THE HOLOGRAPHIC SWAMPLAND

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(based on JC, Quevedo 1811.06276, JC, Revello 2006.01021 JC, Ning, Revello in progress)

MODULI: WHAT?

$$\mathcal{L} = \mathcal{L}_{GR} + \mathcal{L}_{SM} + \mathcal{L}_{BSM}$$

- What is in \mathcal{L}_{BSM} ?
- Simple concept: massive scalar Φ with gravitationally suppressed couplings to ordinary matter such as $\frac{\Phi}{M_P}F_{\mu\nu}F^{\mu\nu}$
- Such moduli are well motivated from e.g. string theory and extra-dimensional theories

MODULI: WHY?

- String theory is a theory of dynamical extra dimensions
- In 4d theory, geometry of extra dimensions (size and shape)



parametrised by moduli - such as Kahler and complex structure moduli.

- Unstabilised, these lead to fifth forces, varying couplings or (fatal) decompactification.
- Essential to develop moduli potentials that fix this geometry
- Stabilisation also provides a minimum in which to compute couplings

MODULI: WHY?

• In an expanding universe

$$\rho_{matter} \sim \frac{1}{a^3} \qquad \rho_{radiation} \sim \frac{1}{a^4}$$

- As matter dominates over radiation, reheating is dominated by the *last* fields to decay *not* the first
- The weaker the coupling, the longer the lifetime.... $\tau_{\Phi} \sim \frac{8\pi M_P^2}{g^2 m_{\Phi}^3} \sim \left(\frac{10 \text{ TeV}}{m_{\Phi}}\right)^3 10^{-3} s$ $\tau \sim \frac{8\pi}{g^2 m}$ • Moduli potentials are everyone's business

Much work on developing moduli potentials (LVS, KKLT) and studying their dynamics with regards to Supersymmetry breaking
 Cosmology - late time de Sitter
 Cosmology - inflation
 Particle physics







LARGE VOLUME SCENARIO

Balasubramanian, Berglund, JC, Quevedo

Perturbative corrections to K and non-perturbative corrections to W

$$W = \int G_3 \wedge \Omega + \sum_i A_i e^{-2\pi a_i T_i}$$

$$K = -2\ln\left(\mathcal{V} + \xi'\right) + \ln\left(\int \Omega \wedge \overline{\Omega}\right) - \ln(S + \overline{S})$$

Resulting scalar potential has minimum at exponentially large values of the volume

$$V = \frac{A\sqrt{\tau_s}e^{-2a_s\tau_s}}{\mathcal{V}} - \frac{B\tau_s e^{-a_s\tau_s}}{\mathcal{V}^2} + \frac{C}{\mathcal{V}^3}$$



Figure 1: $\ln(V)$ for $P_{[1,1,1,6,9]}^4$ in the large volume limit, as a function of the divisors τ_4 and τ_5 . The void channel corresponds to the region where V becomes negative and $\ln(V)$ undefined. As $V \to 0$ at infinite volume, this immediately

WHY LVS?

- In LVS, volume is exponentially large can easily be $\mathcal{V} \sim 10^{50} (2\pi \sqrt{\alpha'})^6$
- This generates interesting hierarchies and ensures superb parametric decoupling of heavy modes (KK modes, heavy moduli)
- Decoupling also has a clear geometric origin large volume $\langle \mathcal{V} \rangle \sim e^{\xi/g_s}$
- $\mathcal{V} \! \rightarrow \! \infty$ limit of LVS also leads to a unique effective theory

LARGE VOLUME SCENARIO

• LVS effective theory for volume modulus Φ and axion a

$$V_{potential} = V_0 e^{-\lambda \Phi/M_p} \left(-\left(\frac{\Phi}{M_p}\right)^{3/2} + A \right) \qquad (\lambda = \sqrt{27/2}$$
$$\mathcal{L}_{kinetic} = \frac{1}{2} \partial_\mu \Phi \partial_\mu \Phi + \frac{3}{4} e^{-\sqrt{\frac{8}{3}}\Phi} \partial_\mu a \partial^\mu a$$

• Other terms are subleading in infinite volume limit by $O\left(\frac{1}{\ln V}\right)$

LARGE VOLUME SCENARIO

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Solve for minimum and expand about it to determine masses and couplings

THE SWAMPLAND (Vafa et al)

- Which low-energy Lagrangians are forbidden by quantum gravity?
- Do de Sitter vacua exist in string theory?
- Can large (trans-Planckian) field inflation occur?



String theory

- IO-dimensional supergravity with alpha' corrections
- 4-dimensional supergravity of moduli and matter
- Integrate out heavy modes to get potential for lightest moduli J EFT
- Find vacuum as minimum of effective potential, construct de Sitter space in string theory.....



Find vacuum as minimum of effective potential

- String theory
- 10-dimensional supergravity with alpha' corrections
- 4-dimensional supergravity of moduli and matter
- Integrate out heavy modes to get potential for lightest moduli
 EFT
- Find vacuum as minimum of effective potential

HOLOGRAPHY

• CFT dimensions of dual operators: $\Delta(\Delta - 3) = m_{\Phi}^{2} R_{AdS}^{2}$ • In infinite volume limit can classify modes as heavy $m_{\Phi}^{2} \gg R_{AdS}^{-2}, \Delta \to \infty$ as $\mathcal{V} \to \infty$ light $m_{\Phi}^{2} \ll R_{AdS}^{-2}, \Delta \to 3$ as $\mathcal{V} \to \infty$ interesting $m_{\Phi}^{2} \sim R_{AdS}^{-2}, \Delta \to \mathcal{O}(1-10)$ as $\mathcal{V} \to \infty$

LVS MASS SPECTRUM

- In LVS we have
- **Heavy:** KK modes, complex structure moduli, all Kahler moduli except overall volume
- Light: Graviton, overall volume axion
- Interesting: overall volume modulus

LVS HOLOGRAPHY

Mode	Spin	Parity	Conformal dimension
$T_{\mu u}$	2	+	3
a	0	-	3
Φ	0	+	$8.038 = rac{3}{2} \left(1 + \sqrt{19} ight)$

Table 1. The low-lying single-trace operator dimensions for CFT duals of the Large Volume Scenario in the limit $\mathcal{V} \to \infty$.

In minimal LVS, AdS effective theory has small number of fields which correspond to specific predictions for dual conformal dimensions

No Landscape! (not true of KKLT)

LVS HOLOGRAPHY

- LVS is attractive as it offers a well-motivated Generalised Free Field Theory
- Large volume limit $\mathcal{V} \rightarrow \infty$ gives a unique theory
- Two scalars with fixed and radiatively stable anomalous dimensions
- All AdS interactions are also fixed and radiatively stable

LVS AND THE SWAMPLAND

• Now consider this small modification:

$$V_{potential} = V_0 e^{-\lambda \Phi/M_p} \left(-\left(\frac{\Phi}{M_p}\right)^{3/2} + A \right) \qquad \left(\lambda = \sqrt{27/2}\right)$$
$$\mathcal{L}_{kinetic} = \frac{1}{2} \partial_\mu \Phi \partial_\mu \Phi + \frac{3}{4} e^{+\sqrt{\frac{8}{3}}\Phi/M_p}} \partial_\mu a \partial^\mu a$$

 This coupling is equivalent to axion decay constants that *diverge* in the decompactification limit - must be in the swampland!

$$f_a / M_P \to \infty$$
 as $\mathcal{V} \to \infty$

HOLOGRAPHIC SWAMPLAND

n-point self interactions of volume modulus

$$\mathcal{L}_{n-pt} = (-1)^{n-1} \lambda^n (n-1) \left(-3 \frac{M_P^2}{R_{AdS}^2} \right) \frac{1}{n!} \left(\frac{\delta \Phi}{M_P} \right)^n \left(1 + \mathcal{O} \left(\frac{1}{\ln \mathcal{V}} \right) \right) \left(\lambda = \sqrt{27/2} \right)^n$$

Mixed interactions of volume modulus and axion

$$\mathcal{L}_{\Phi^{n}aa} = \left(-\sqrt{\frac{8}{3}}\right)^{n} \frac{1}{2n!} \left(\frac{\delta \Phi}{M_{P}}\right)^{n} \partial_{\mu} a \partial^{\mu} a$$

 The higher-point interaction define 3- and higher point-correlators within a dual CFT

HOLOGRAPHIC SWAMPLAND

 $\left(\lambda = \sqrt{27/2}\right)$

n-point self interactions of volume modulus

$$\mathcal{L}_{n-pt} = (-1)^{n-1} \lambda^n (n-1) \left(-3 \frac{M_P^2}{R_{AdS}^2} \right) \frac{1}{n!} \left(\frac{\delta \Phi}{M_P} \right)^n \left(1 + \mathcal{O} \left(\frac{1}{\ln \mathcal{V}} \right) \right)$$

Now modify interactions of volume modulus and axion

$$\mathcal{L}_{\Phi^{n}aa} = \left(+\sqrt{\frac{8}{3}} \right)^{n} \frac{1}{2n!} \left(\frac{\delta \Phi}{M_{P}} \right)^{n} \partial_{\mu} a \partial^{\mu} a$$

 This defines a perturbation to the Generalised Free Field CFT with axion decay constants that *diverge* in the decompactification limit must be in the swampland!

$$f_a / M_P \to \infty$$
 as $\mathcal{V} \to \infty$

HOLOGRAPHIC SWAMPLAND

• The problem:

I. Generalised Free Field + (some corrections)
- consistent theory
2. Generalised Free Field + (other corrections)
- swampland!

Where does the difference lie? Can one correlate any properties of the CFT with this change from the consistent theory to the swampland theory?