



Intrinsic edge rotation as governed by momentum transport due to neutrals



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Edge rotation

- Flow shear suppresses turbulence
- Rotation stabilizes MHD modes (e.g. resistive wall mode, neoclassical tearing mode)
- Edge sets the boundary condition for core rotation profile

Need to understand momentum transport in edge

- Neutrals present inside separatrix, large cross-field mobility

Role of neutrals?

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JET results:

- Pedestal toroidal rotation is higher in “corner” divertor configuration, and it is affected by fueling

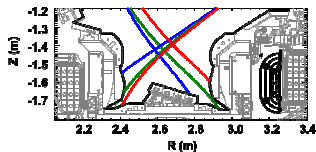


Fig 1a (above): Divertor geometry used in JET (horizontal=blue; corner=green; vertical=red). The cryogenic pump is also pictured in black.

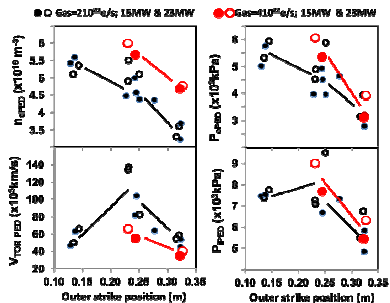


Fig 5: Clockwise from top to bottom: pedestal density, electron and ion pedestal pressure and toroidal velocity at the top of the pedestal for the three divertor configurations at different power and gas injection rate.

E. Joffrin et al., *25th IAEA Fusion Energy Conference* (2014).

Effect of geometry? TCV data

TCV results:

- Rotation profile shifts by changing X-point location

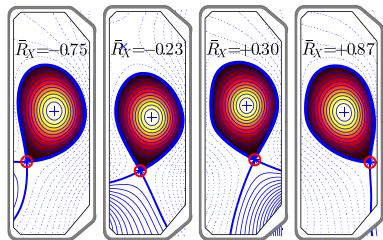


FIG. 1. Representative plasma geometries, varying \bar{R}_X , in LSN (lower row) and USN (upper row) configurations.

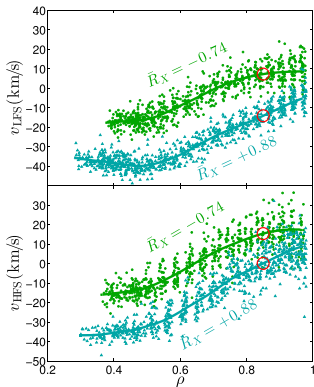


FIG. 6. LFS and HFS carbon velocity measurement points (dots) and fitted profiles (lines) for an HFS ($\bar{R}_X = -0.74$) and LFS ($\bar{R}_X = +0.88$) X-point. The approximate core-edge boundary location $\rho = 0.85$ is marked with red circles. CXRS measurement error bars are comparable to point scatter.

T. Stoltzfus-Dueck et al., *Physics of Plasmas* **22**, 056118 (2015).

Momentum transport and E_r

- Standard neoclassical theory, for low flow (i.e. subsonic)
 - Viscosity is small, higher order quantity
 - Particle fluxes automatically ambipolar to lowest order, independent of E_r
 - Need f_i very accurately to set E_r
- At the edge other effects can give sizable momentum transport
 - ion orbit losses
 - toroidal field ripple or resonant magnetic perturbations
 - sharp profile variations, $\rho_p/L_\perp \sim \mathcal{O}(1)$
 - **transport by neutral particles**

Neutral momentum transport

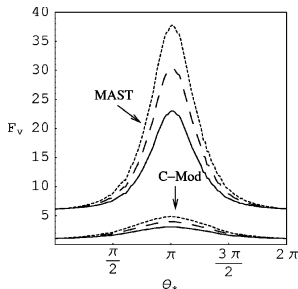
- Low concentration, but high mobility
 - Can affect momentum transport strongly
- **Result: E_r is determined by major radius of neutrals**
 - Control possible
 - e.g. fueling location¹
- Neutrals may be concentrated near X-point
 - Move X-point to change location of neutrals
- Consider intrinsic rotation (i.e. no torque) and steady state
 - ITER-relevant case
 - Momentum flux ($\propto n_n$) vanishes, independent of n_n

¹H-mode access in MAST easier with inboard gas-puff: A.R. Field et al., *PPCF* **46**, 981 (2004).

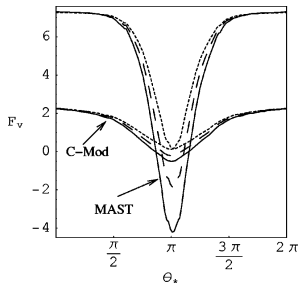
Analytical results

- Analytical results: P. Helander, T. Fülöp, and P. J. Catto, *Physics of Plasmas* **10**, 4396 (2003).

Pfirsch-Schlüter



banana



$$V_\zeta \propto \frac{dT_i}{d\psi} F_V, \text{ neutral density is localized to } \theta_*$$

Neutral momentum flux

- Charged particle momentum equation, no torque, steady state

$$\langle e_a \Gamma_a \cdot \nabla \psi \rangle = \underbrace{\langle R \hat{\phi} \cdot \nabla \cdot \pi_a \rangle}_{\approx 0} - \sum_b \langle R \hat{\phi} \cdot \mathbf{R}_{a,b} \rangle$$

- For neutrals, viscosity balanced by friction on ions

$$0 = \langle R \hat{\phi} \cdot \nabla \cdot \pi_n \rangle - \langle R \hat{\phi} \cdot \mathbf{R}_{n,i} \rangle$$

- Require ambipolar flux

$$0 = \sum_{a \& n} \langle e_a \Gamma_a \cdot \nabla \psi \rangle \approx \langle R \hat{\phi} \cdot \nabla \cdot \pi_n \rangle$$

$$\Rightarrow \langle R \hat{\phi} \cdot \pi_n \cdot \nabla \psi \rangle = 0 \quad \text{with no external torque}$$

- Gives constraint for E_r

Short mean-free-path neutral kinetic equation

- Neutral dynamics is dominated by charge exchange with ions

$$\mathbf{v} \cdot \nabla f_n = C_{CX}(f_n) \approx \frac{1}{\tau} \left(\frac{n_n}{n_i} f_i - f_n \right)$$

where $\tau^{-1} = n_i \langle \sigma v \rangle_{CX}$

- Solve perturbatively in short CX mean-free-path
 - $\lambda_n \approx 0.8 \text{ cm}$ for $n_i = 10^{20} \text{ m}^{-3}$, $\sigma_{CX} = 6 \times 10^{-15} \text{ cm}^2$
- f_n is calculated in terms of f_i
- Momentum flux moment of f_n written as moment of f_i

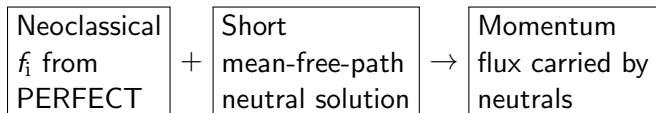
[P.J. Catto, P. Helander, J.W. Connor, and R.D. Hazeltine, *Phys. Plasmas* **5**, 3961 (1998)]

Solution for ion distribution

- Use PERFECT code² to compute f_i
 - neglect finite orbit-width effects here (radially local)
- Inputs to each PERFECT call:
 - magnetic geometry
 - ion temperature gradient $d_\psi T_i$
 - radial electric field $d_\psi \Phi_0$
- Outer loop iterates to find $d_\psi \Phi_0$, calling PERFECT to compute f_i

²M. Landreman, F.I. Parra, P.J. Catto, D.R. Ernst, and I. Pusztai, *PPCF* **56**, 045005 (2014).

Neutral momentum transport



- Iterate to find E_r so that momentum flux vanishes
- Limitations
 - Neutrals dominate momentum transport

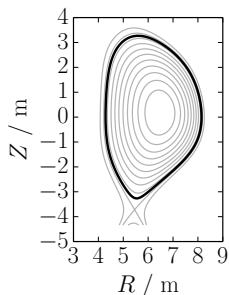
$$\frac{n_n}{n_i} \gtrsim 3 \cdot 10^{-7} \frac{L_n}{L} \frac{n_i^{2[10^{19}/\text{m}^3]} q^2}{B_{[T]}^2 T_{i[\text{keV}]}} \sim 10^{-6} - 10^{-4}$$

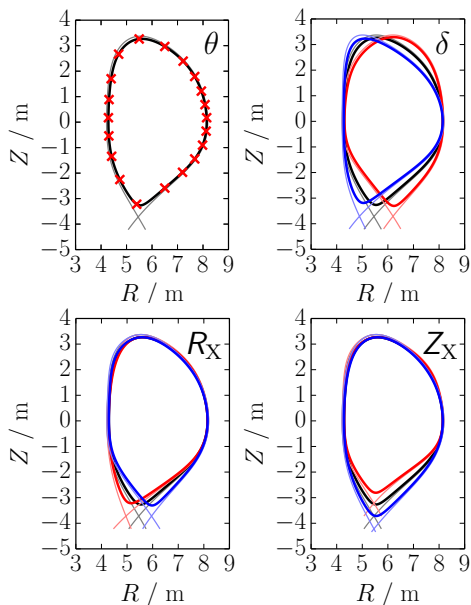
- Ion distribution is not directly affected by neutrals

$$\frac{n_n}{n_i} \lesssim 0.08 \frac{L_n}{R} \frac{n_i^{[10^{19}\text{m}^{-3}]} \sqrt{T_{i[\text{keV}]}}}{B_{[T]}} \sim 10^{-3}$$

- Short neutral CX mean-free-path, $\lambda_{\text{CX}} \ll L$
- Local neoclassical solutions $L \gg \rho_{pi}$ (no pedestal)

- Investigate effect of poloidally localized neutral density
- ITER-like parameters, near separatrix
 - Model equilibria [Cerfon & Freidberg]
 - Simulate $\psi_N = 0.95$ surface
 - Baseline parameters $n_i = 10^{20} \text{ m}^{-3}$,
 $T_i = 300 \text{ eV} \Rightarrow \nu_* = 6.95$
- Scanning magnetic geometry and/or neutral location

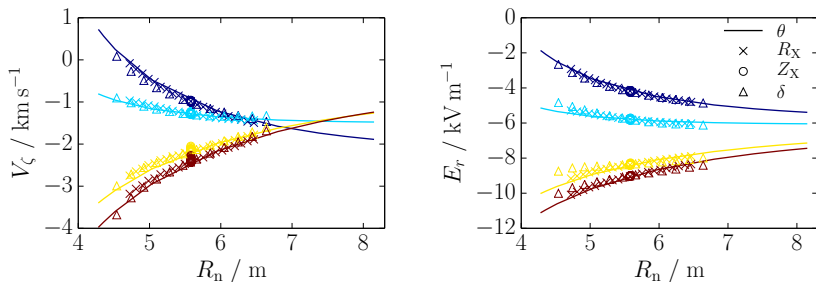




- θ : baseline geometry, neutrals located at different poloidal angles
- δ , R_X and Z_X : neutrals located at bottom of the flux surface (close to X-point), triangularity, and the horizontal and vertical position of the X point is varied, respectively.

Results: major radius of neutrals determines E_r , rotation

- All 4 types of scan collapse, R_n determines the behavior³



Baseline parameters: cyan $n = 10^{20} \text{ m}^{-3}$, $T_i = 300 \text{ eV}$, $\nu_* = 6.95$

Others: blue $\nu_* = 0.695$; yellow $\nu_* = 69.5$; red $\nu_* = 695$

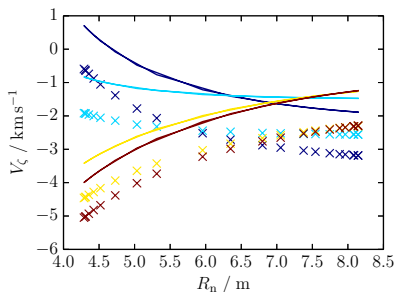
$1/L_T = 10 \text{ m}^{-1}$, $1/L_n = 0$

(arbitrary choices, but $V_\zeta \propto 1/L_T$ and E_r is simply offset by a/L_n)

³J. Omotani, I. Pusztai, T. Fülöp (2016) "Plasma rotation from momentum transport by neutrals in tokamaks" *arXiv:1604.08028*

Impurity rotation

- We can also add an impurity species, here trace Carbon ($Z = 6$, $n_C/n_e = 0.0004$)
- Show scan in poloidal position of neutrals:



Lines Deuterium; Crosses Carbon

Baseline parameters: cyan $n = 10^{20} \text{ m}^{-3}$, $T_i = 300 \text{ eV}$, $\nu_* = 6.95 Z_{\text{eff}}$

Others: blue $\nu_* = 0.695 Z_{\text{eff}}$; yellow $\nu_* = 69.5 Z_{\text{eff}}$; red $\nu_* = 695 Z_{\text{eff}}$

Summary

- Calculate E_r and rotation when neutral viscosity dominates
- Major radius of neutrals controls solution
- Ways to influence edge rotation:
 - move X-point
 - change fueling location

J. Omotani, I. Pusztai, T. Fülöp (2016) “Plasma rotation from momentum transport by neutrals in tokamaks” *arXiv:1604.08028*

Future work

- Experimental modeling
 - Rotation measurements
 - Inboard vs. outboard fueling?
 - Move X-point (but same target configuration)?
- Theory/code developments
 - Allow for torque from neutral beam injection
 - Higher neutral density
 - Affects ion distribution function
 - Density pedestal
 - Requires radially-global solution
 - Non-linear solve to obtain $\Phi_0(\psi)$