

# Correlation integrals and how to use them: a discussion

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Hosking & Schekochihin (2021), *Phys. Rev. X* **11**, 041005

Hosking & Schekochihin (2022), arXiv:2203.03573

Hosking & Schekochihin (2022), arXiv:2202.00462

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Fluid (and kinetic) systems support the conservation of certain quantities, i.e.,

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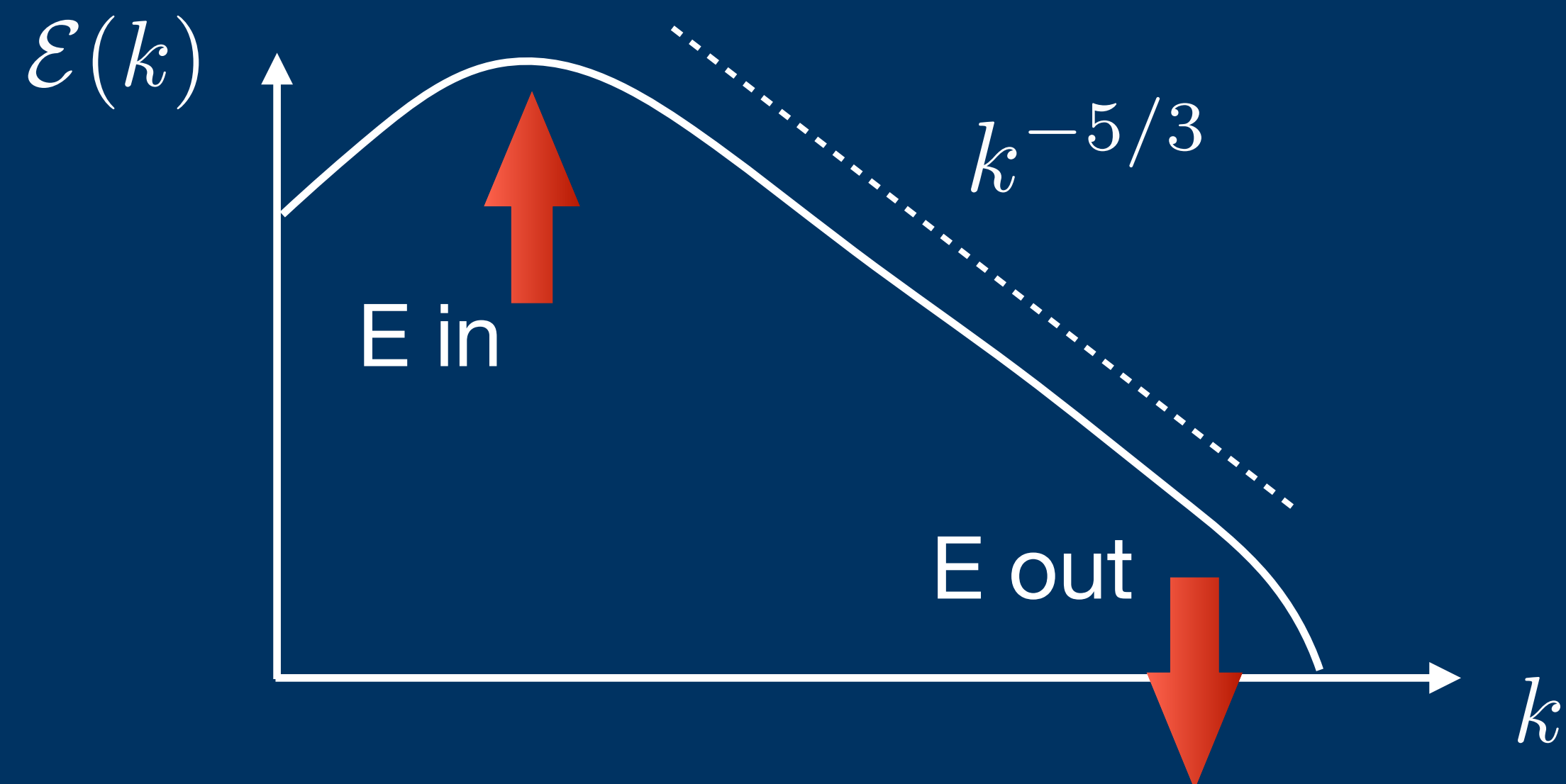
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At some level, turbulence is just the nonlinear processing of these quantities between injection and dissipation scales.



# Turbulence and conservation laws

The conserved quantities  $\psi$  come in two distinct types.

**Type 1:**  $\psi$  is a generalised energy, i.e., the norm of a dynamical field.  $\langle \psi \rangle > 0$  for any possible realisation of the system. Examples are:

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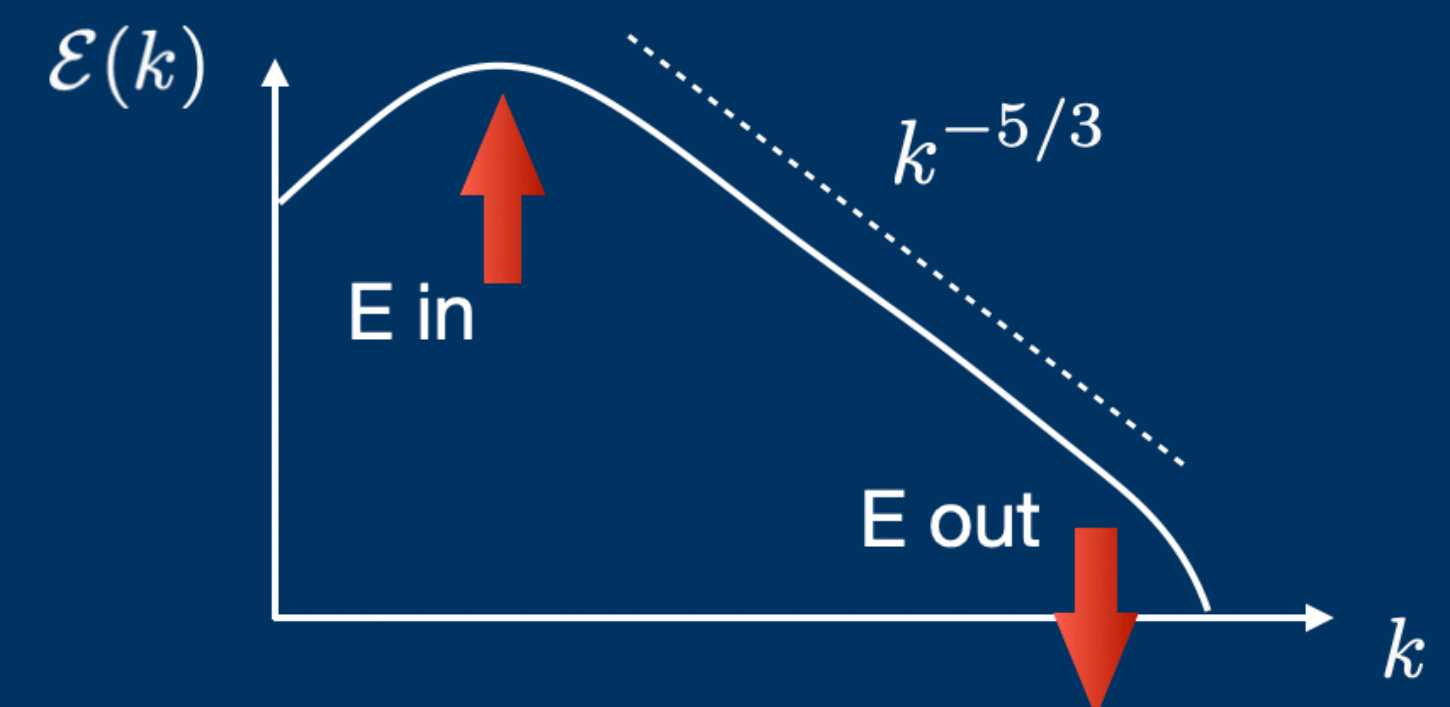
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Properties of generalised energies: cascades

$$\varepsilon_\psi \sim \frac{\psi_l}{\tau_l} \sim \text{const} \implies \mathcal{E}(k)$$

e.g., 3D hydro:  
(K41):

$$\varepsilon_E \sim u_l^2 \frac{u_l}{l} \implies u_l \sim (\varepsilon_E l)^{1/3} \implies \mathcal{E}(k) \propto k^{-5/3}$$



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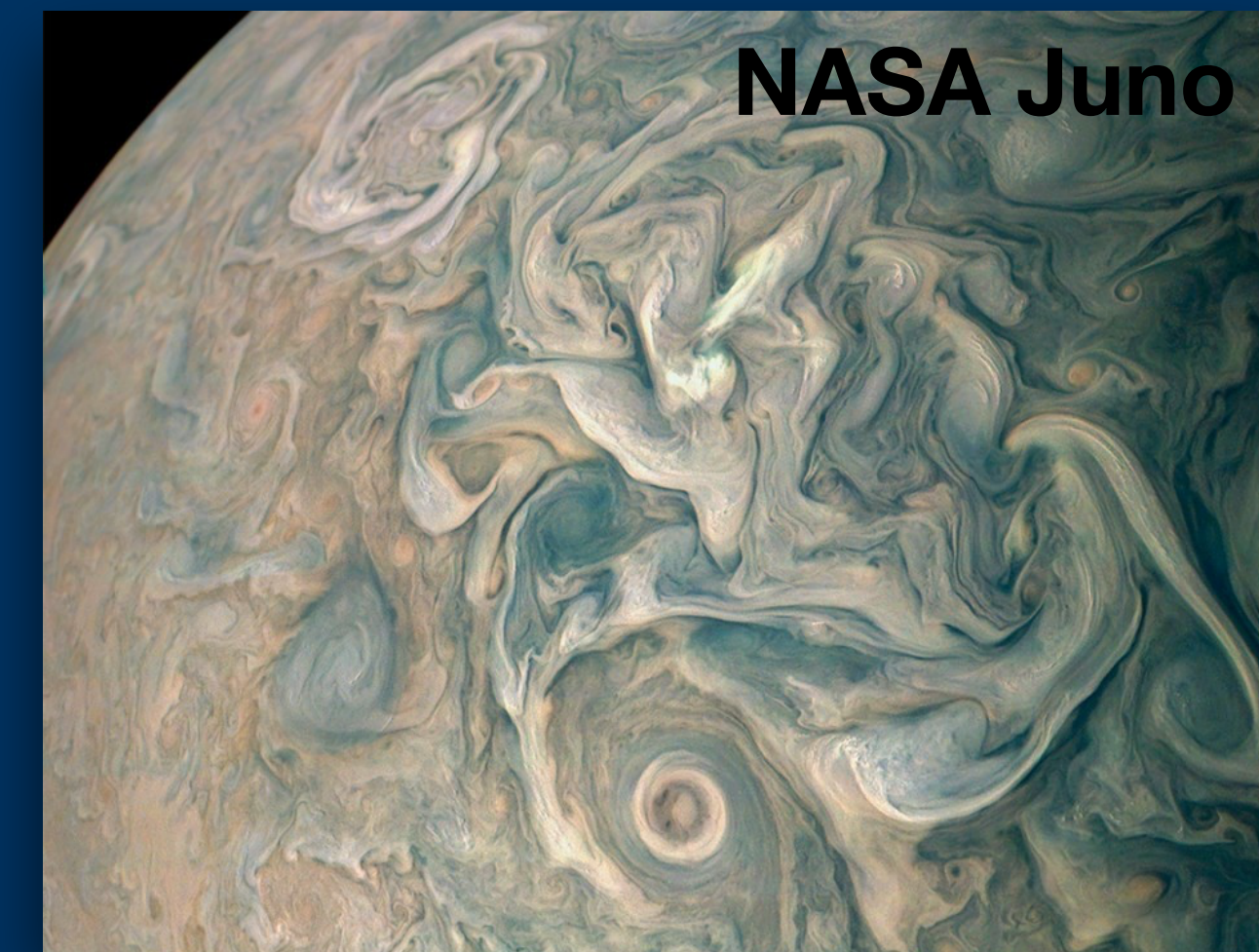
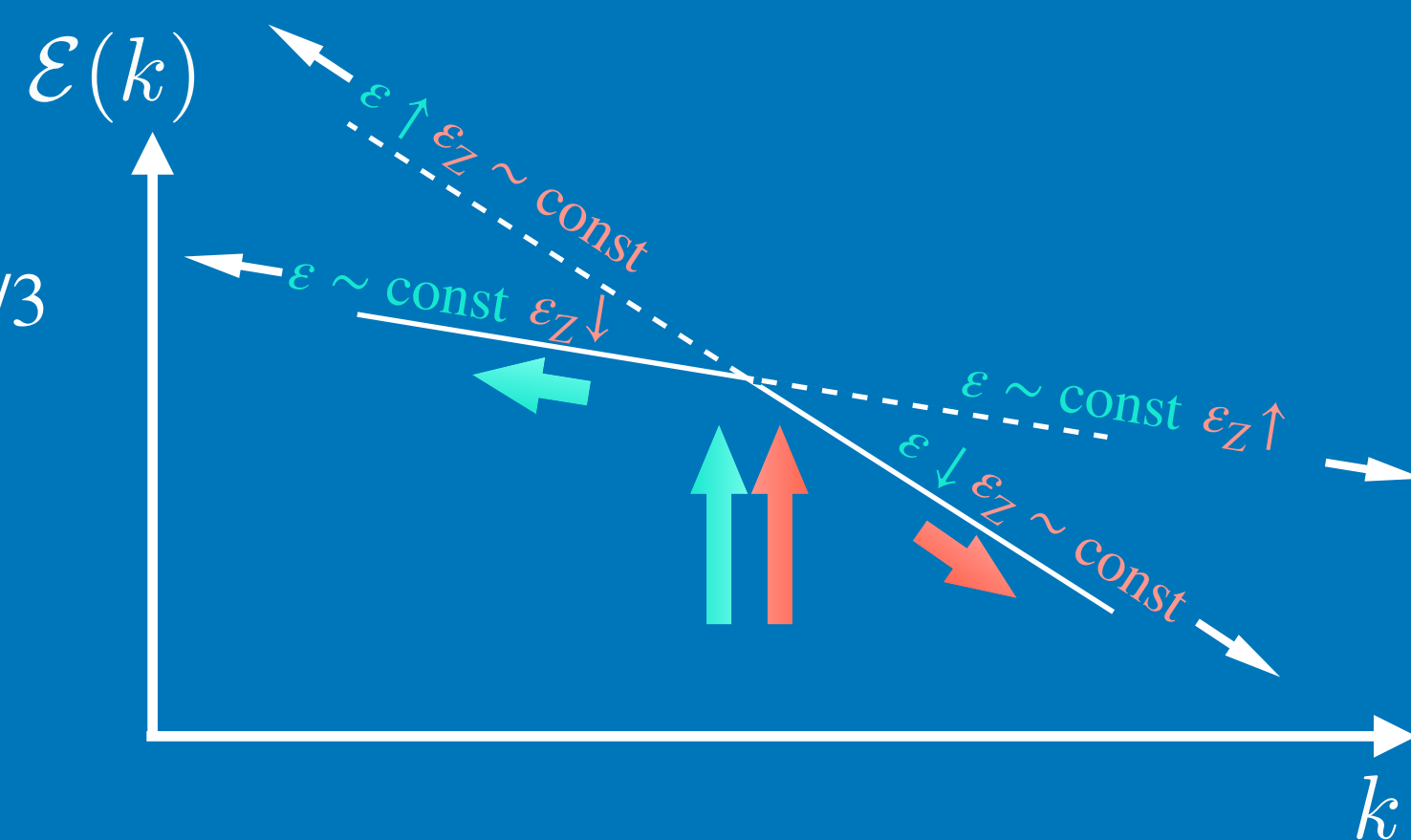
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$$\varepsilon_Z \sim \frac{u_l^2}{l^2} \frac{u_l}{l} \implies u_l \sim \varepsilon_Z^{1/3} l \implies \mathcal{E}(k) \propto k^{-3}$$





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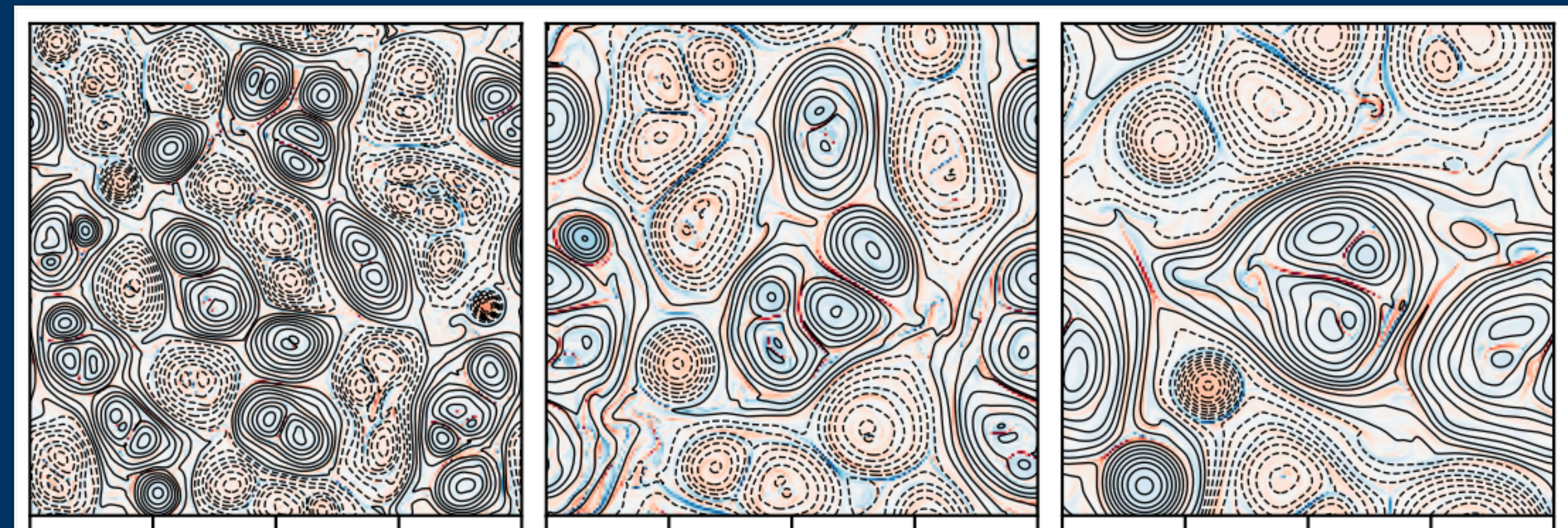
Properties of generalised energies: constrain turbulent decay

Zhou *et. al.* 2019

$$\frac{dE}{dt} \sim -\frac{E}{\tau_L}, \quad \langle \psi \rangle \sim U^\alpha L^\beta \sim \text{const} \implies E \propto t^{-p}$$

e.g., 2D MHD:

$$\tau_L \sim \frac{L}{B} \epsilon_{\text{rec}}^{-1}, \quad \langle A_z^2 \rangle \sim B^2 L^2 \sim \text{const} \implies B^2 \propto t^{-1}$$





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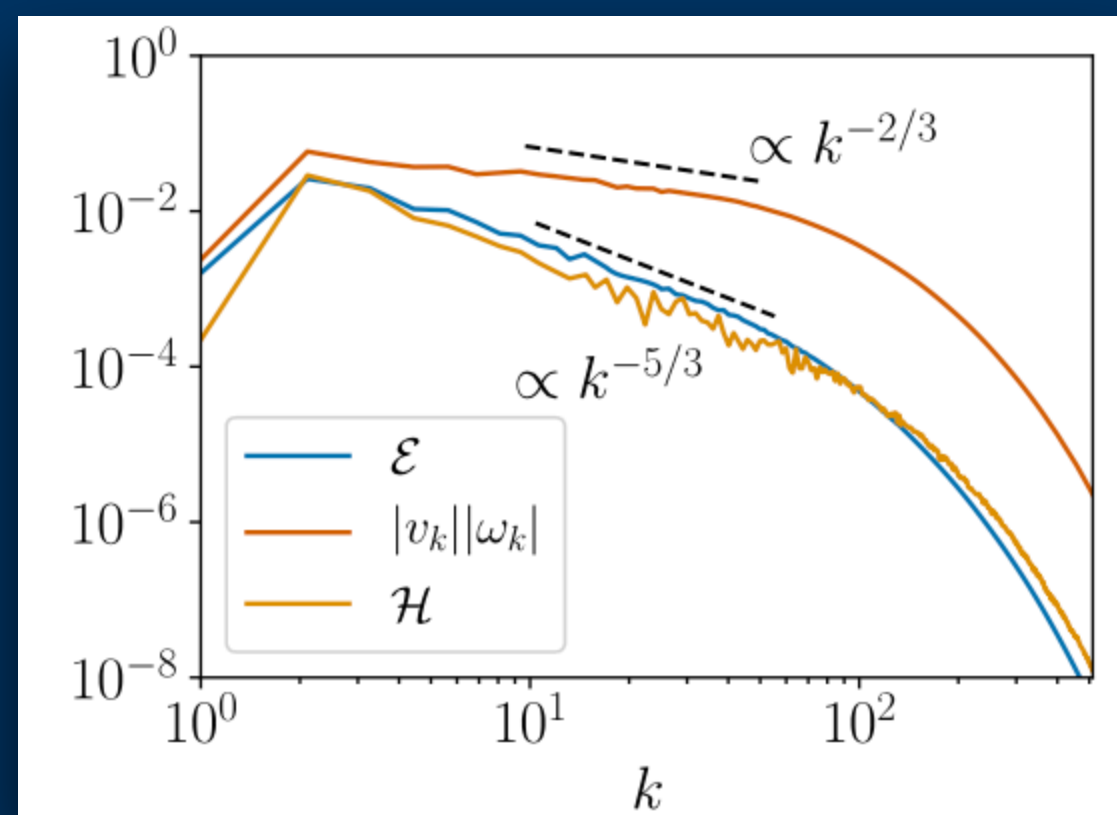
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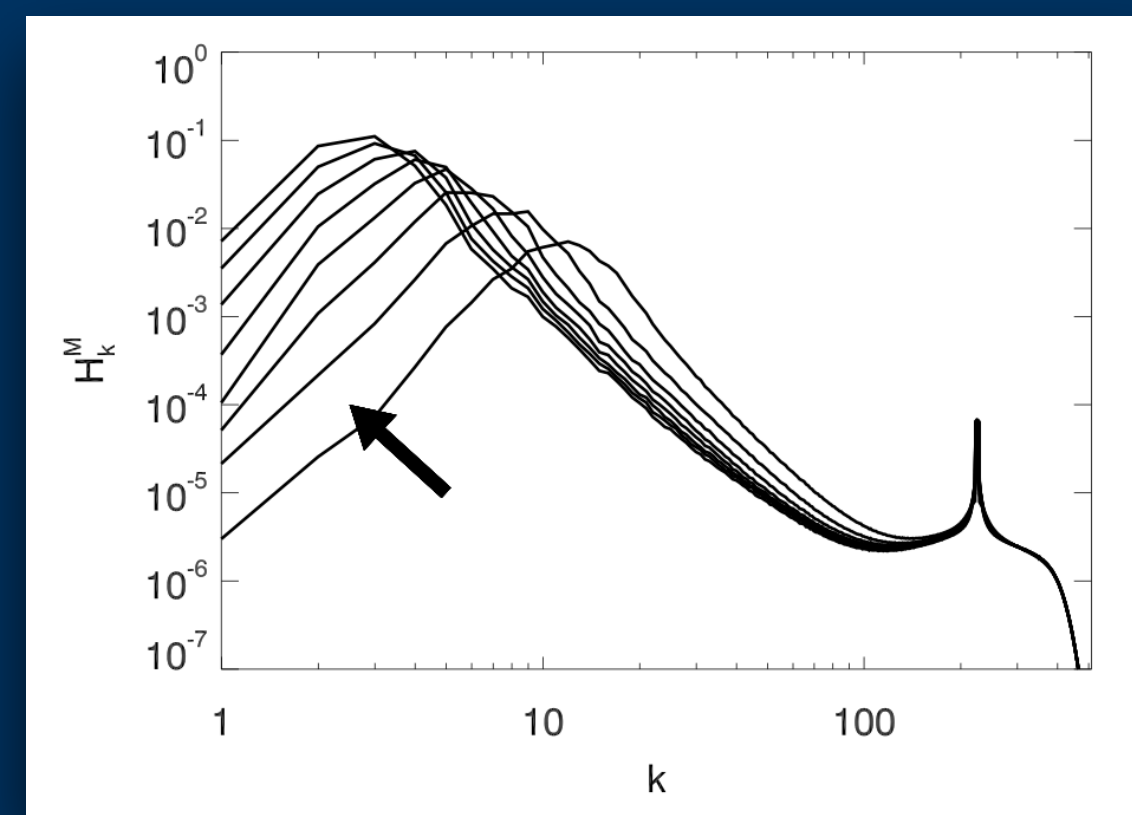
Properties of directed conserved quantities for finite  $\langle \psi \rangle$ :

Milanese et. al. 2022



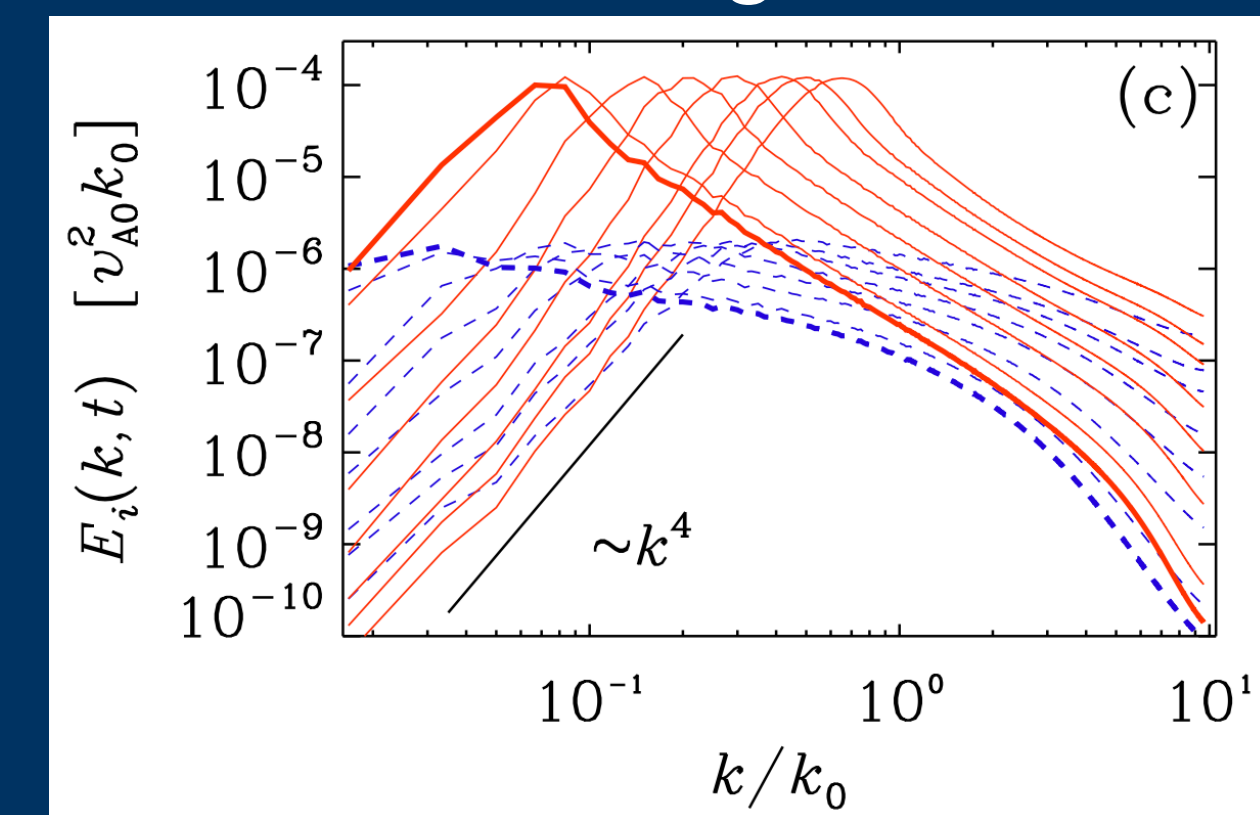
Cascades

Müller et. al. 2012



Dual/inverse cascade/transfer

Brandenburg et. al. 2017



Constrain turbulent decay

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→ correlation integrals.

# Correlation integrals

$$\frac{\partial \psi}{\partial t} + \nabla \cdot \mathbf{F} \simeq 0 \quad \Longrightarrow \quad \frac{\partial}{\partial t} \langle \psi(\mathbf{x}) \psi(\mathbf{x} + \mathbf{r}) \rangle + 2 \frac{\partial}{\partial \mathbf{r}} \cdot \langle \psi(\mathbf{x}) \mathbf{F}(\mathbf{x} + \mathbf{r}) \rangle \simeq 0$$

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$$\Longrightarrow \quad \frac{d}{dt} \underbrace{\int d^3 r \langle \psi(\mathbf{x}) \psi(\mathbf{x} + \mathbf{r}) \rangle}_{\equiv I_\psi} \simeq 0$$

Some examples:

$\psi = \mathbf{u}$ , Saffman (1967) integral

$\psi = \mathbf{r} \times \mathbf{u}$ , Loitsyansky (1939) integral\*

$\psi = \mathbf{u} \cdot \boldsymbol{\omega}$ , “Levich-Tsinober (1983) integral”

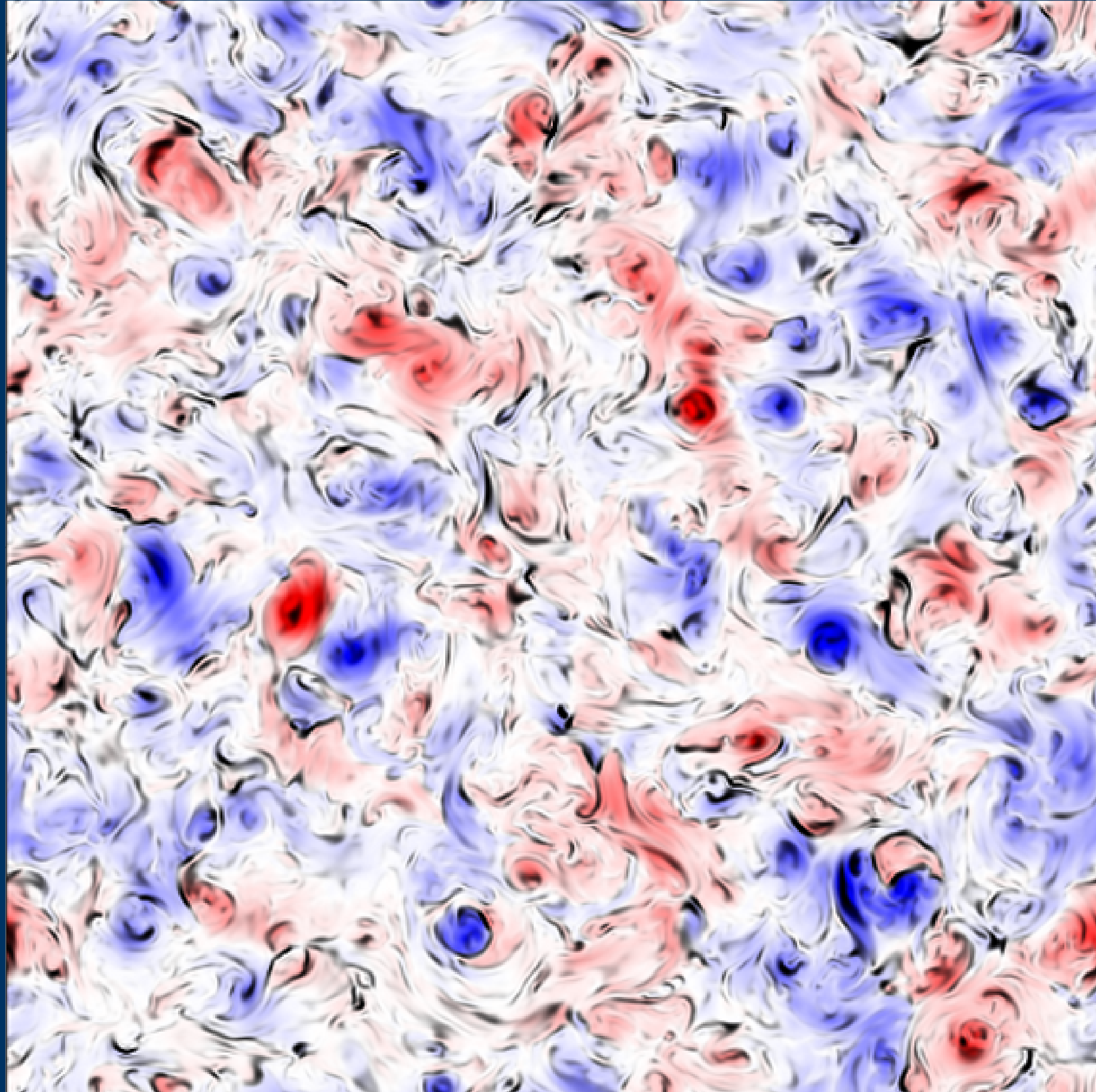
$\psi = \theta$ , Corrsin (1950) integral

$\psi = \mathbf{B}$ , “magnetic Saffman integral”

$\psi = \mathbf{A} \cdot \mathbf{B}$ , “H-S (2021) integral”

$\psi = \mathbf{u} \cdot \mathbf{B}$ , “Bershadskii (2019) integral”

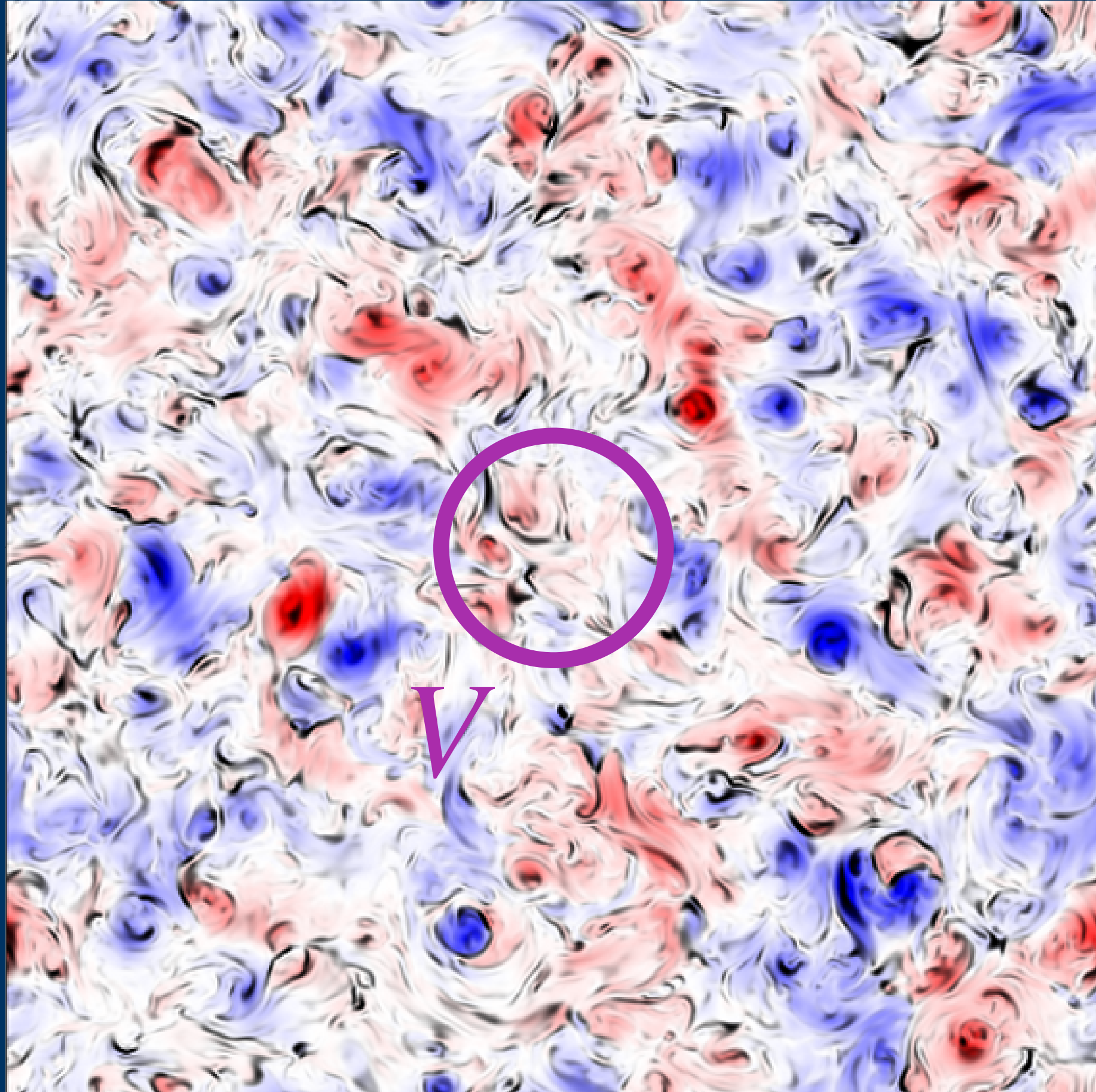
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$$I_\psi = \int d^3r \langle \psi(\mathbf{x}) \psi(\mathbf{x} + \mathbf{r}) \rangle = \lim_{V \rightarrow \infty} \frac{\left\langle \left[ \int_V d^3x \psi(\mathbf{x}) \right]^2 \right\rangle}{V}$$

2D slice of  $\psi = h = \mathbf{A} \cdot \mathbf{B}$  from simulation of MHD decay

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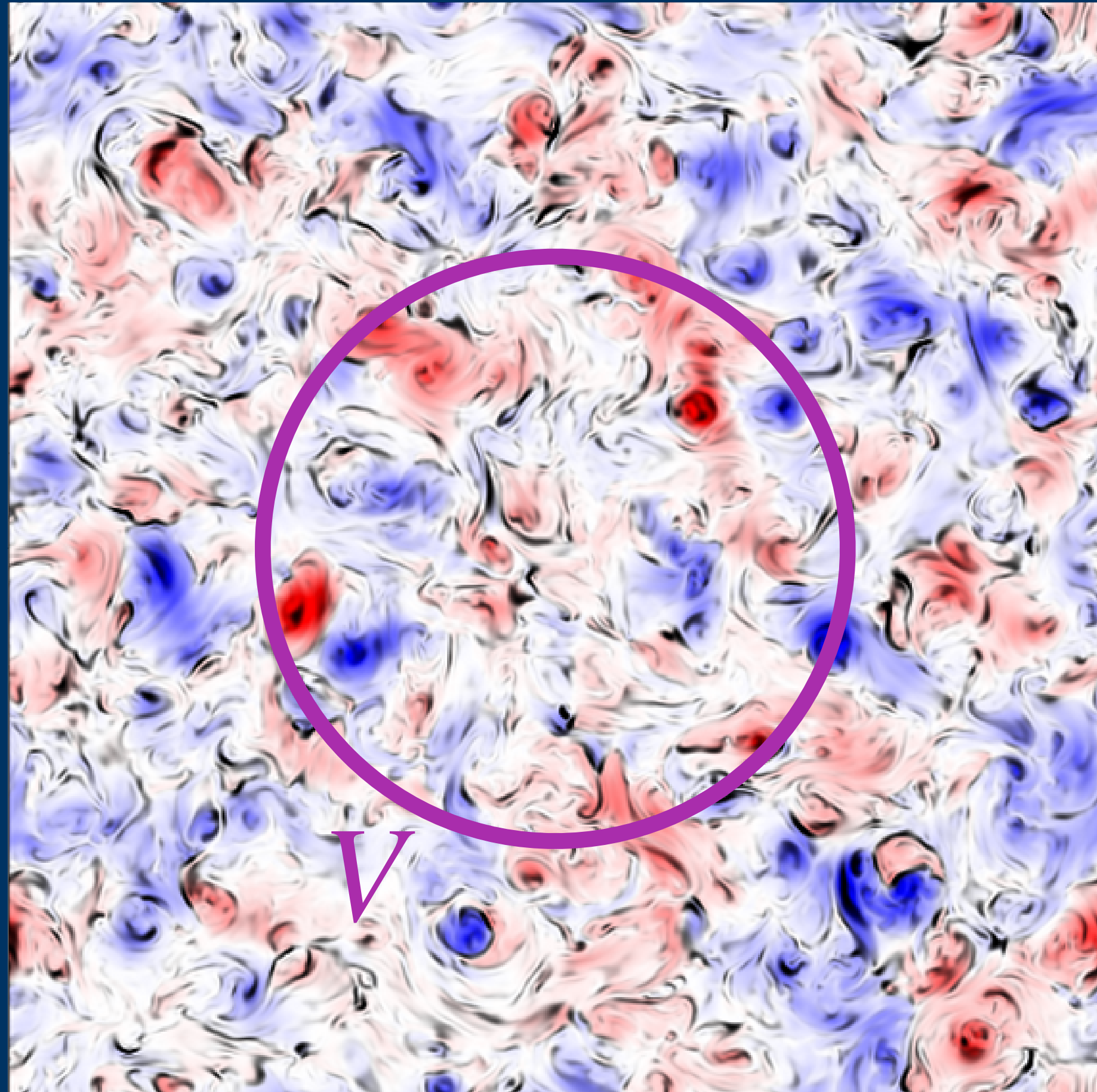


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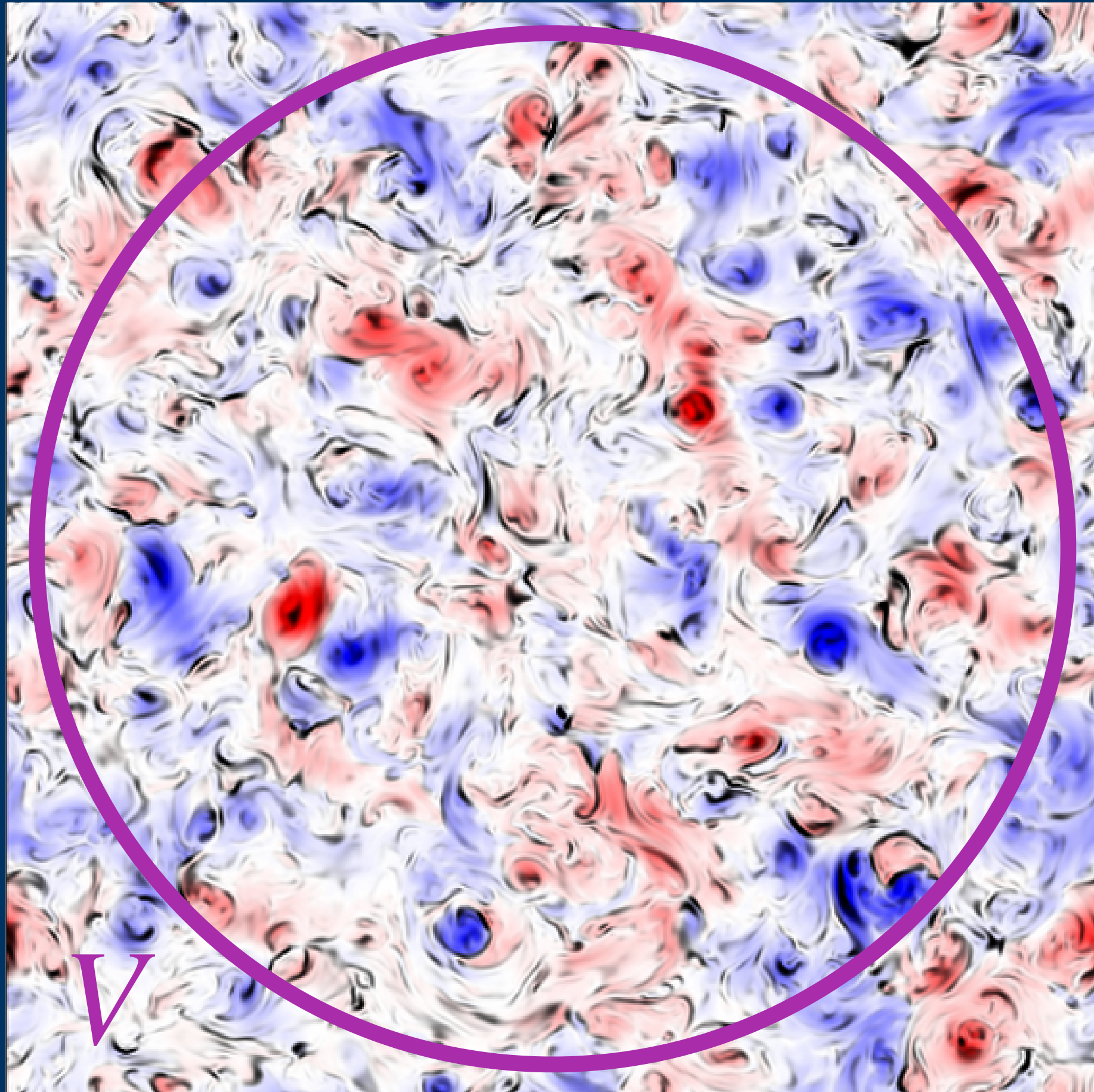


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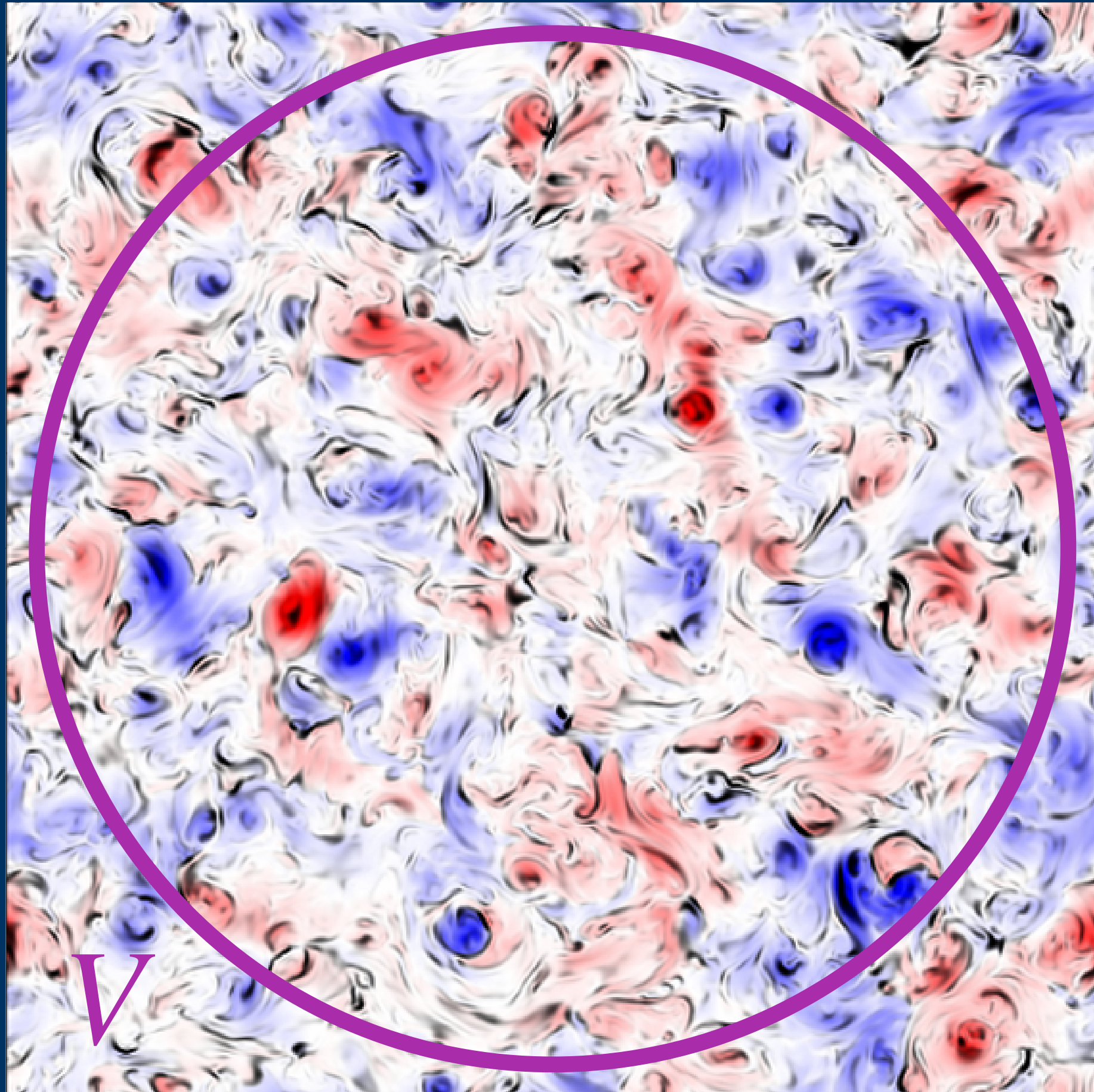
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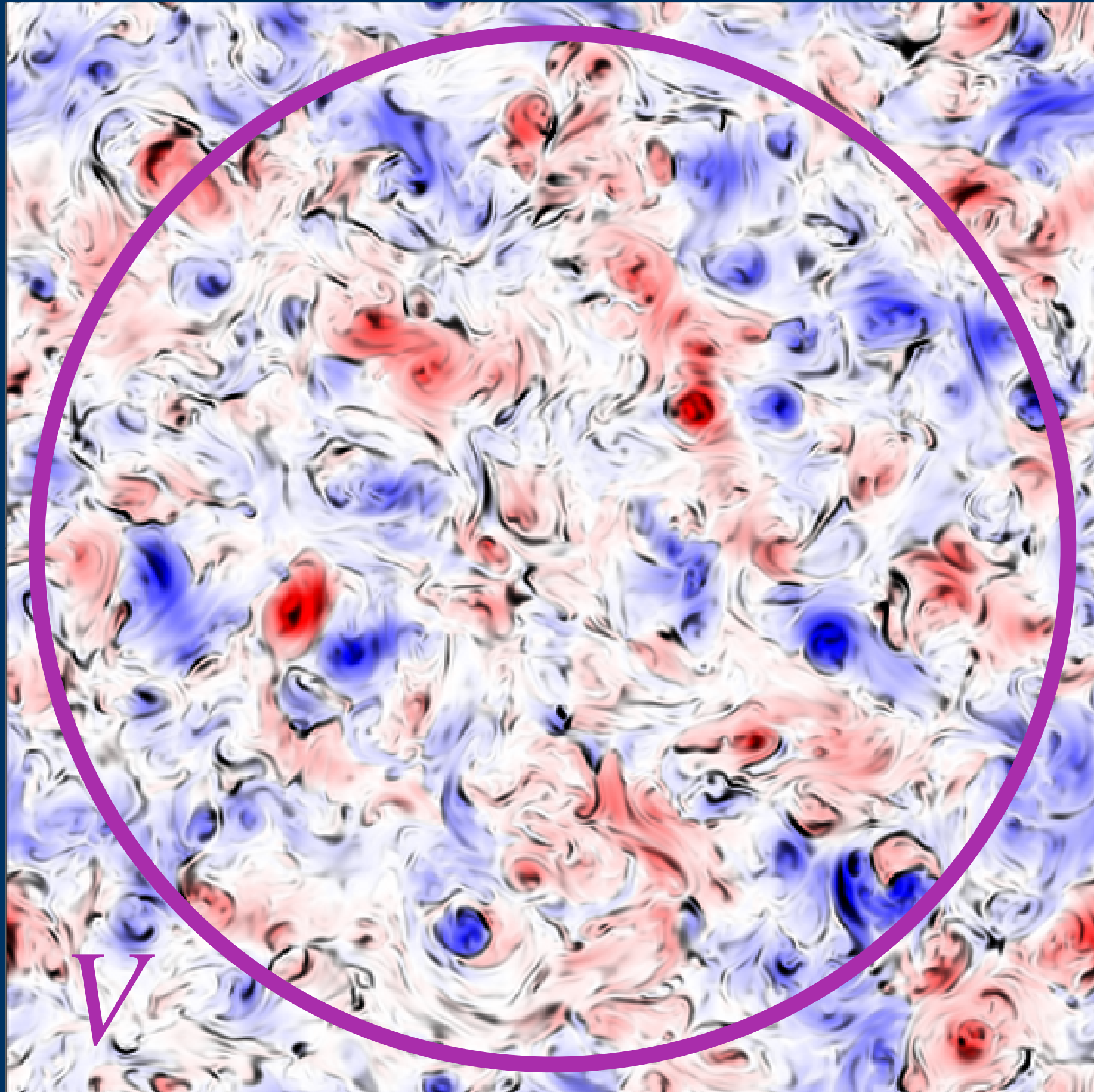


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$$\implies \frac{1}{I_\psi} \frac{dI_\psi}{dt} \propto \lim_{V \rightarrow \infty} V^{-2/3} = 0$$

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# Kinematics: correlation integrals and spectra

The power spectrum of  $\psi$  is

$$\mathcal{E}_\psi(k) = \frac{k^2}{2\pi^2} \int d^3r e^{ik \cdot r} \langle \psi(\mathbf{x}) \psi(\mathbf{x} + \mathbf{r}) \rangle$$

If the correlation function decays sufficiently quickly with  $r$ , then

$$\mathcal{E}_\psi(k \rightarrow 0) = \frac{k^2}{2\pi^2} \int d^3r \langle \psi(\mathbf{x}) \psi(\mathbf{x} + \mathbf{r}) \rangle - \frac{k^4}{2\pi^2} \int d^3r r^2 \langle \psi(\mathbf{x}) \psi(\mathbf{x} + \mathbf{r}) \rangle + O(k^6).$$

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If not, we can show that, for  $0 < a \leq 3$ ,

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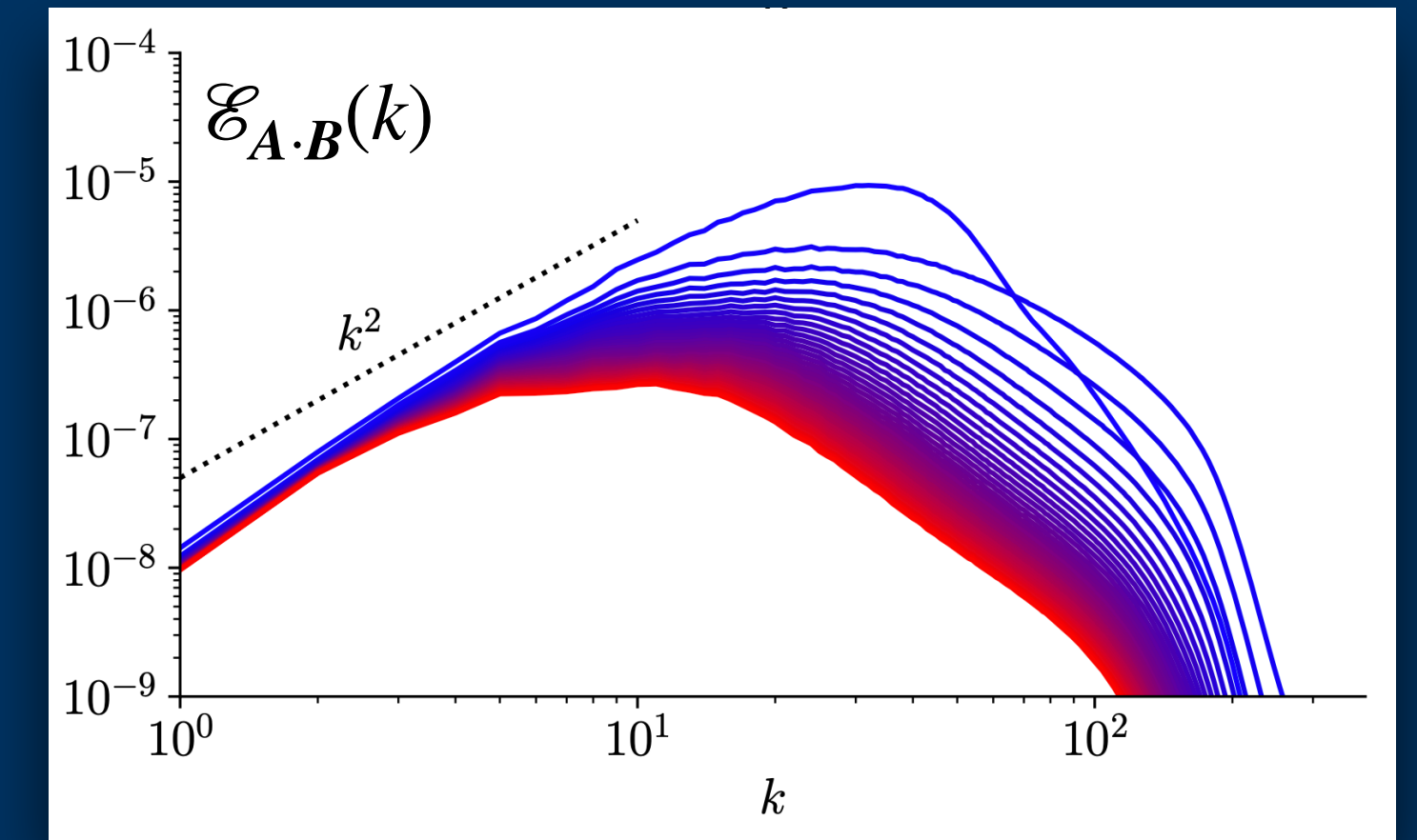
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Simulation of MHD decay

# Correlation integrals and turbulent dynamics

Cascades / inertial-range phenomenology

Inverse cascades and transfers / infra-red phenomenology

Control of decaying turbulence



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$\mathcal{E}_u(k \rightarrow 0)$  grows as  $U$  decays if  $5\alpha > 2\beta + 3$ .

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**Consequence: inverse transfer of energy**

Example: Non-helical MHD

$$I_h = \int d^3r \langle h(\mathbf{x})h(\mathbf{x} + \mathbf{r}) \rangle \sim B^4 L^5 \sim \text{const}$$

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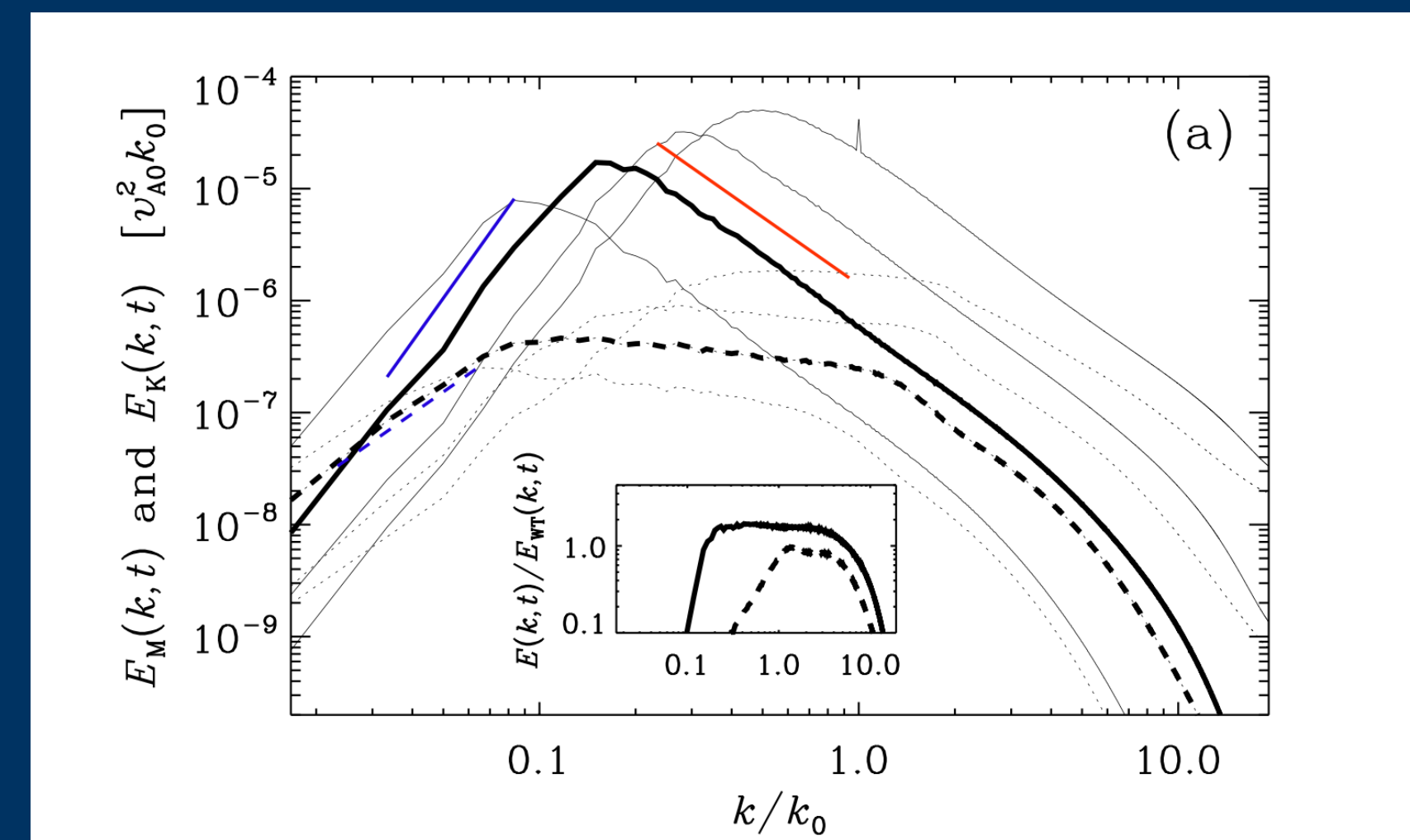
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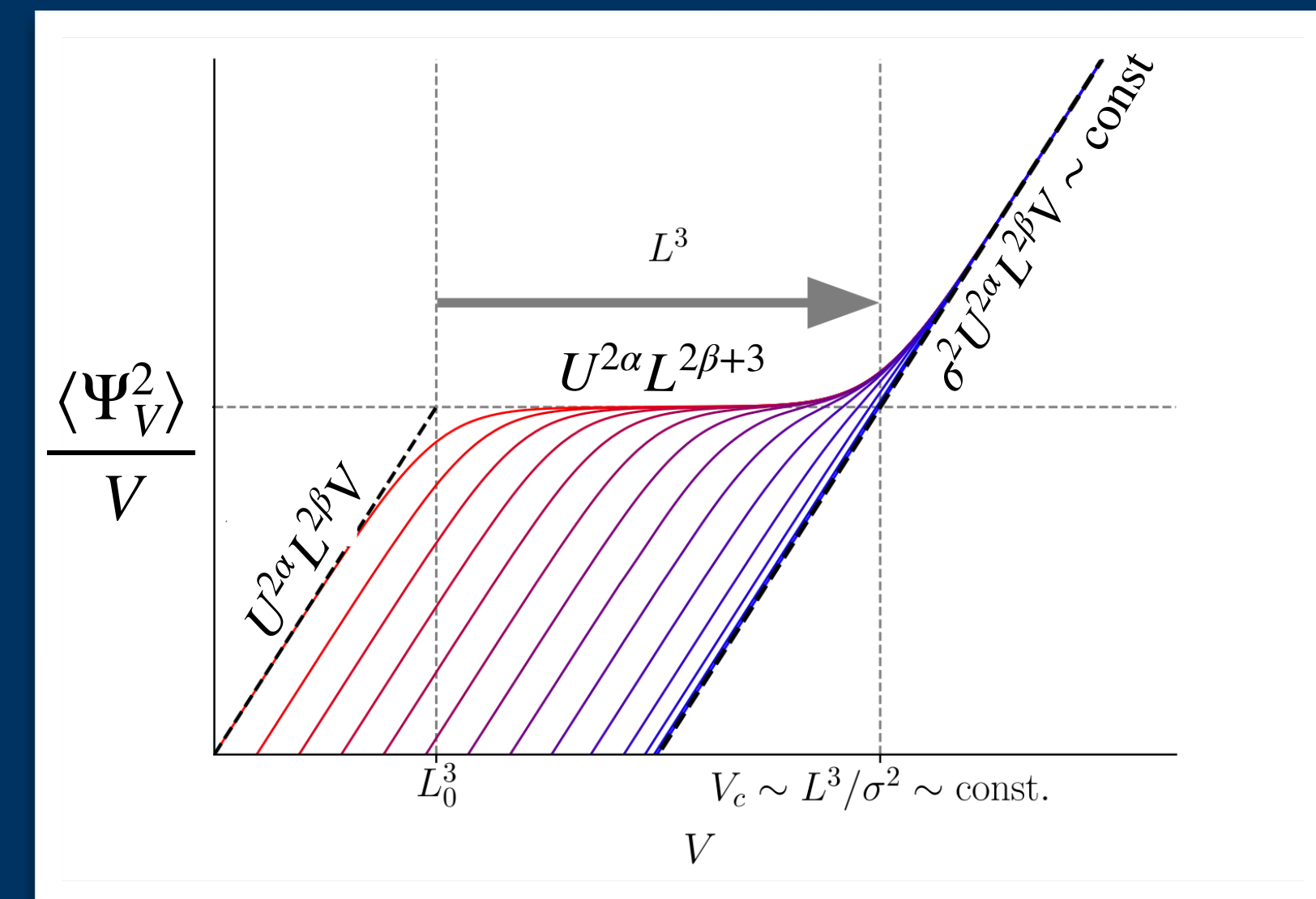
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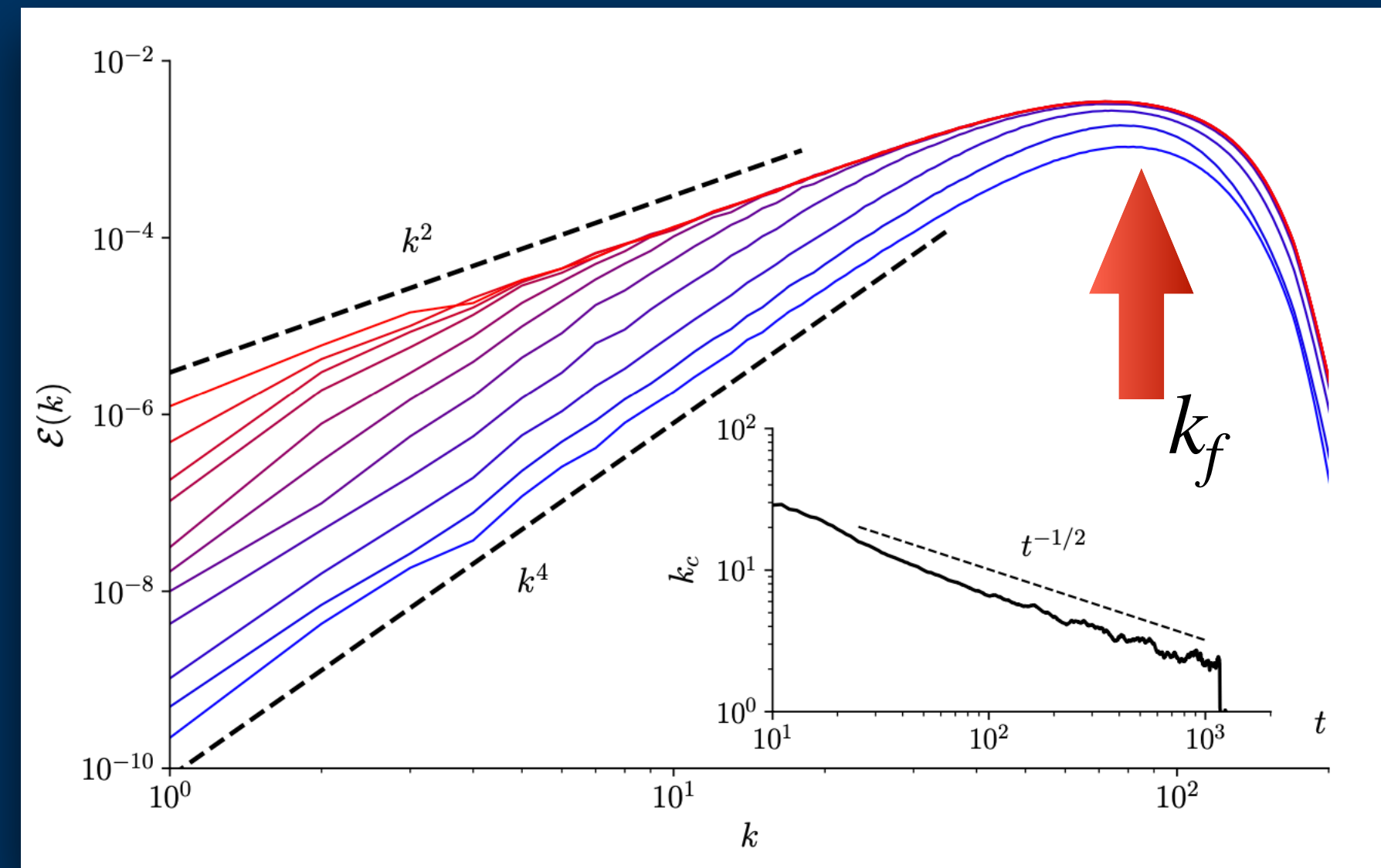
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However, the connection between correlation integrals, large-scale spectra and the scaling of  $\langle \Psi_V^2 \rangle$  with volume can be used to explain certain infra-red phenomena.

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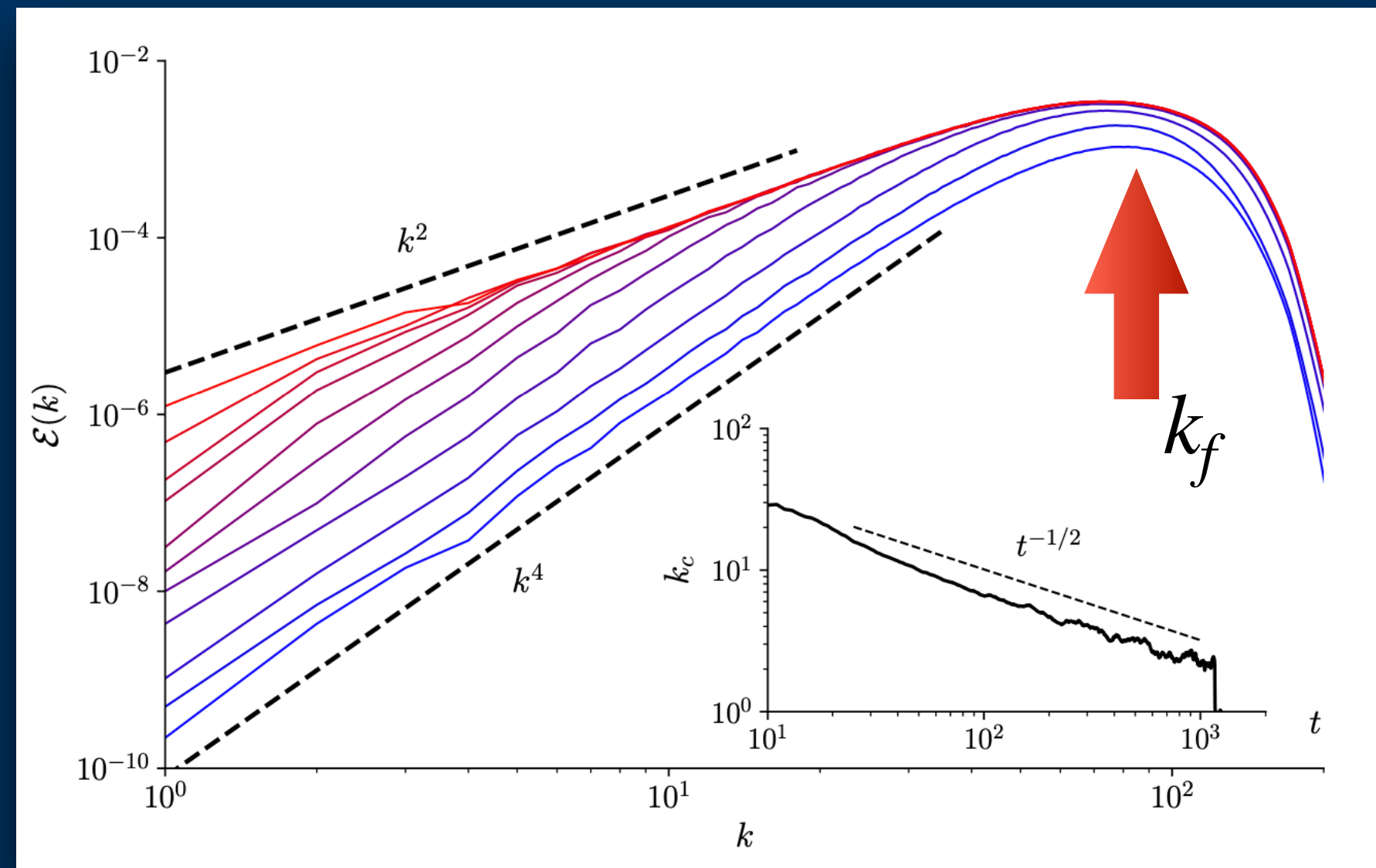
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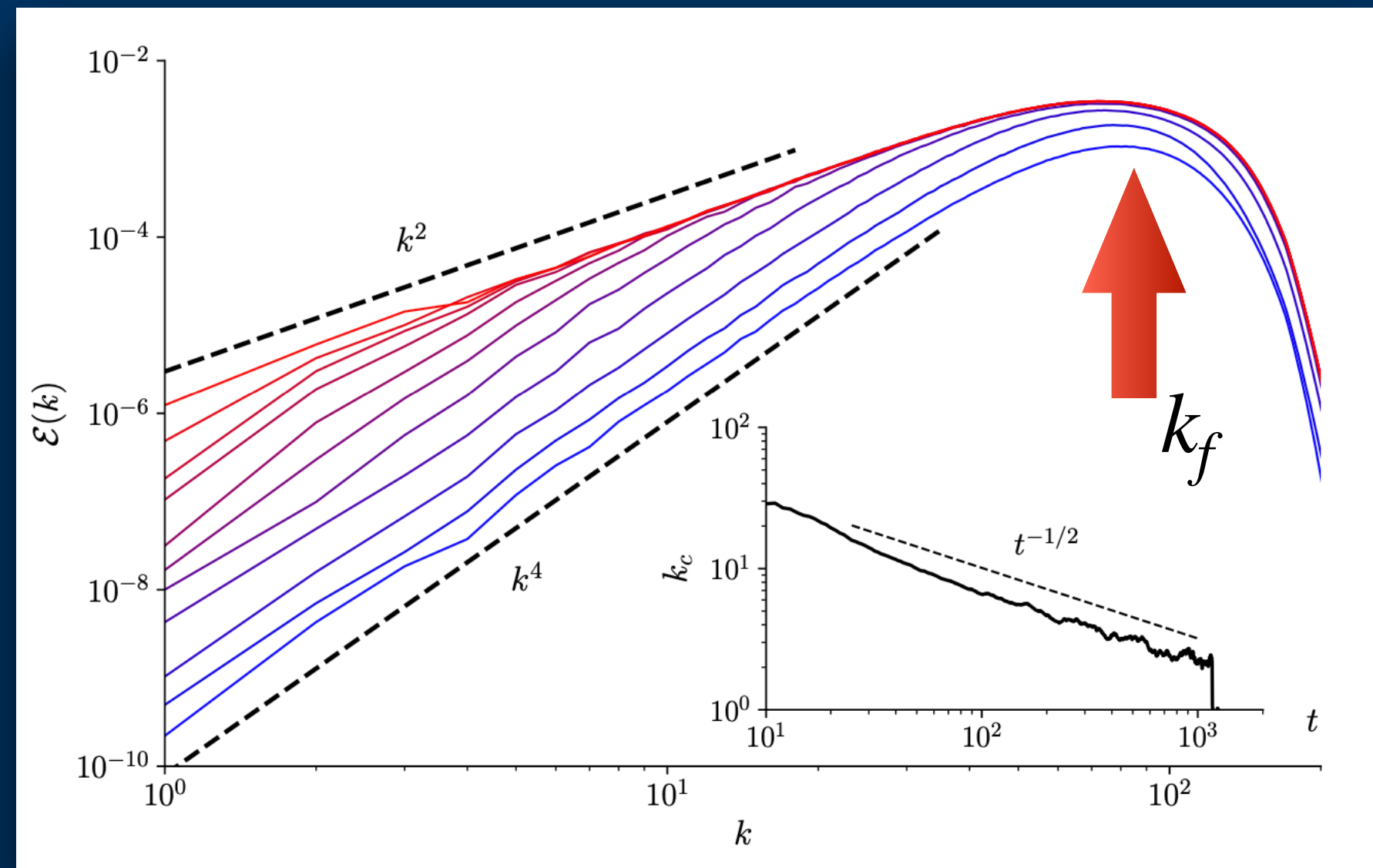


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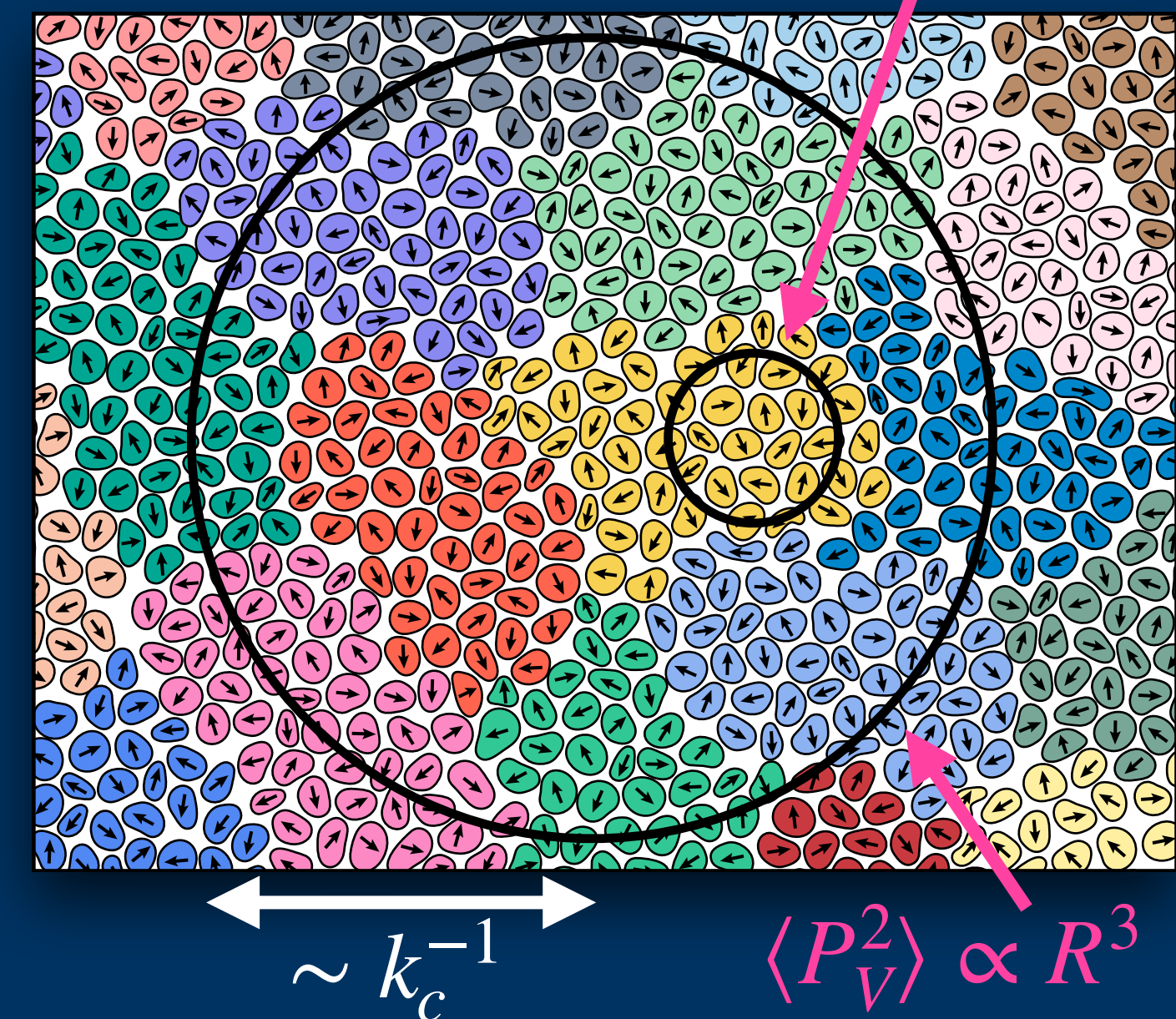


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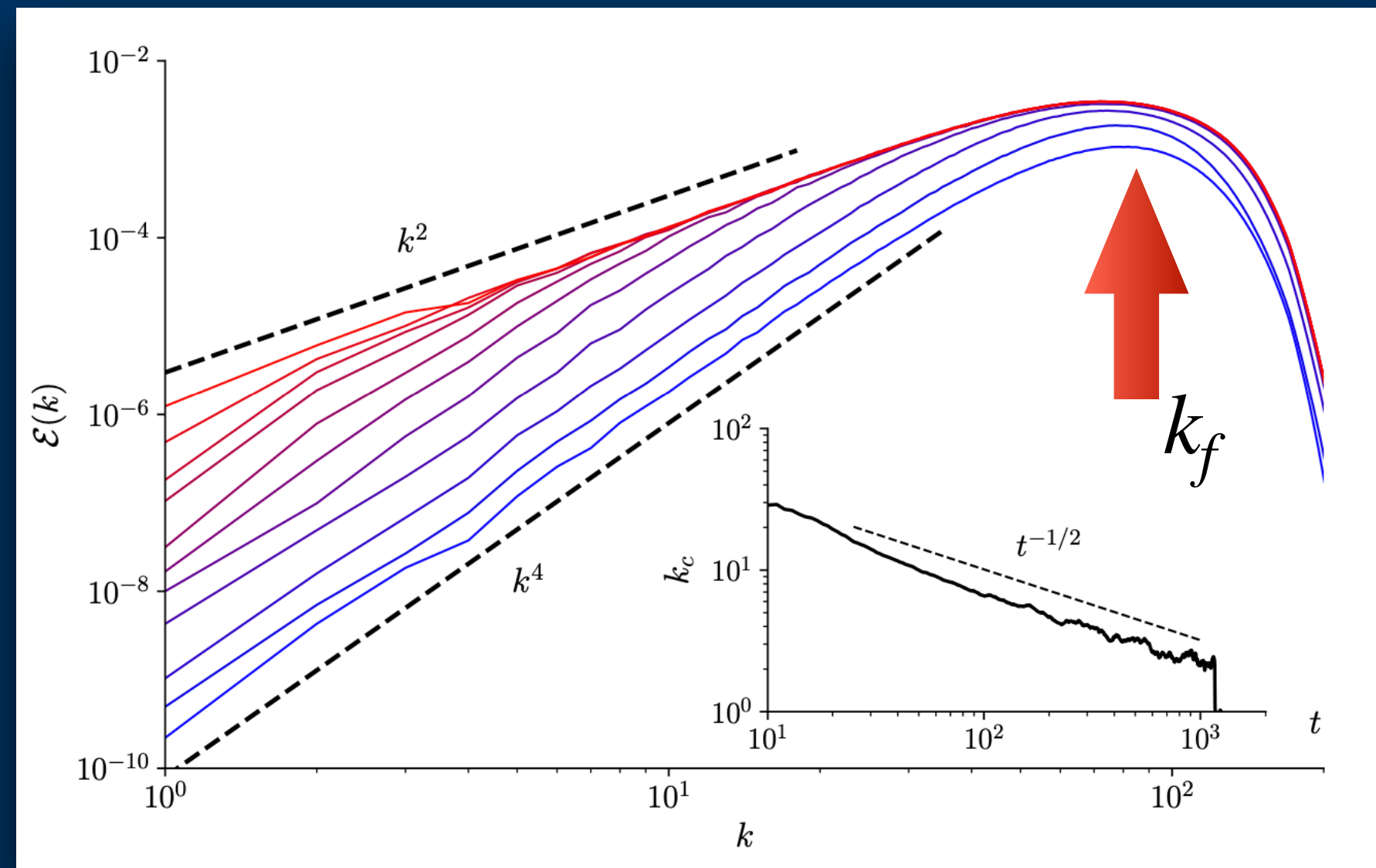
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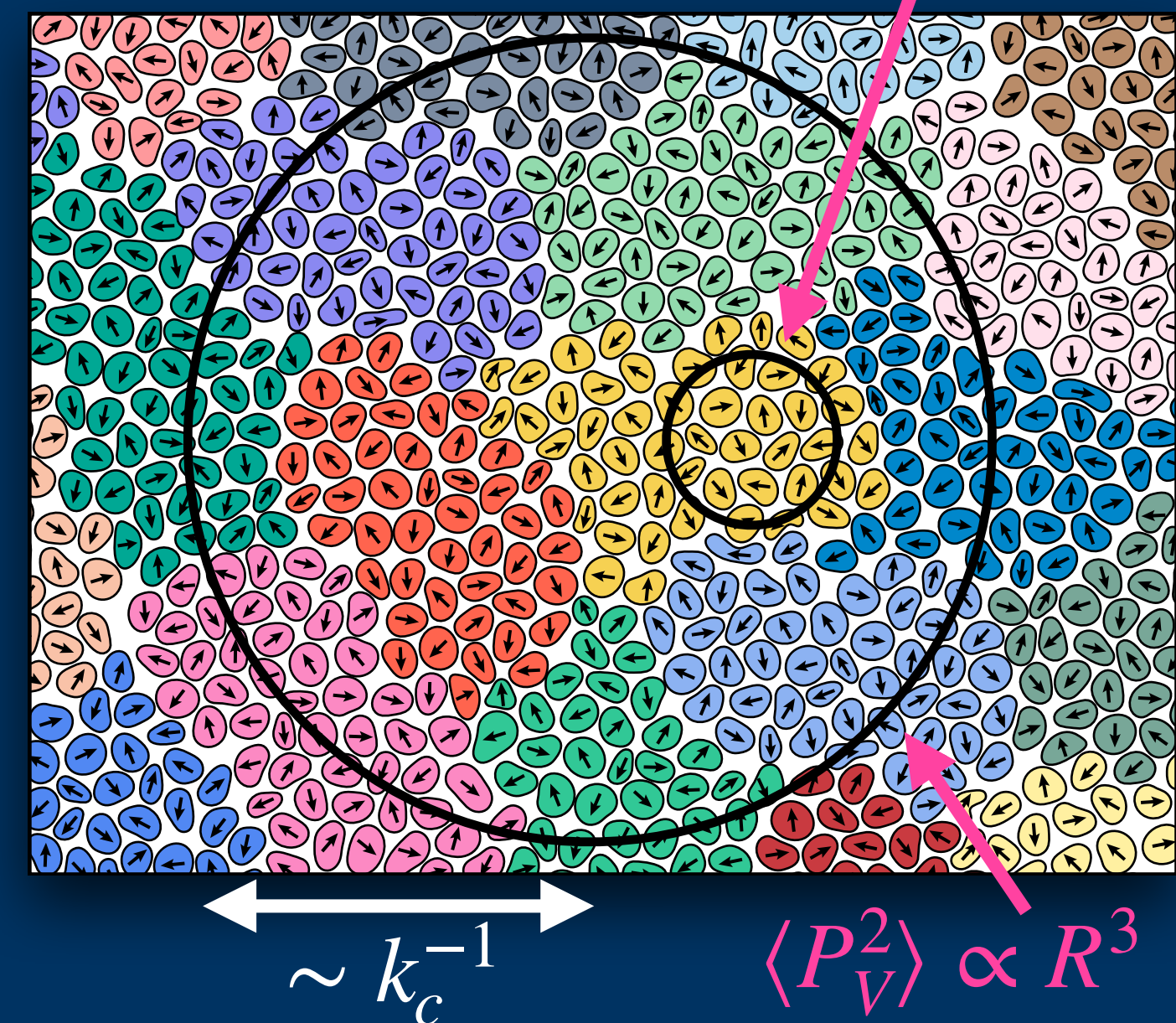
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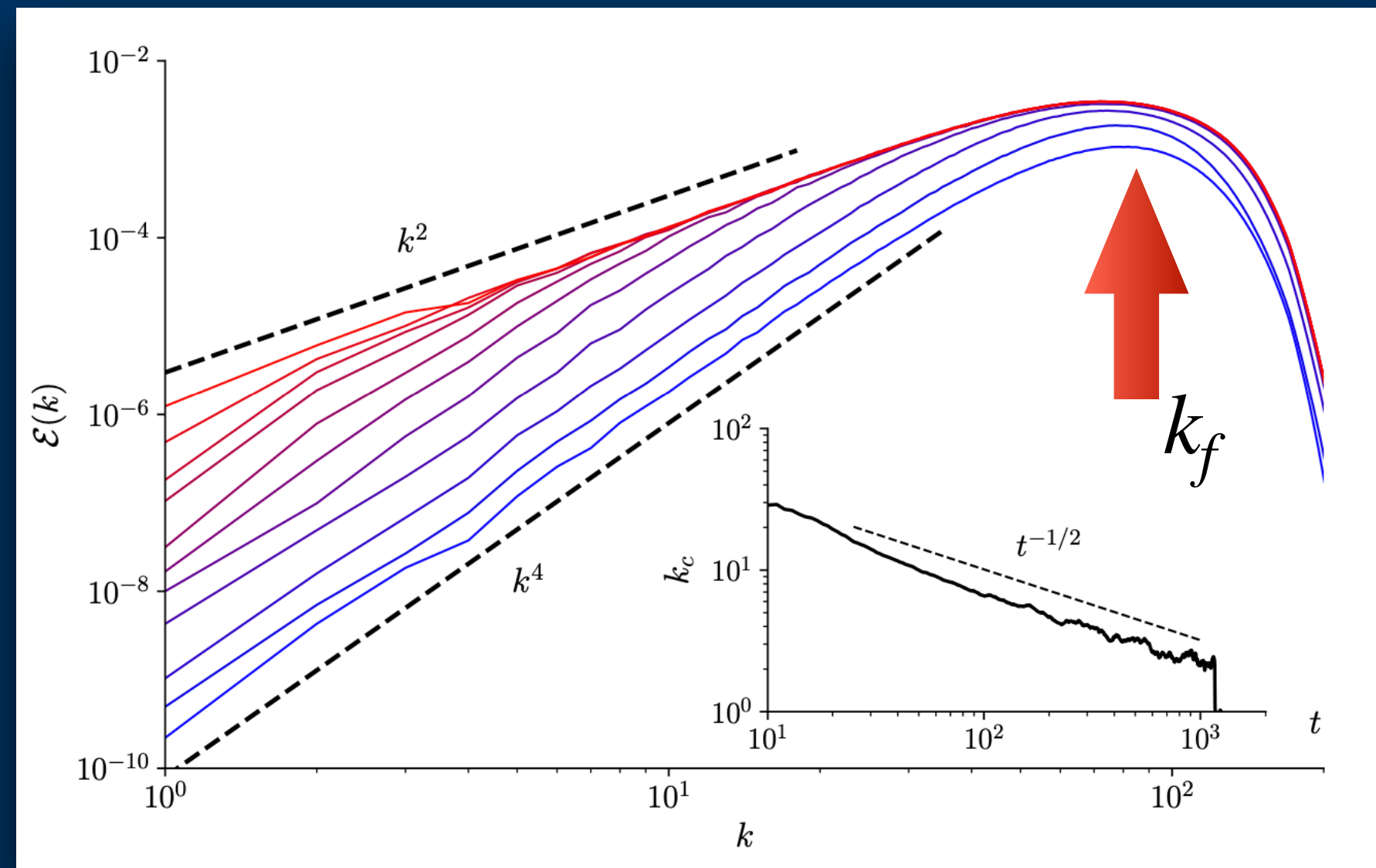
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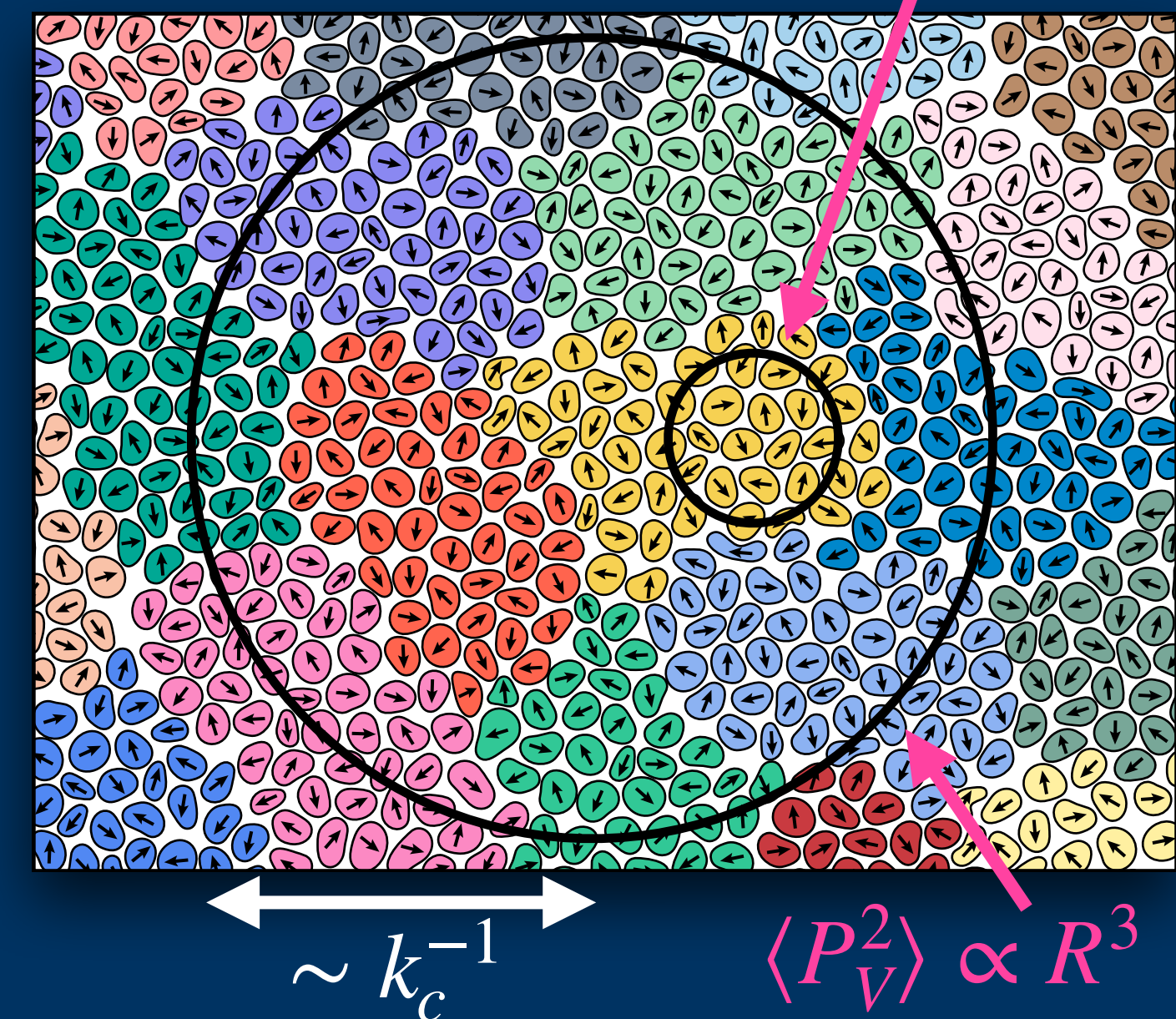
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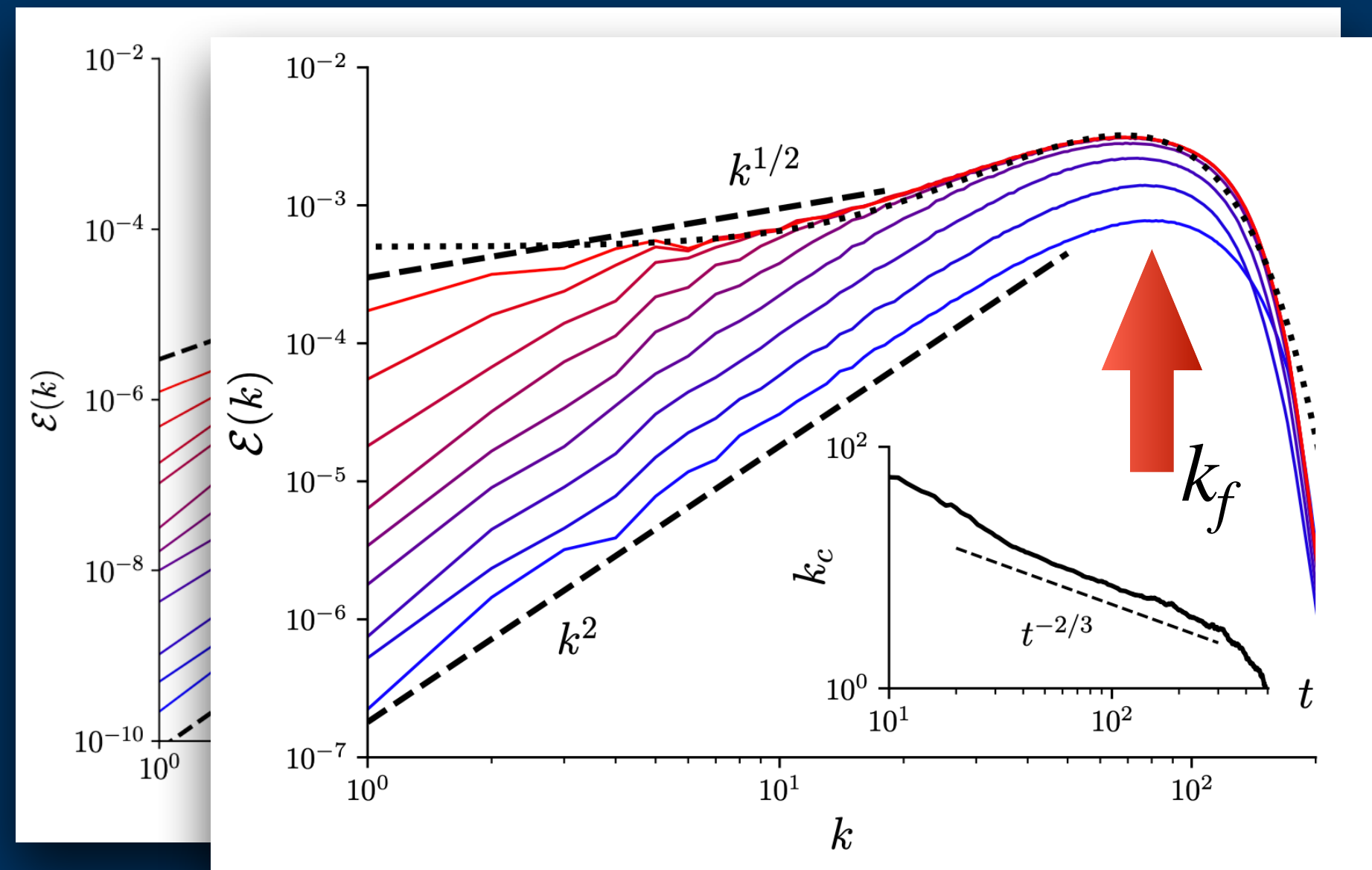
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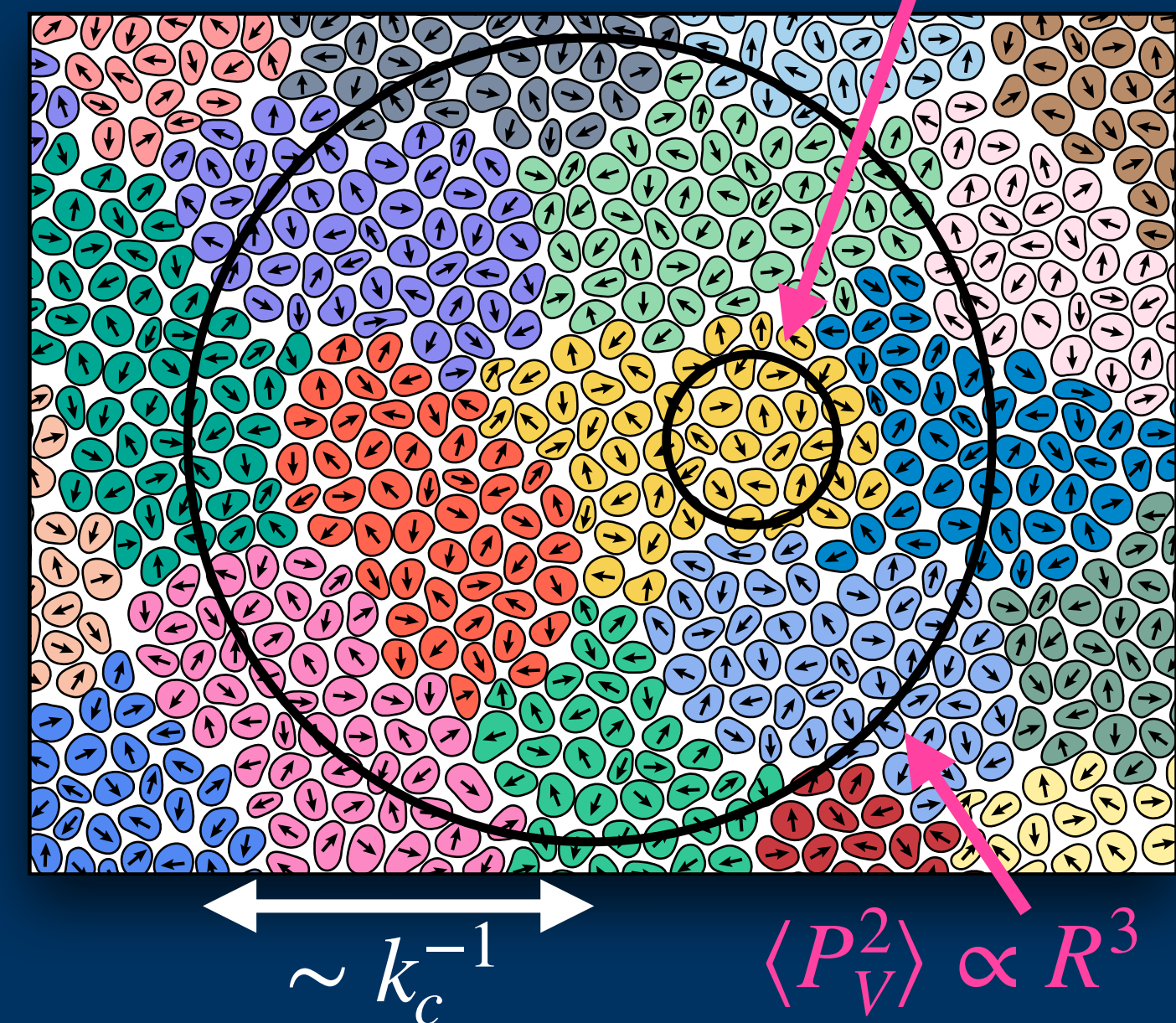
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Conclusion is that, loosely, there must be anomalously large fluctuations in kinetic helicity at small scales in hydro turbulence (cf. Milanese *et. al.* 2022).

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- Levich 2009 presents an correlation-integral-based argument for the development of anomalously strong kinetic-helicity fluctuations at the small scales of hydrodynamic turbulence. Might similar ideas be applicable to cross-helicity? I.e., to scale-dependent alignment of  $\mathbf{u}$  and  $\mathbf{B}$ ?