathrm{d}\alpha}$

(2)



• Demand that any function, $A(t, x, y, z) = \sum_k \hat{A}_{k_x, k_y}(t, z) e^{ik_y(y-y_0) + ik_x(x-x_0)}$, be periodic at the same poloidal location, θ

$$\textbf{But } y = y(\psi(x), \theta, \zeta) = \zeta - q(\psi)\theta \\ A(t, x, y(x, \theta, z), z) = A(t, x, y(x, \theta, z + 2p\pi), z + 2p\pi)$$
(1)

- \blacktriangleright Let k' be the wavenumber that satisfies the above: $\hat{A}_{k}(z) = \hat{A}_{k'}(z+2p\pi)$ (phase factor) (2)
- $\flat \delta k_x = 2\pi \hat{s} k_u$
- $\blacktriangleright \delta k_x = N \Delta k_x$
- $\triangleright \ 2\pi \hat{s}k_y = N\Delta k_x$ ▶ $j_{\text{twist}} = \frac{2\pi \hat{s} L_x}{L_y} \in \mathbb{Z}$

Tokamaks are trivial



 Fluxtube simulations are sufficient for Tokamak because we can stitch our fluxtubes together

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Stepping outside the axisymmetric box



- Geometry sampled varies with field line
- Method of stitching together fluxtubes no longer holds
- \blacktriangleright Currently no δf poloidally-global code that can correctly deal with kinetic electrons





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-----> kz S;

• We have a k_y spectrum for a given fieldlines





▶ Different fieldlines have different spectra



▶ Full flux effects should sample all fieldlines

Geometry is at the root of all evil...

Geometry is no longer trivial

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▶ But how does geometry enter our code?

Geometry is at the root of all evil...

- Geometry is no longer trivial
- ▶ But how does geometry enter our code?

Simplified notation GK:

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$$\frac{\partial g}{\partial t} = (\text{geometric factors}) \cdot (\nabla g + \nabla \langle \phi \rangle_{\mathbf{R}})$$
(3)

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$$\frac{\partial g}{\partial t} = \underbrace{(\text{geometric factors})}_{\text{e.g} \,\hat{\boldsymbol{b}} \cdot \boldsymbol{\nabla} \boldsymbol{z}} \cdot (\boldsymbol{\nabla} g + \boldsymbol{\nabla} \underbrace{\langle \boldsymbol{\phi} \rangle_{\boldsymbol{R}}}_{J_{0,\boldsymbol{k}} \,\hat{\boldsymbol{\phi}}_{\boldsymbol{k}}}) \tag{3}$$

- ▶ Bessel functions $J_0(a_k)$ with $a_k = \frac{k_\perp v_\perp}{\Omega_s}$
- ▶ Geometric factors are α -dependent

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▶ Bessel functions $J_0(a_k)$ with $a_k = \frac{k_\perp v_\perp}{\Omega_s}$

• Geometric factors are α -dependent

▶ Gyro-averaging introduces coupling between different k_y -modes → no longer a local operation. But want to retain spectral accuracy

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$$\frac{\partial g}{\partial t} = \underbrace{(\text{geometric factors})}_{\mathbf{e}, \mathbf{g} \ \hat{\mathbf{b}} \cdot \nabla z} \cdot (\nabla g + \nabla \underbrace{\langle \phi \rangle_R}_{J_{0,k} \ \hat{\phi}_k}) \tag{3}$$

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▶ Geometric factors are α -dependent

- Gyro-averaging introduces coupling between different k_y -modes \rightarrow no longer a local operation. But want to retain spectral accuracy
- Geometric factors are introducing non-linear terms in k_y

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• Geometric factors are α -dependent

- Gyro-averaging introduces coupling between different k_y -modes \rightarrow no longer a local operation
- Geometric factors are introducing non-linear terms in k_y
- ▶ Domain changes for each fieldline

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▶ Fourier Decompose functions once at the beginning of code

$$\langle \phi \rangle_{\mathbf{R}} = \sum_{\mathbf{k}^{\prime\prime}} e^{i\mathbf{k}^{\prime\prime} \cdot \mathbf{R}} J_0(a_{\mathbf{k}^{\prime\prime},s}) \hat{\phi}_{\mathbf{k}^{\prime\prime}}, \qquad (4)$$

$$(\langle \phi \rangle_{\mathbf{R}})_{\mathbf{k}} = \int d^2 \mathbf{R} \sum_{\mathbf{k}^{\prime\prime}, k_{\alpha}^{\prime}} e^{i \left(k_{\psi}^{\prime\prime} - k_{\psi}\right) \psi} e^{i \left(k_{\alpha}^{\prime} + k_{\alpha}^{\prime\prime} - k_{\alpha}\right) \alpha} \hat{J}_{\mathbf{k}^{\prime\prime}, k_{\alpha}^{\prime}, s} \hat{\phi}_{\mathbf{k}^{\prime\prime}}, \tag{5}$$

where we have used

$$J_0(a_{\mathbf{k}^{\prime\prime},s}) = \sum_{k_{\alpha}^{\prime}} \hat{J}_{\mathbf{k}^{\prime\prime},k_{\alpha}^{\prime},s}(z,\mu) e^{ik_{\alpha}^{\prime}\alpha}.$$
(6)

- ▶ Problem: A full flux δf -gyrokinetic code couples together all k_y 's
- Solution: Michael Barnes
- ▶ Bessel functions, QN
- Geometric factors
- Non-constant domain
- Work with $\bar{g} = \frac{g}{F_0}$
- Multiply coefficients in real space
- Modify implicit terms

- ▶ Problem: A full flux δf -gyrokinetic code couples together all k_y 's
- Solution: Michael Barnes
- ▶ Bessel functions, QN
- ▶ Geometric factors
- Non-constant domain
- Fixed 2π poloidal domain.in θ , now set Δk_y using ρ_*
- ▶ Fieldlines will vary in path length around device
- Modify coordinates accordingly, and make "large aspect ratio approximation" for terms treated implicitly

$$\hat{b} \cdot \boldsymbol{\nabla} z = \langle \hat{b} \cdot \boldsymbol{\nabla} z \rangle_{\alpha} + \left[\hat{b} \cdot \boldsymbol{\nabla} z - \langle \hat{b} \cdot \boldsymbol{\nabla} z \rangle_{\alpha} \right]$$
(4)

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Rhostar scans



Sanity Check

▶ Miller Geometry:



Miller Nonlinear, kinetic electrons

Sprinkle in some geometry



Sprinkle in some geometry



▶ W7-X, nonlinear with kinetic electrons



WARNING

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Zonal modes may not be being treated correctly



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Why do we care?

Want to answer some global questions

- ▶ Look at effects on transport
- ▶ Look at global modes, e.g. TEM

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- ▶ Look at effects on transport
- ▶ Look at global modes, e.g. TEM
- ▶ Can stellar ators support zonal flows





Summary

- \blacktriangleright Aim for poloidally-global δf -gyrokinetic code
- ▶ Making progress but zonal flows are being dealt with incorrectly

Thank You