

Turbulence in magnetized pair plasmas

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Magnetized pair plasmas

Fluid equations for a low beta, non-relativistic pair plasma



Start from two-fluid equations for electrons and positrons:

$$n^{\pm}m \left(\frac{\partial \mathbf{v}^{\pm}}{\partial t} + \mathbf{v}^{\pm} \cdot \nabla \mathbf{v}^{\pm} - \mu^{\pm} \nabla^2 \mathbf{v}^{\pm} \right) = \pm n^{\pm}e \left(\mathbf{E} + \frac{\mathbf{v}^{\pm} \times \mathbf{B}}{c} \right) - \nabla p^{\pm} \mp \mathcal{R}$$

Sum and subtract:

$$nm \left(\frac{\partial \mathbf{v}}{\partial t} + \mathbf{v} \cdot \nabla \mathbf{v} + \frac{1}{n^2 e^2} \mathbf{j} \cdot \nabla \mathbf{j} - \mu \nabla^2 \mathbf{v} \right) = \frac{\mathbf{j} \times \mathbf{B}}{c} - \nabla p$$

\mathbf{v} is center-of-mass velocity

$$\frac{m}{ne^2} \left(\frac{\partial \mathbf{j}}{\partial t} + \mathbf{v} \cdot \nabla \mathbf{j} + \mathbf{j} \cdot \nabla \mathbf{v} - \mu \nabla^2 \mathbf{j} \right) = \mathbf{E} + \frac{\mathbf{v} \times \mathbf{B}}{c} - \eta \mathbf{j}.$$

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Impose reduced-MHD-like ordering: $\mathbf{B} = B_0 \hat{z} + \mathbf{B}_\perp$, with $B_\perp/B_0 \sim \epsilon$

Introduce potentials: $\mathbf{v}_\perp = \hat{z} \times \nabla_\perp \phi$; $\frac{\mathbf{B}_\perp}{\sqrt{4\pi\rho}} = \hat{z} \times \nabla_\perp \psi$,

$$\frac{\partial}{\partial t} \nabla_\perp^2 \phi + \{\phi, \nabla_\perp^2 \phi\} = \{\psi, \nabla_\perp^2 \psi\} + V_A \frac{\partial}{\partial z} \nabla_\perp^2 \psi + \mu \nabla_\perp^2 \phi, \quad \text{momentum equation}$$

$$\frac{\partial}{\partial t} (1 - d_e^2 \nabla_\perp^2) \psi + \{\phi, (1 - d_e^2 \nabla_\perp^2) \psi\} = V_A \frac{\partial \phi}{\partial z} + \eta \nabla_\perp^2 \psi - \mu d_e^2 \nabla_\perp^4 \psi \quad \text{Ohm's law}$$

These reduce to the familiar RMHD equations when electron inertia is neglected (unsurprisingly).

These equations have two *exact* invariants at *all* scales:

$$\mathcal{E} = \frac{1}{2} \int dV \left\{ (\nabla_{\perp} \psi)^2 + d_e^2 (\nabla_{\perp}^2 \psi)^2 + (\nabla_{\perp} \phi)^2 \right\} \quad \text{energy}$$

$$\mathcal{H}^C = \int dV \left\{ \nabla_{\perp}^2 \phi (1 - d_e^2 \nabla_{\perp}^2) \psi \right\} \quad \text{(generalized) cross helicity}$$

Only one wave is supported:

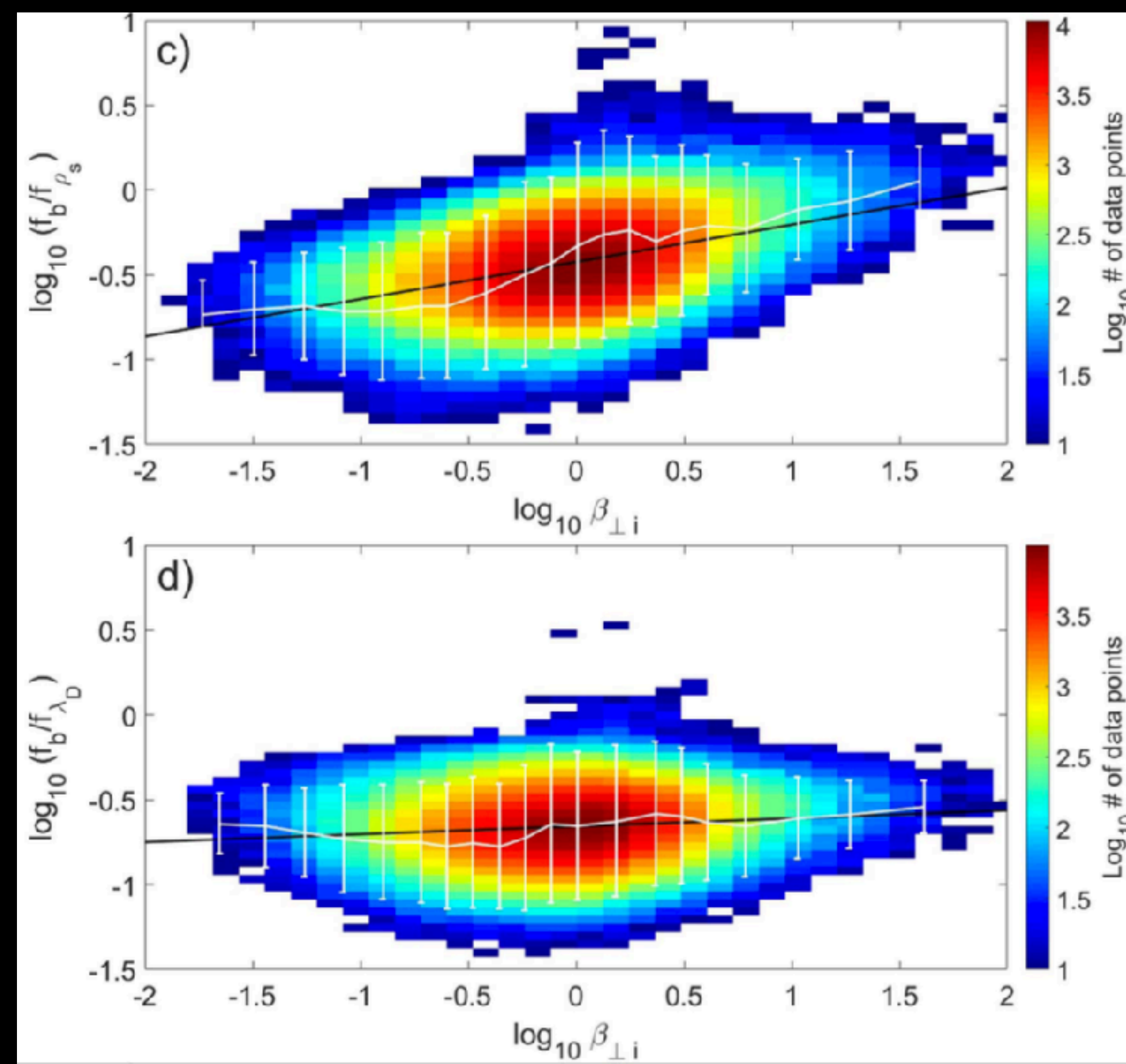
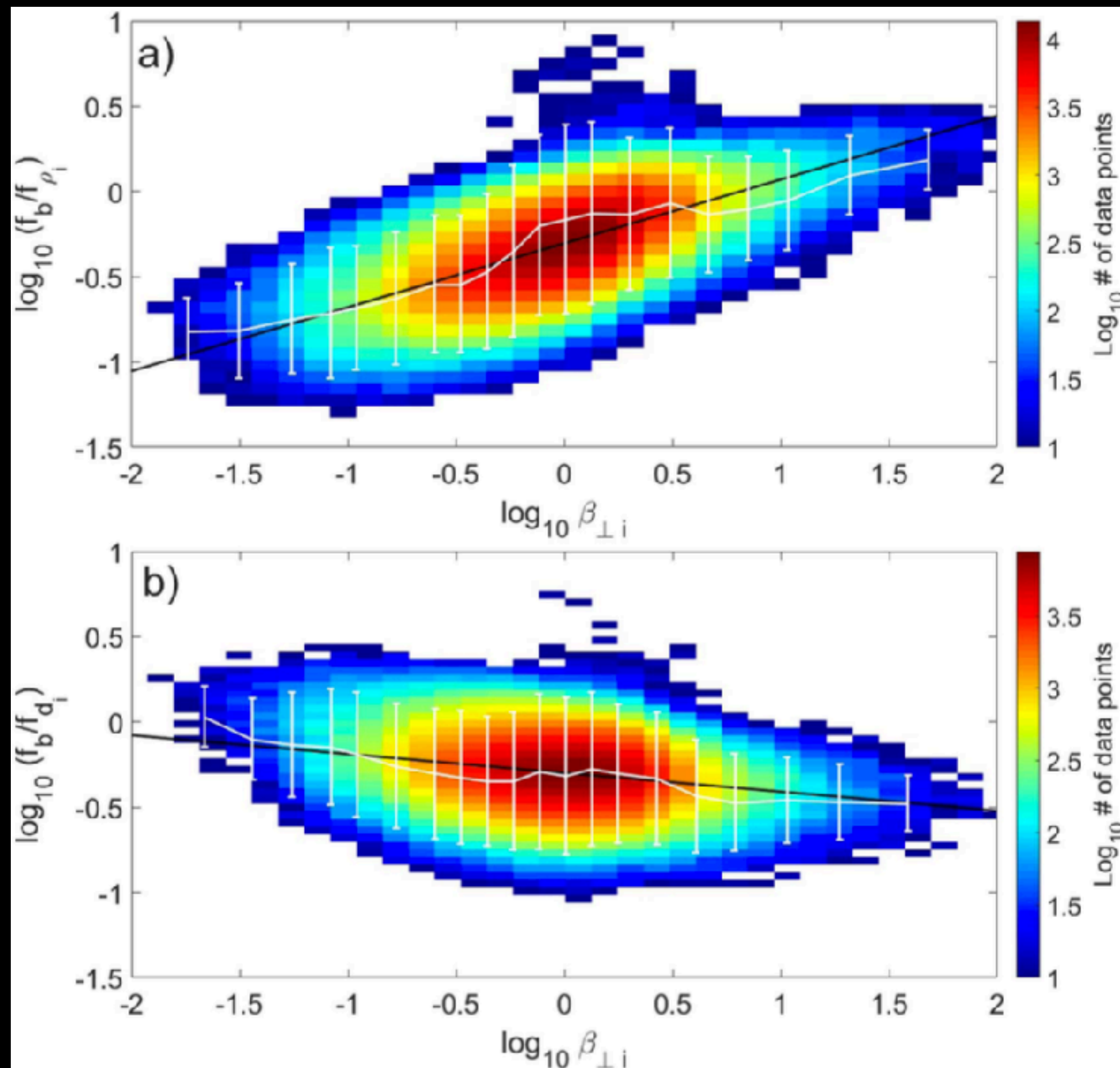
$$\omega_l = \pm \frac{k_z V_A}{\sqrt{1 + k_{\perp}^2 d_e^2}} \quad \text{Alfvén wave modified at kinetic scales by electron inertia}$$

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Turbulence at MHD scales

At **MHD scales** ($k d_e \ll 1$), the equations are the same as for ion-electron plasmas. So turbulence will be the same:

1. expect $k^{-3/2}$ up until the reconnection scale,
2. followed by a transition to a k^{-3} (or $-8/3$) due to reconnection (?)



Vech et al., ApJ 2018

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Turbulence at kinetic scales



Let us work out the spectrum expected from a Kolmogorov-like energy cascade at **kinetic scales** ($kd_e \gg 1$).

$$\mathcal{E} = \frac{1}{2} \int dV \left\{ (\nabla_{\perp} \psi)^2 + d_e^2 (\nabla_{\perp}^2 \psi)^2 + (\nabla_{\perp} \phi)^2 \right\}$$

$$\Rightarrow \frac{1}{2} \int dV \left\{ d_e^2 (\nabla_{\perp}^2 \psi)^2 + (\nabla_{\perp} \phi)^2 \right\}$$

Expect equipartition between these two terms (parallel and perpendicular kinetic energies) at these scales.

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Turbulence at kinetic scales (cont'd)

$$k_{\perp}^2 \phi_{\lambda}^2 / \tau_{\lambda} \sim \varepsilon$$

$$\tau_{\lambda} = 1/\omega_{nl} \sim 1/(k_{\perp}^2 \phi_{\lambda})$$

This results in:

$$\phi_{\lambda} \sim \varepsilon^{1/3} k_{\perp}^{-4/3} \Rightarrow E_{\phi}(k_{\perp}) dk_{\perp} \sim \varepsilon^{2/3} k_{\perp}^{-11/3} dk_{\perp}$$

and same scaling for magnetic energy, by equipartition

Finally, declare that the fluctuations are *critically balanced* at these scales: $\omega_l \sim \omega_{nl}$

$$k_{\parallel} \sim \varepsilon^{1/3} d_e V_A^{-1} k_{\perp}^{5/3}$$

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Dealing with the cross helicity



Recall cross-helicity:

$$\mathcal{H}^C = \int dV \{ \nabla_{\perp}^2 \phi (1 - d_e^2 \nabla_{\perp}^2) \psi \}$$

It is *not positive definite*.

Estimate the flux of cross helicity as

$$(k_{\perp}^2 \phi_{\lambda}) (d_e^2 k_{\perp}^2 \psi_{\lambda}) R_{\lambda} / \tau_{\lambda} \sim \varepsilon^c$$

where R is a dimensionless cancelation factor at scale λ . From the energy invariant:

$$k_{\perp}^2 d_e^2 \psi_{\lambda}^2 \sim \phi_{\lambda}^2 \quad \text{so the cross helicity flux becomes}$$

$$k_{\perp} d_e (k_{\perp}^2 \phi_{\lambda}^2) R_{\lambda} / \tau_{\lambda} \sim \varepsilon^c$$

But energy flux (constant) is

$$k_{\perp}^2 \phi_{\lambda}^2 / \tau_{\lambda} \sim \varepsilon$$

Thus:

$$R_{\lambda} \propto 1 / (k_{\perp} d_e)$$

Presented a tentative theory of turbulence in low beta, strongly magnetized, non-relativistic pair plasmas:

- same as MHD at fluid scales
- Energy cascade implies $k_{\perp}^{-11/3}$ spectrum prediction at kinetic scales.
- Critical balance then gives $k_{\parallel} \sim k_{\perp}^{5/3}$
- Notice that reconnection is not allowed at $k_{\perp} d_e \gg 1$ (flux is no longer frozen) so no need to worry about that.

Loureiro & Boldyrev, arXiv:1805.09224

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