

# **Dissipation/plasma in compact objects**

Andrei Beloborodov  
Columbia University

# ”Compactness” in astrophysics

Gravitational:  $\ell_g = \frac{R_S}{R} \sim \frac{\Phi}{c^2}$

Radiative: cooling time vs. light crossing time

$$\frac{t_{\text{cool}}}{s/c} \sim \frac{1}{\gamma_e \ell}$$

$$\ell \equiv \frac{U}{m_e c^2} \sigma_T s \sim \frac{m_p}{m_e} \frac{L}{L_{\text{Edd}}} \frac{R_S}{R} \gg 1$$

(accreting neutron stars, black holes,  
gamma-ray bursts)

## **Dissipative plasma processes:**

reconnection, shocks, turbulence damping, discharge

In compact objects these processes occur in dense radiation

⇒

- 1) dominant radiative losses
- 2) “bulk Comptonization” of photons  
(turns out more important than thermal Comptonization)
- 3) extension of radiation spectrum above  $m_e c^2$
- 4) copious  $e^+e^-$  pair creation
- 5) non-linear self-regulated state  $(n_{\pm}/n_{\text{ion}}, n_{\text{ph}}/n_{\text{ion}})$

# I. Shocks

Plasma shocks in dense radiation become mediated by photons

$\Rightarrow$  fundamental change in the shock dissipation mechanism

upstream

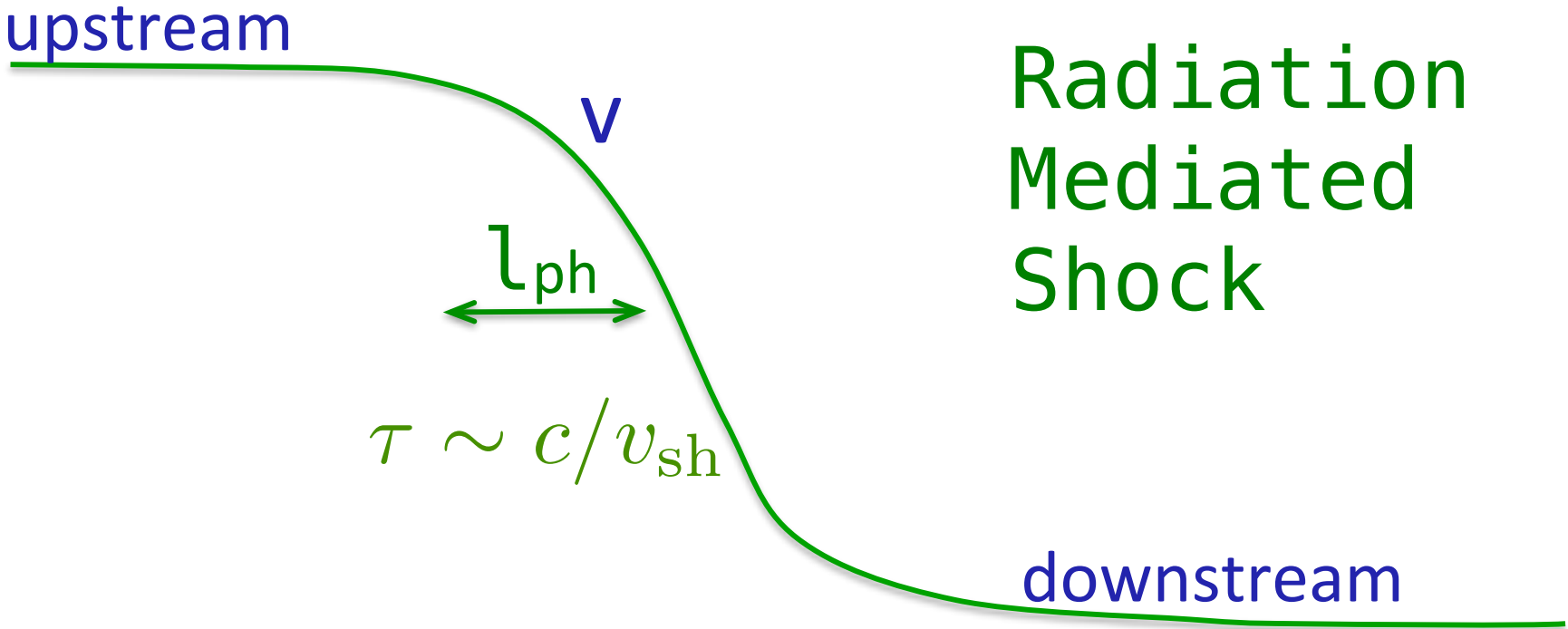
$v$

$\ell_{\text{ph}}$

$$\tau \sim c/v_{\text{sh}}$$

Radiation  
Mediated  
Shock

downstream



upstream

$v$

$\ell_{\text{ph}}$

$$\tau \sim c/v_{\text{sh}}$$

Radiation  
Mediated  
Shock

downstream

nonthermal  
radiation  
(Fermi mechanism)

plasma stays cold

# Radiative MHD from first principles: “Photon In Cell”

Fluid motion: Lagrangian grid

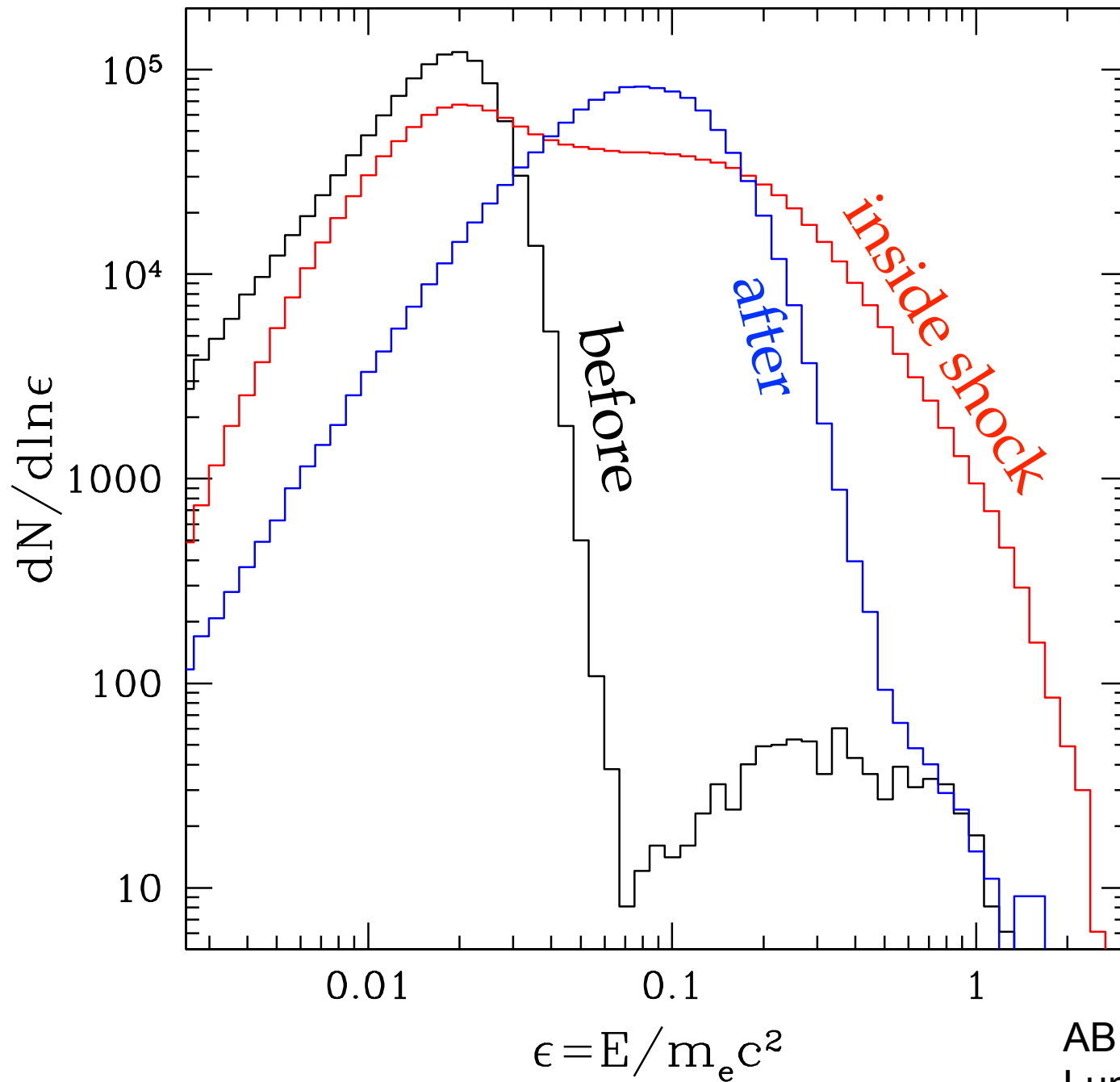
Radiative transfer: individual photons (Monte-Carlo)

**goal:** self-consistent solution: radiation + shock structure

AB 2017

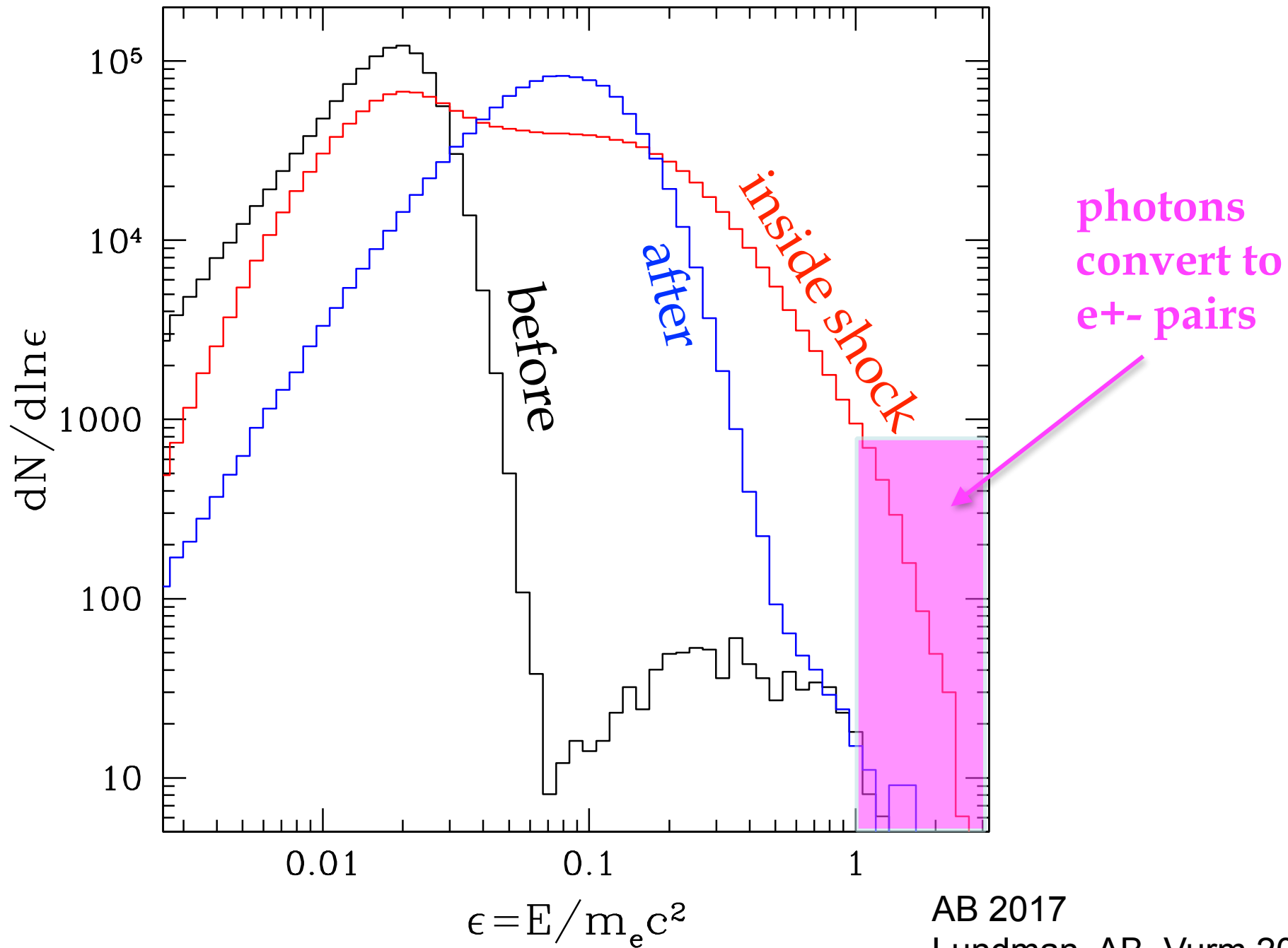
Lundman, AB, Vurm 2018

Ito et al. 2018



AB 2017  
Lundman, AB, Vurm 2018





# Bulk Comptonization in non-relativistic shocks

$$\left. \begin{array}{l} \text{inverse} \\ \text{Compton:} \quad \frac{\Delta\epsilon}{\epsilon} \sim \frac{v^2}{c^2} \\ \\ \text{recoil:} \quad \frac{\Delta\epsilon}{\epsilon} \sim -\epsilon \end{array} \right\} \Rightarrow \epsilon_{\text{max}} \sim \frac{v^2}{c^2}$$

$$\epsilon = \frac{h\nu}{m_e c^2}$$

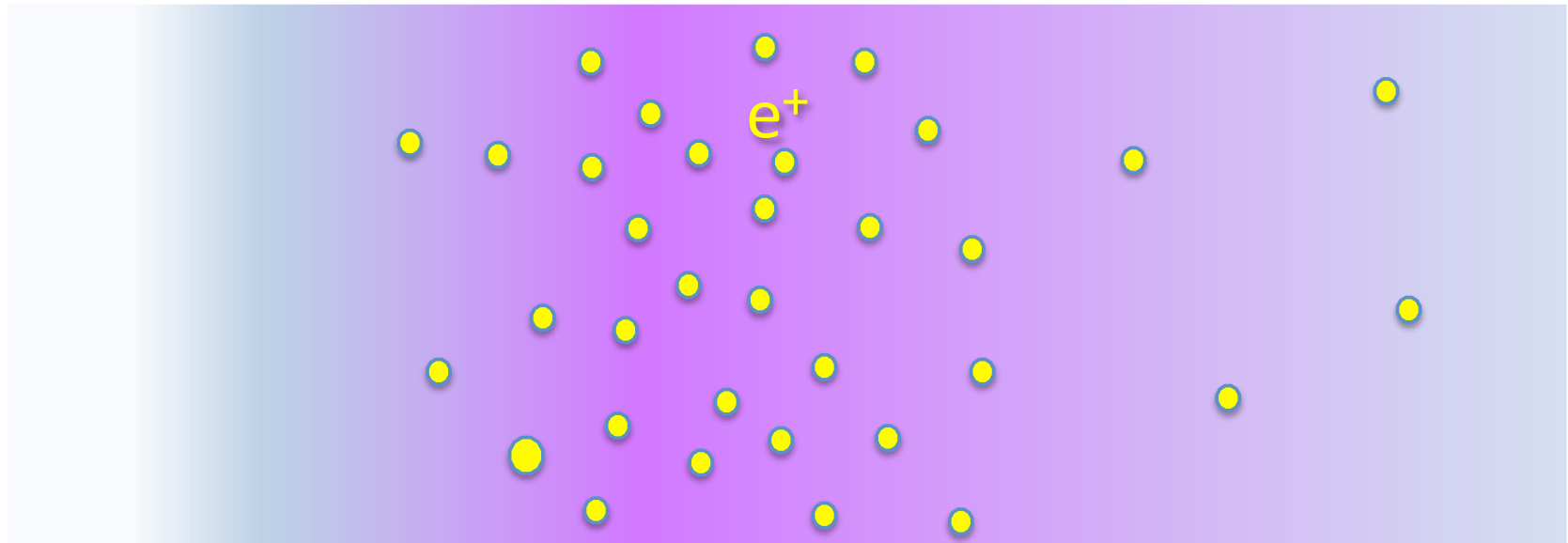
upstream

$v$

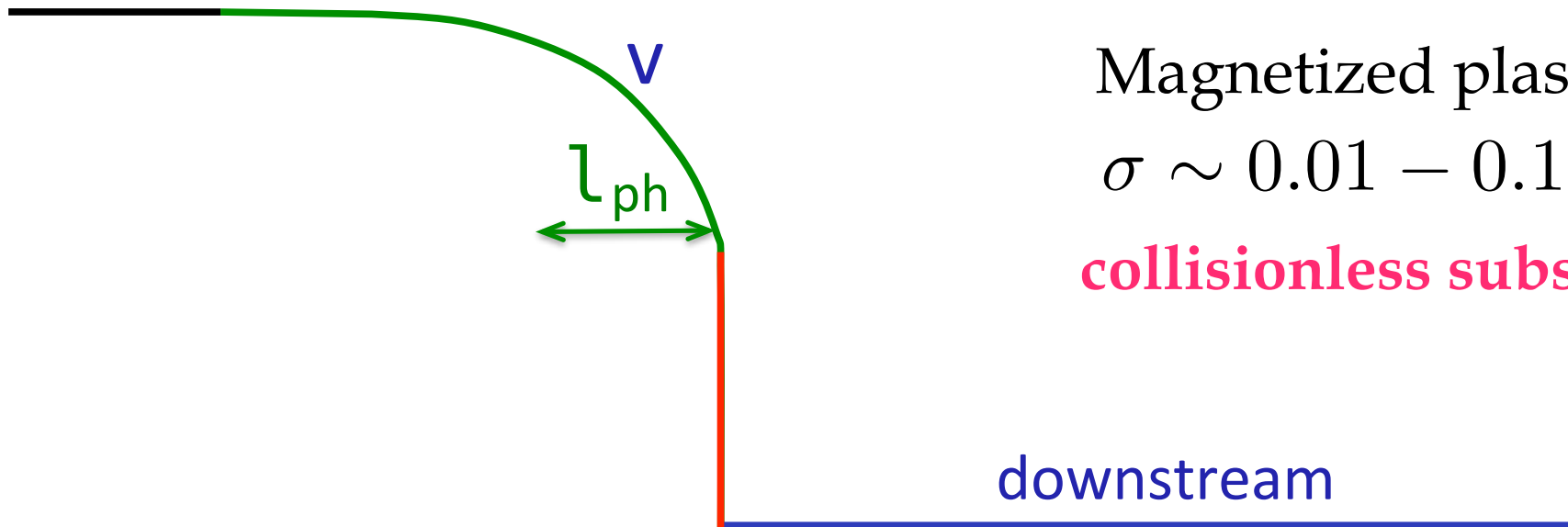
$l_{ph}$

$e^+e^-$  creation in  
relativistic RMS:  
 $l_{ph}$  shrinks by  $\sim 100$

downstream



$$Z_{\pm} = n_{\pm}/n_{\text{ion}} \sim 100 - 300 \quad \text{if} \quad \gamma_{\text{sh}}\beta_{\text{sh}} > 1$$

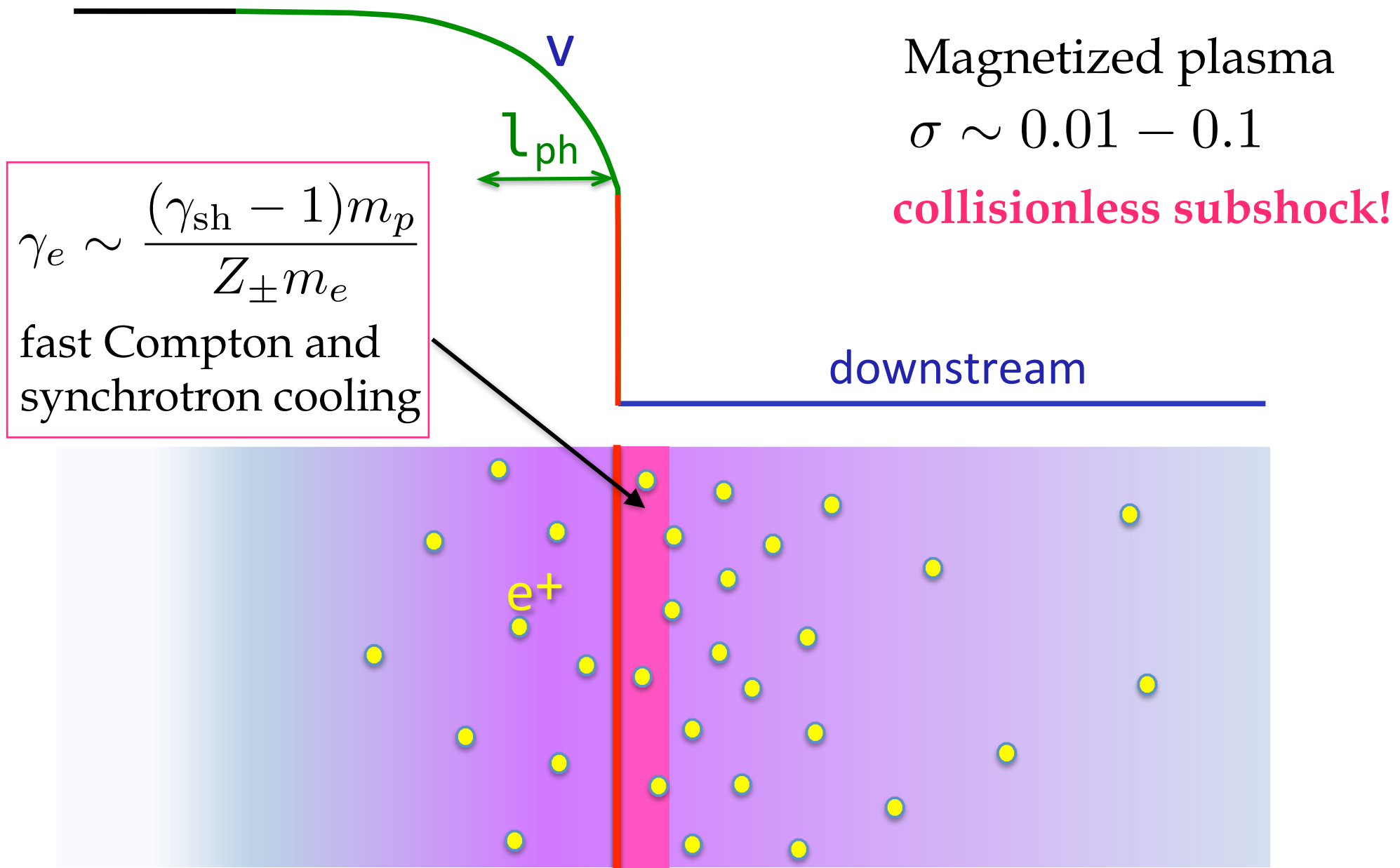


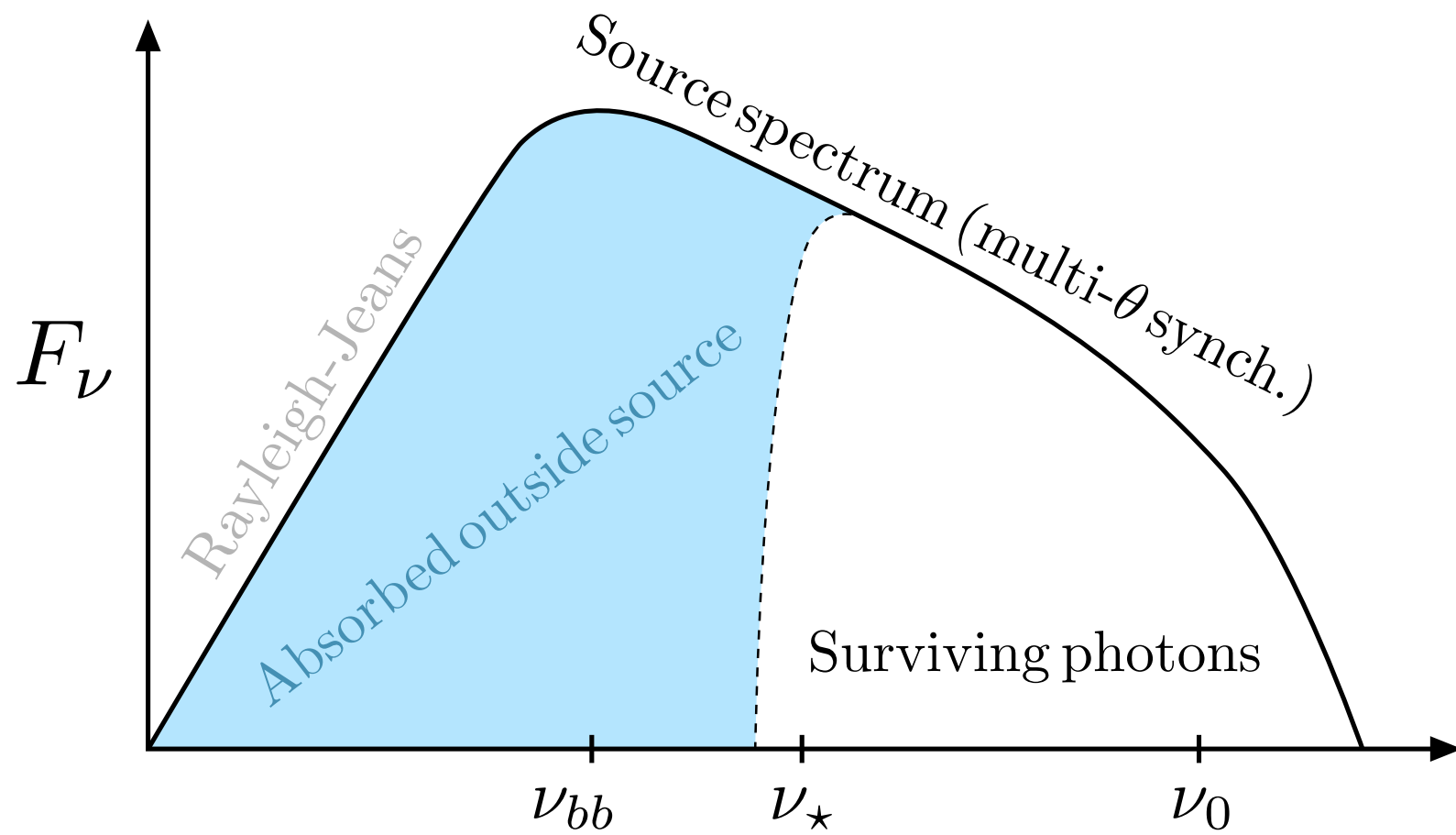
Magnetized plasma

$$\sigma \sim 0.01 - 0.1$$

**collisionless subshock!**

downstream





## Synchrotron photon number:

- peaks at low photon energies
- controlled by **induced down-scattering limit**:

$$\frac{kT_b}{m_e c^2} \tau_T^2 \sim 1 \quad (\tau_T \sim c/v_{\text{sh}})$$

$$\Rightarrow \boxed{\frac{n_{\text{ph}}}{n_{\text{ion}}} \approx \alpha \frac{m_p v_{\text{sh}}^2}{m_e c^2} \sigma \gamma_e^4}$$

$$\gamma_e \sim 30 \quad \left[ \sqrt{\gamma_e} - 2 - \ln \frac{\gamma_e}{4} \approx 1 \right]$$

(regulated by e<sup>+</sup>- creation)

# Shock self-organization

- Dresses itself in pair plasma:  $n_{\pm}/n_{\text{ion}} \sim 100$

=> “carries” the explosion photosphere  
(delayed shock breakout)

- Feeds itself with photons that mediate the shock:

$$n_{\text{ph}}/n_{\text{ion}} \sim 10^5 - 10^6$$



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2 dimensionless parameters:

$v_{\text{sh}}/c$  and medium magnetization  $\sigma$

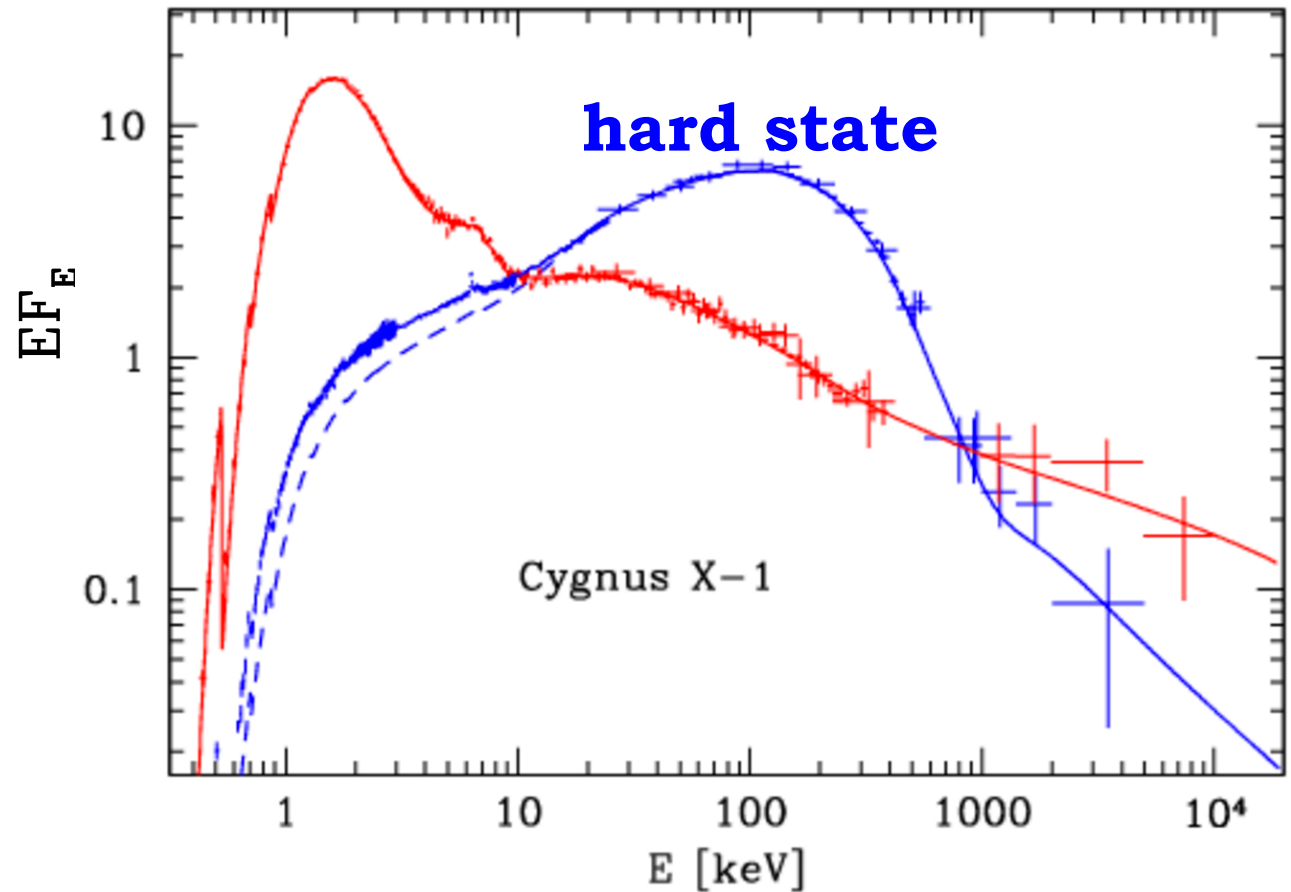
current work: shock spectrum emerging from photosphere

(Lundman & AB)

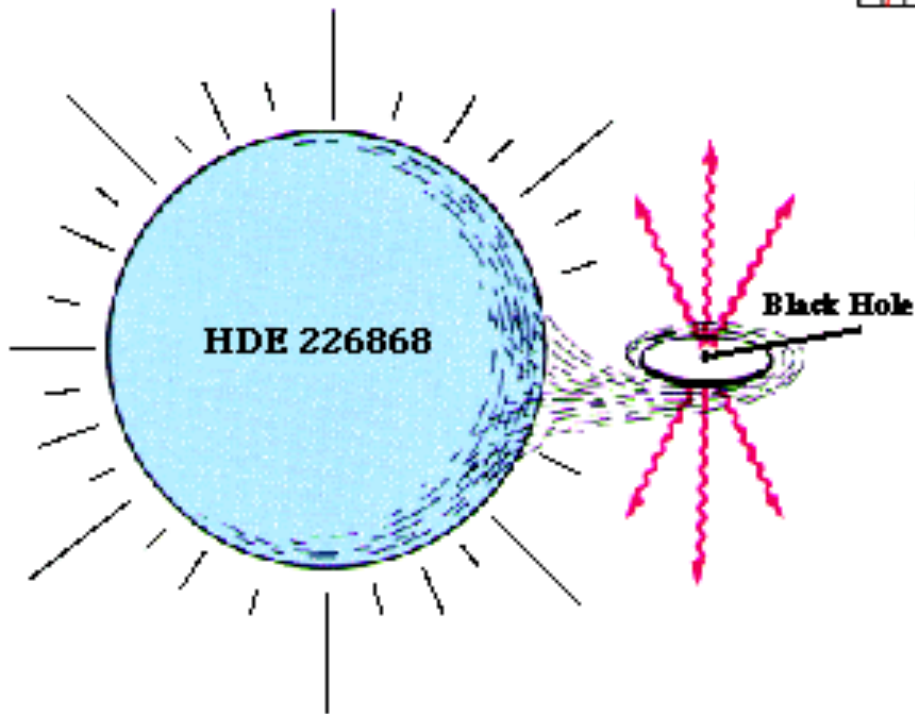
## II. Magnetic reconnection

- Magnetic flares near accreting black holes
- Flares in magnetars (fireballs)
- Reconnection in pulsar winds and BH corona/jets  
(X-ray binaries, AGN, GRBs)

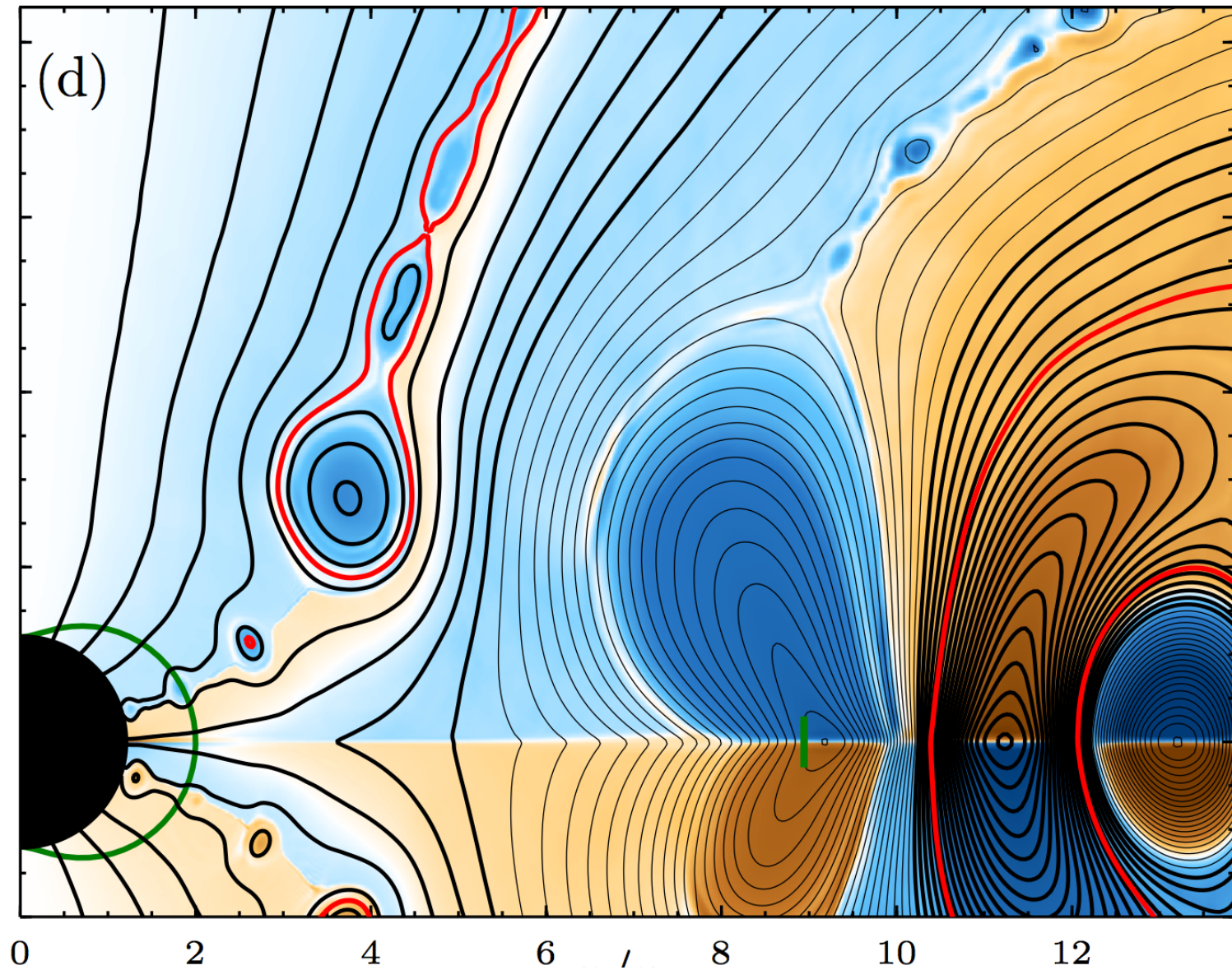
# Hard X-rays from accreting black holes



McConnell et al. 2002

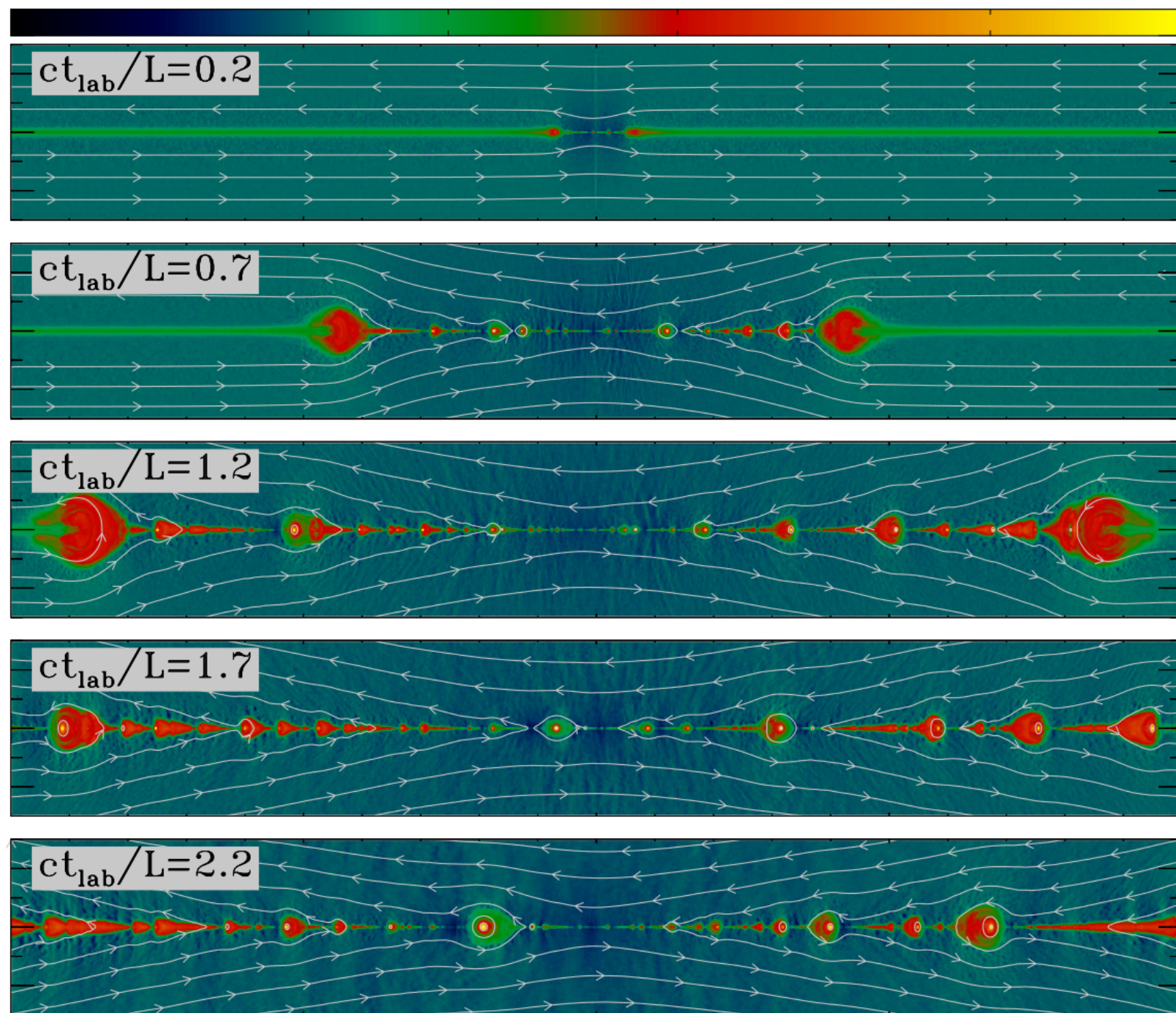


# Reconnection near black holes





# ab initio (PIC) simulations



Self-similar  
chain of  
plasmoids

$$v_{\text{rec}} \approx 0.1c$$

Uzdensky et al. 10

Melzani et al. 14

Sironi, Spitkovsky 14

Guo et al.16

Sironi et al.16

Werner et al.16

# Parameters of relativistic reconnection

1. Magnetization:  
(MHD)

$$\sigma = \frac{B^2}{4\pi\rho c^2} = \frac{2U_B}{\rho c^2}$$

$\gg 1$   
in accretion disk  
corona/jet

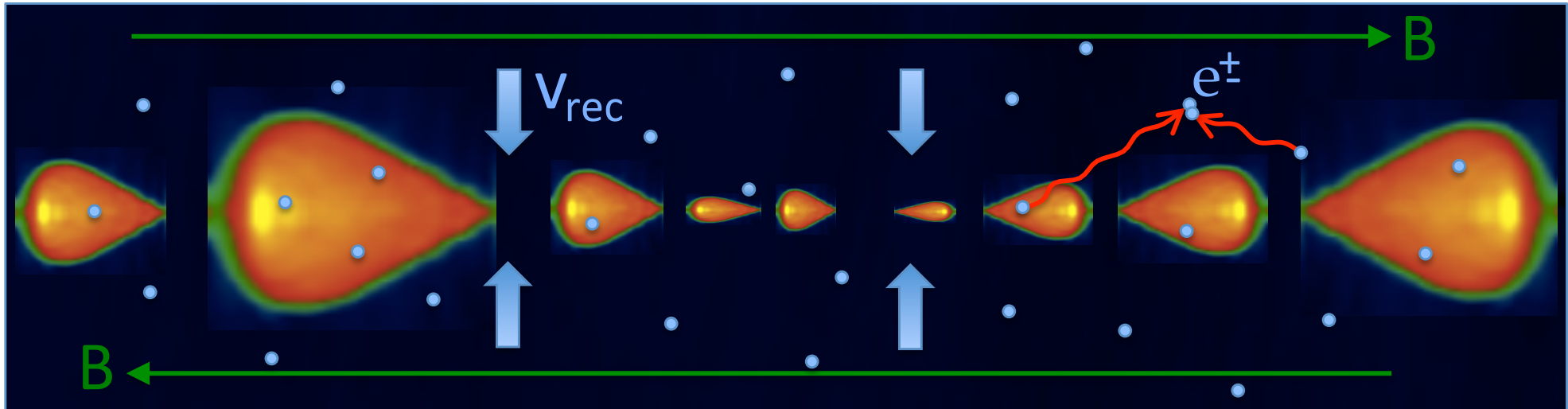
2. Compactness: cooling time vs. light crossing time  
(radiative)

$$\frac{t_{\text{cool}}}{s/c} \sim \frac{1}{\gamma_e \ell}$$

$$\ell \equiv \frac{U}{m_e c^2} \sigma_T s \sim \frac{m_p}{m_e} \frac{L}{L_{\text{Edd}}} \frac{R_S}{R} \sim 10^3$$

# Radiative reconnection (cooling time $\ll$ light crossing time)

AB 2017



1. Plasmoids are cool + fast  $\Rightarrow$  “chain Comptonization”
2. Energetic photons ( $>1$  MeV) convert to  $e^\pm$  pairs  
 $\Rightarrow$  reconnection layer self-feeds with plasma

$$\Rightarrow \tau_T \sim 1$$

# High-energy particles from X-points

X-point acceleration:  $t_X \sim \frac{\bar{\gamma}_e m_e c}{eB} \sim 10^{-9} \left( \frac{M}{10M_\odot} \right)^{-1/2} \frac{r_g}{c}$

Synchrotron cooling:  $t_{\text{syn}}^{\text{min}} \sim \frac{m_e c}{\sigma_T U_B \bar{\gamma}_e} \sim 10^{-5} \frac{r_g}{c} \quad (\gamma_e \sim \sigma)$



# Compton cooling of growing plasmoids

Inverse Compton:  $t'_{\text{IC}} \lesssim \frac{10}{\ell_B} \frac{r_g}{c}$

$$t'_{\text{age}} \sim 10(w/c) \quad \text{Sironi et al. 2016}$$

=> Strong cooling of plasmoids of size  $w \gg r_g/\ell_B$

# Plasmoid bulk motion

Magnetic stresses push plasmoids:  $f_{\text{push}} = \xi \frac{U_B}{w}$

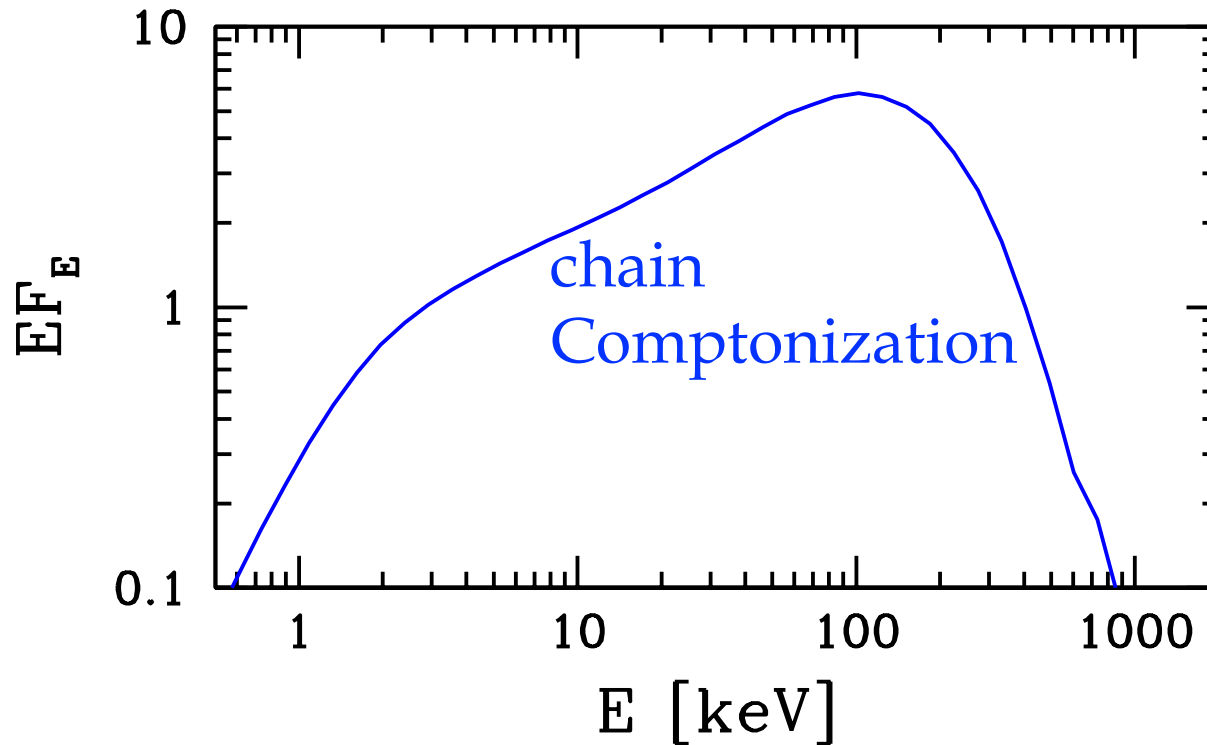
$$\frac{f_{\text{drag}}}{f_{\text{push}}} = \beta \gamma^2 \frac{\tau_{\text{pl}}}{\tau_{\star}} \quad \tau_{\star} \equiv \xi \frac{U_B}{U_{\text{rad}}} \approx \frac{\xi}{\beta_{\text{rec}}} \sim 1$$

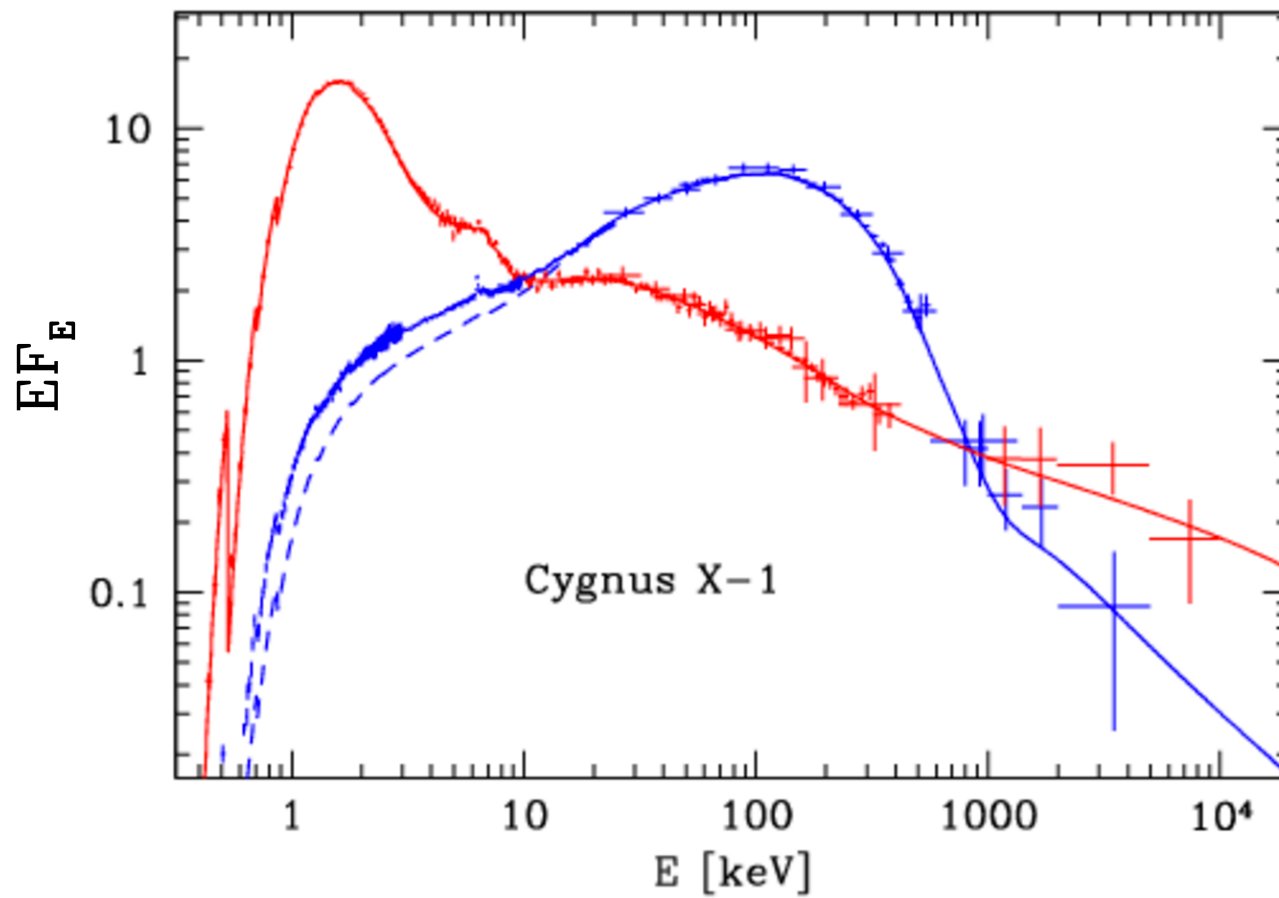
$$\tau_{\text{pl}} = n_{\pm} \sigma_{\text{T}} w$$

Drag-limited motion:  $\gamma \approx (\tau_{\star} / \tau_{\text{pl}})^{1/2} \quad (\gamma \leq \sigma^{1/2})$

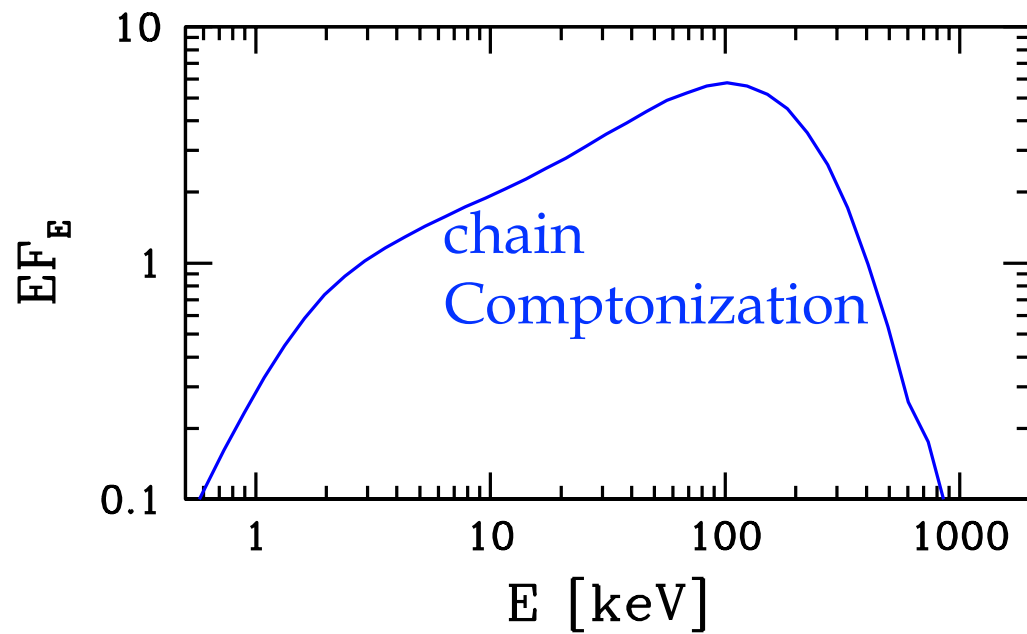
Power deposited in the plasmoid chain:  $\frac{dL}{d \ln a} \propto \begin{cases} a^{q/2}, & a < 1 \\ a^{q-1}, & a > 1 \end{cases} \quad a \equiv \beta^2 \gamma^2$

$$\frac{dN}{d \ln w} \propto w^{-q}$$

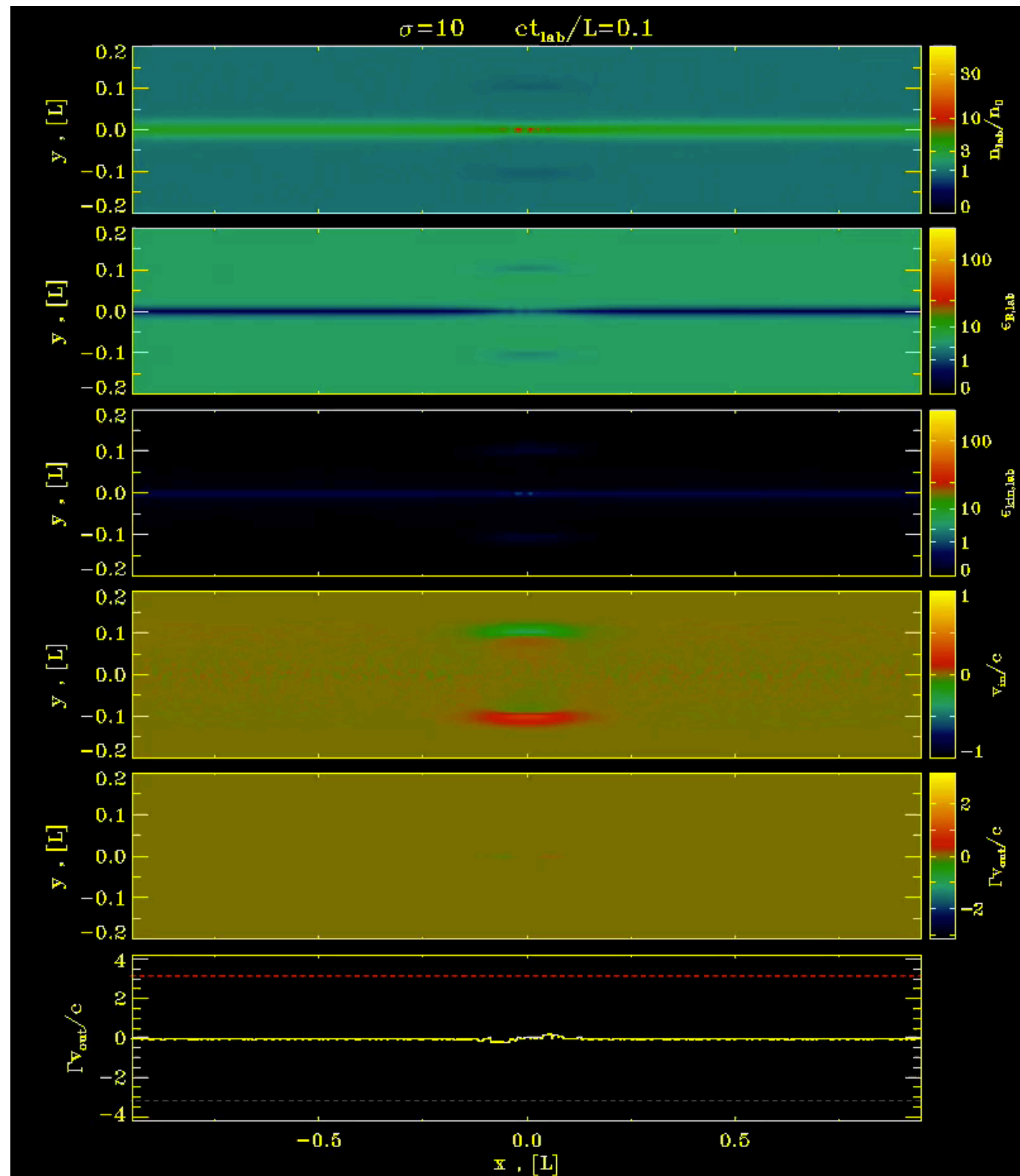




McConnell et al. 2002



Radiative  
PIC simulations  
(Sironi, AB in prep.)



### **III. Turbulence**

**(accretion disks, jets, mergers, magnetar flares)**

# Radiation damping of turbulence cascade

cascade power:  $\dot{Q} \sim \frac{h\rho v_0^3}{\ell_0}$       spectrum:  $v(\ell) = v_0 \left( \frac{\ell}{\ell_0} \right)^{1/3}$

$$\dot{U}_{\text{damp}}(\ell) \sim \frac{h\rho v^2(\ell)}{t_{\text{damp}}(\ell)} \qquad t_{\text{damp}}(\ell) \sim \frac{\ell^2}{\nu}$$

$$\nu \sim \frac{U}{U + \rho c^2} l_{\star} c$$

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$\dot{U}_{\text{damp}}(\ell) \sim \frac{h\rho v^2(\ell)}{t_{\text{damp}}(\ell)}$        $t_{\text{damp}}(\ell) \sim \frac{\ell^2}{\nu}$

$\nu \sim \frac{U}{U + \rho c^2} l_\star c$

$\dot{U}_{\text{damp}} = \dot{Q}$

$Re = \frac{\ell_0 v_0}{\nu}$        $\ell_{\text{damp}} \sim \ell_0 Re^{-3/4}$        $\ell_0 = \tau \ell_\star$

$\ell_{\text{damp}} \gg \ell_\star \iff Re \ll \tau^{4/3}$

– “viscous” regime  
shear Comptonization



Collisionless dissipation regime:  $Re > \tau^{4/3}$

$$t_{\text{damp}}(\ell) > t_{\text{turb}} = \frac{\ell}{v(\ell)}$$

if true at  $\ell \geq \ell_*$  then also true at all  $\ell$

$$t_{\text{damp}}(\ell) \sim \frac{U + \rho c^2}{U} \frac{\ell_*}{c} \times \begin{cases} (\ell/\ell_*)^2 & \ell > \ell_* \\ 1 & \ell < \ell_* \end{cases}$$

# Turbulent jets in GRBs

$\xi$  = turbulence injection power  
divided by total jet power

switch from viscous regime  
to collisionless dissipation:  $\tau_{\text{switch}} \sim \frac{1}{8\xi^2}$

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Compton

injected

energy balance:  $(\theta - \theta_C) \tau \approx \frac{\xi}{3} \frac{\rho c^2}{U}$

# Dissipation at high compactness: summary

- Luminosity is powered by dissipation of magnetic/kinetic energy in a compact region filled with dense radiation
- “Dissipation machine” (reconnection layer/shock) self-organizes into a non-linear state, feeding itself with  $e^+$ - plasma, photons, and generating the observed spectrum
- — a first-principle problem with few parameters can be isolated from “mud wrestling” — MHD weather around compact objects
- Best method of study: direct numerical experiment  
(PIC/Vlasov plasma + radiative transfer)