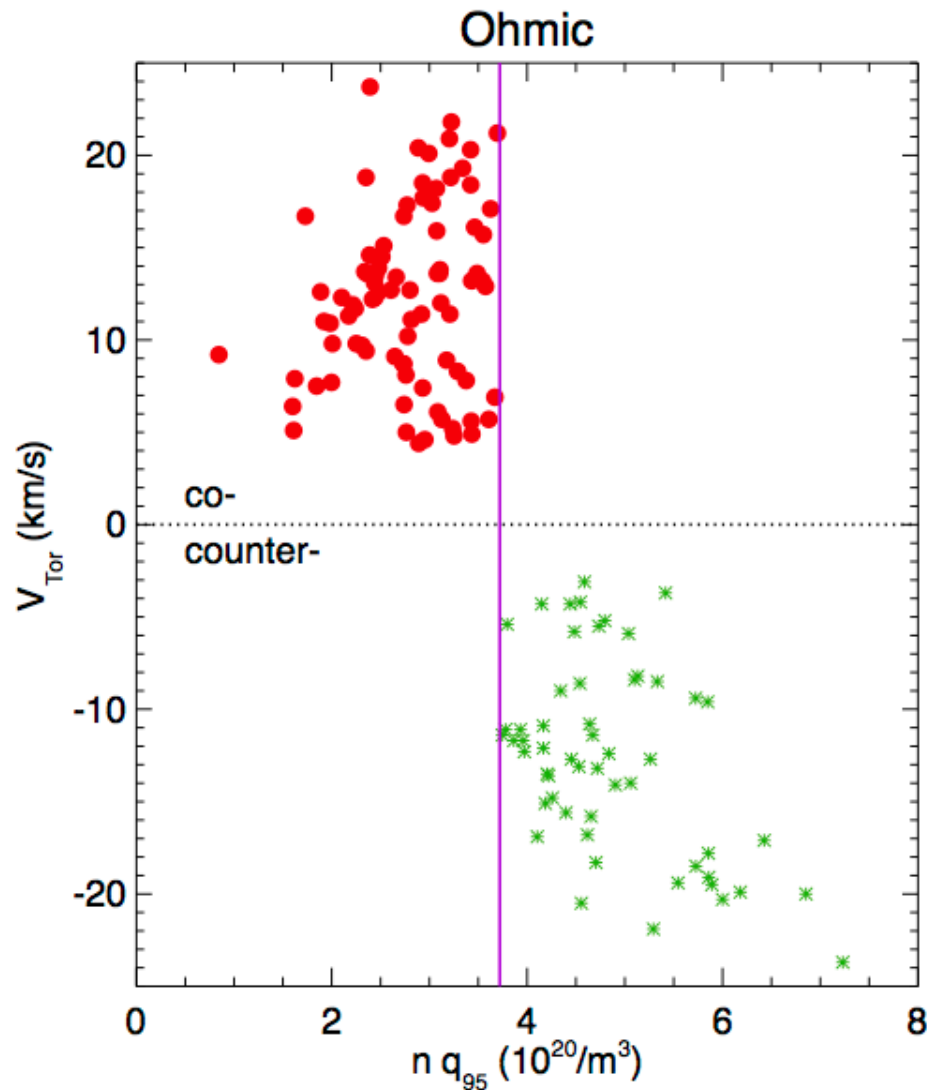


Intrinsic rotation reversals in C-Mod



Rotation changes direction as density and current ($\sim 1/q$) change

$$v_* \doteq \frac{q R_0}{v_{ti}} v_{ii} \propto n q$$

Turbulence model with diamagnetic effects

Solve fluctuation dynamics with gyrokinetic turbulence code GS2:

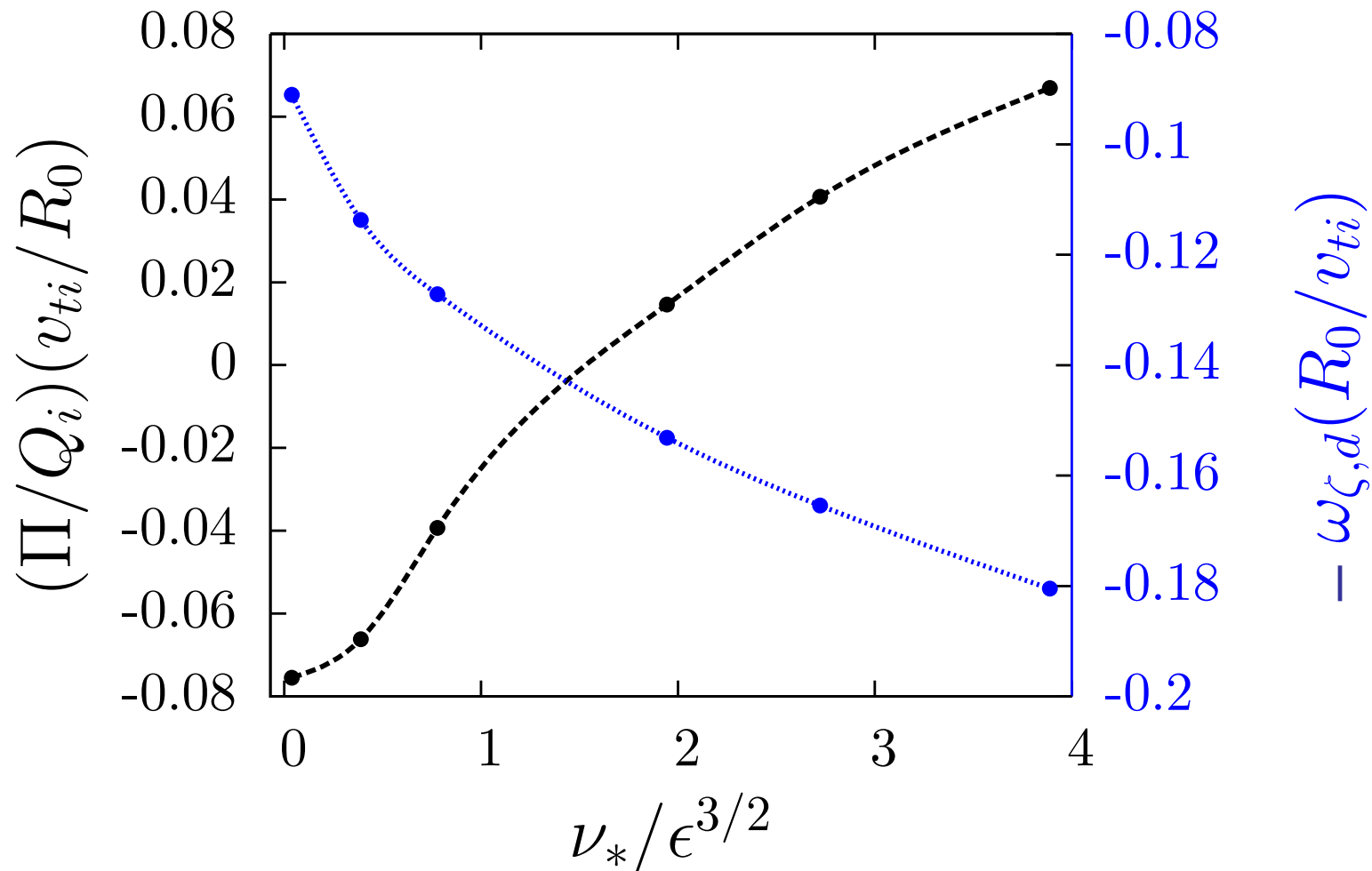
$$\begin{aligned} \frac{Dg_s}{Dt} + (\mathbf{v}_{\parallel} \cdot \nabla_{\parallel} + \mathbf{v}_{Ds} \cdot \nabla_{\perp}) \left(g_s - Z_s e \langle \varphi \rangle \frac{\partial \hat{F}_s}{\partial \varepsilon} \right) + C_s [g_s] \\ = - \langle \delta \mathbf{v}_E \rangle \cdot \left(\nabla \hat{F}_s + \frac{m_s R v_{\parallel}}{T_s} F_{Ms} \nabla \omega_{\zeta, E} \right) + Z_s e \mathbf{v}_{\parallel} \cdot \nabla \hat{\Phi} \frac{\partial g_s}{\partial \varepsilon} \end{aligned}$$

Obtain $F_1(r)$ and $\Phi_1(r)$ of $\hat{F} = F_0 + F_1$ and $\hat{\Phi} = \Phi_0 + \Phi_1$ by solving drift kinetic equation with NEO at multiple radii

Set $\omega_{\zeta, E} = -\omega_{\zeta, d}$, with $\langle (R^2 \mathbf{v} \cdot \nabla \zeta) F_{1i} \rangle_{\Lambda} = n_i \langle R^2 \rangle_{\Lambda} \omega_{\zeta, d}$

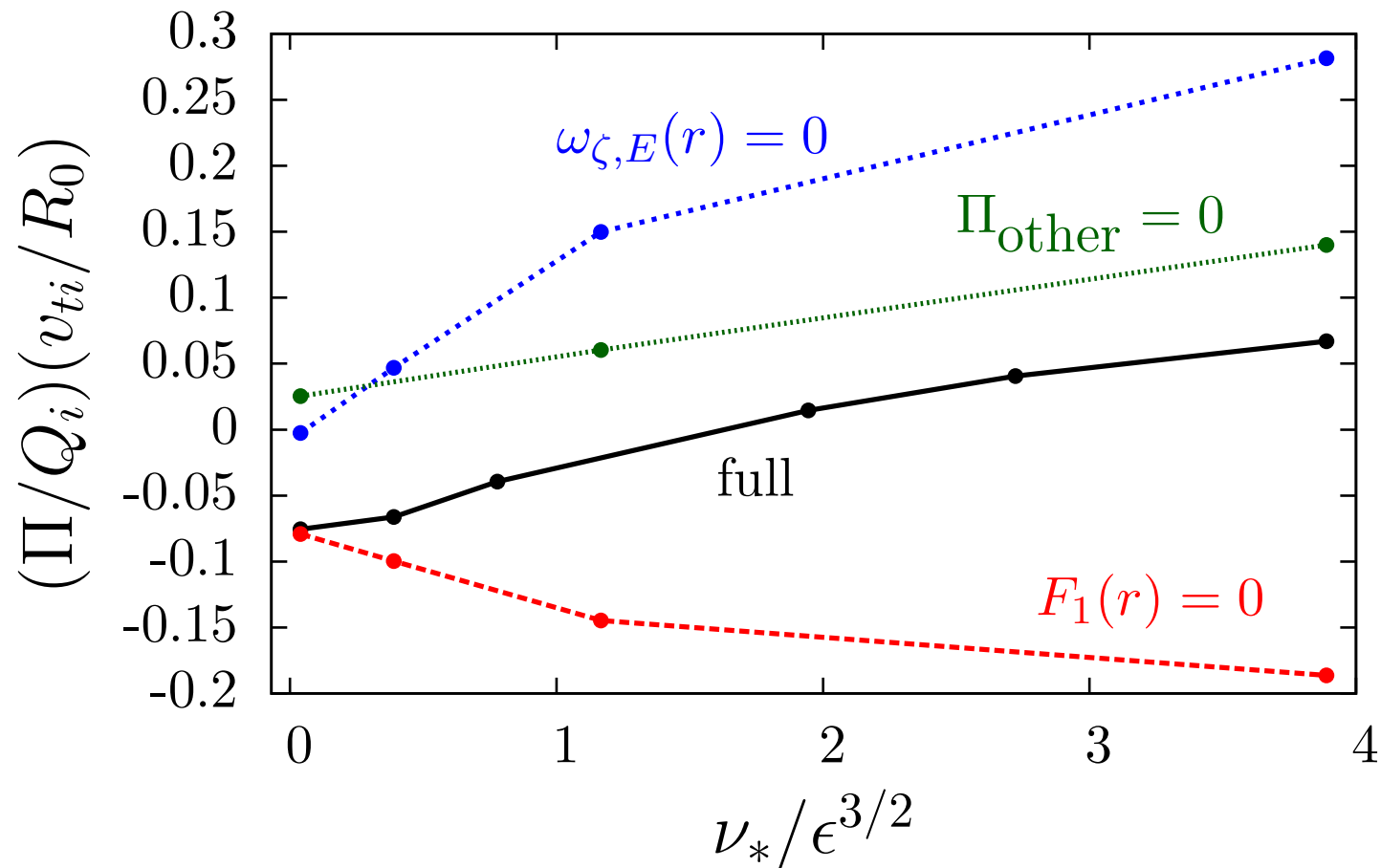
Momentum transport reverses direction with increasing collisionality

Transition occurs for $\nu_* \sim \epsilon^{3/2}$, corresponding to a transition from collisionless to collisional trapped particles (banana to plateau)



Opposing contributions from ExB and F_1

Transport is a result of partial cancellation between contributions from ExB advection-diffusion (red) and diamagnetic effects (blue)



Decomposition into advection and diffusion

$$\begin{aligned}\Pi &= -mR_0^2 \left(\chi_{\phi,d} \frac{\partial \omega_{\zeta,d}}{\partial r} + \chi_{\phi,E} \frac{\partial \omega_{\zeta,E}}{\partial r} \right) \\ &\quad - mR_0 (P_d \omega_{\zeta,d} + P_E \omega_{\zeta,E}) + \Pi_{\text{other}} \\ &= -mR_0^2 \chi_{\phi,\text{eff}} \frac{\partial \omega_{\zeta,d}}{\partial r} - mR_0 P_{\text{eff}} \omega_{\zeta,d} + \Pi_{\text{other}}\end{aligned}$$

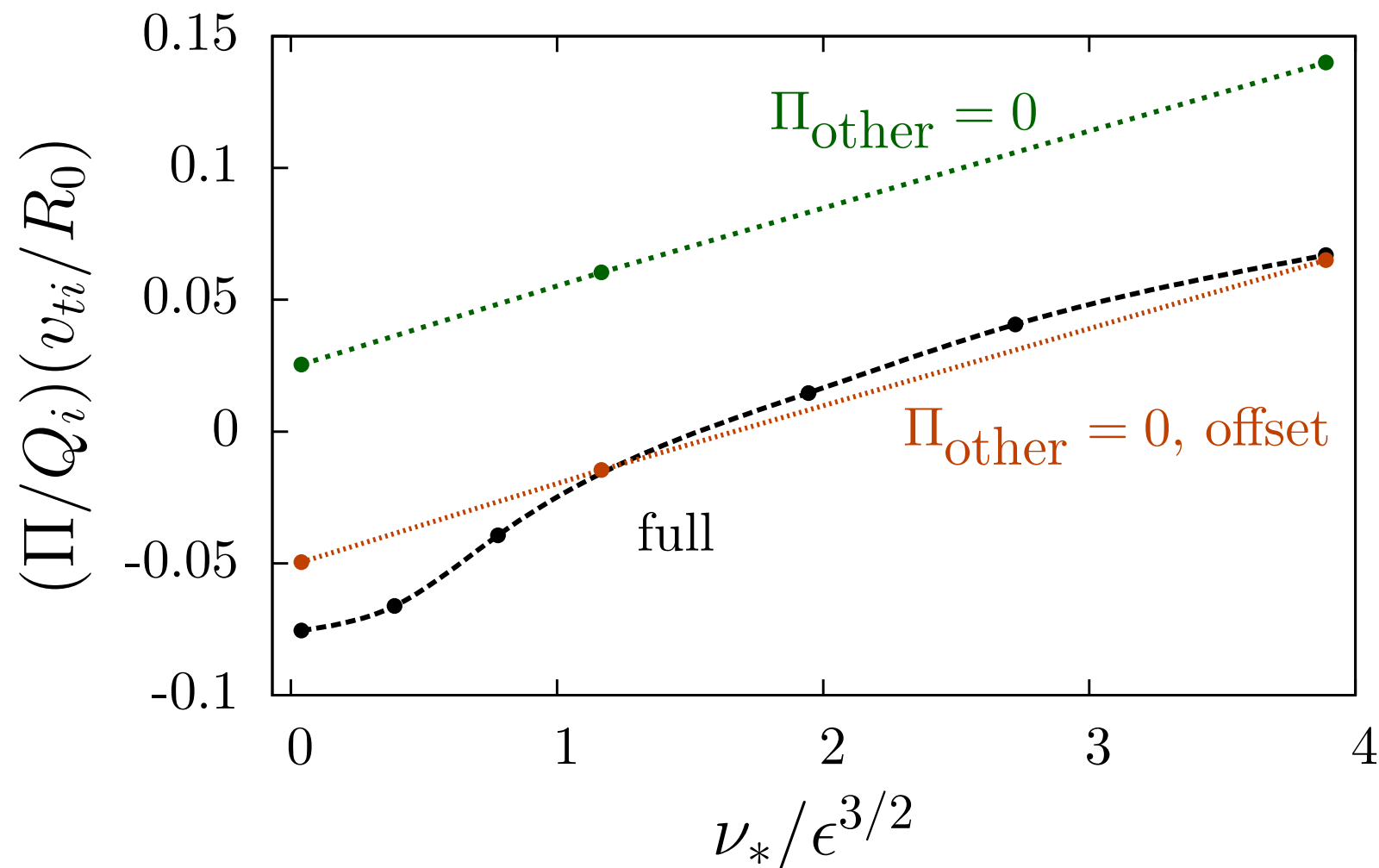
$$\omega_{\zeta,E}(r) = -\omega_{\zeta,d}(r)$$

$$\chi_{\phi,\text{eff}} \doteq \chi_{\phi,d} - \chi_{\phi,E}$$

$$P_{\text{eff}} \doteq P_d - P_E$$

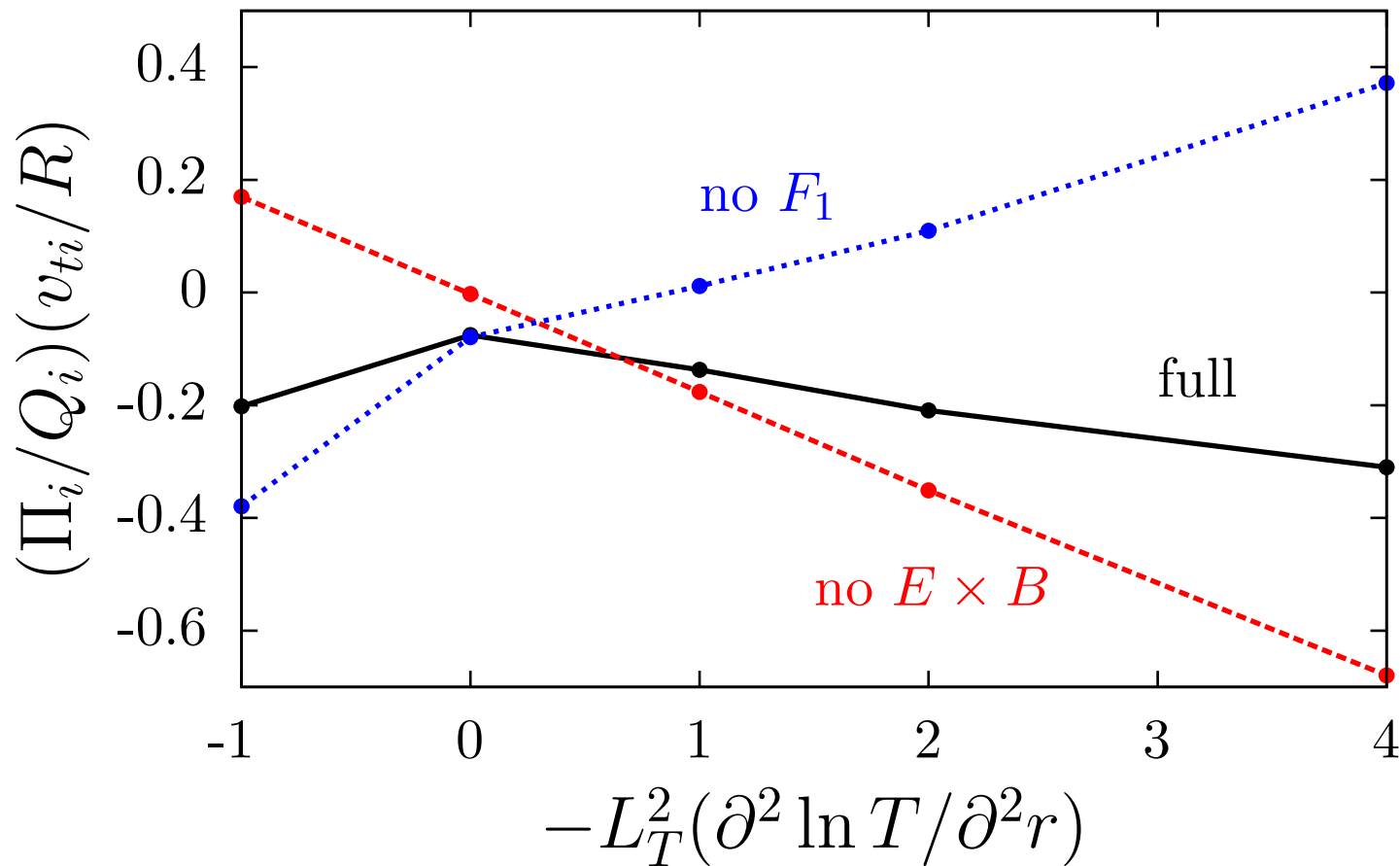
Decomposition into advection and diffusion

$$\Pi = -mR_0^2 \chi_{\phi, \text{eff}} \frac{\partial \omega_{\zeta, d}}{\partial r} - mR_0 P_{\text{eff}} \omega_{\zeta, d} + \Pi_{\text{other}}$$



Dependence on temperature profile

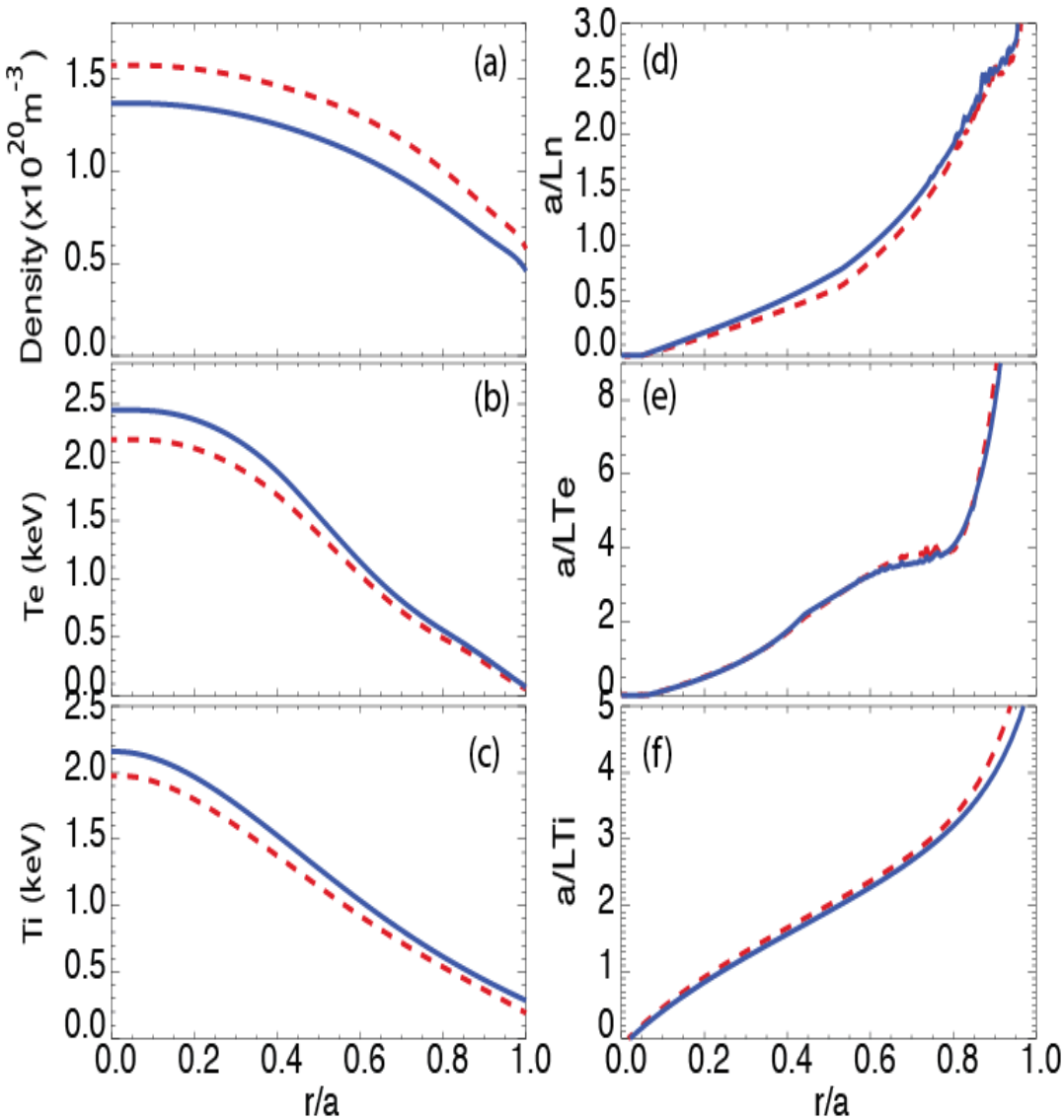
Relatively weak dependence on temperature profile for case considered here



Conclusion

- Inclusion of diamagnetic corrections to equilibrium Maxwellian distribution introduces strong dependence of momentum transport on collisions and weaker dependence on curvature of temperature profiles
- Behavior consistent with experimentally observed intrinsic rotation reversals
- Turbulence drive type (ITG, TEM, etc.) found to have little effect on intrinsic rotation in this study

Intrinsic rotation reversals in C-Mod



Density and temperature profiles modified slightly, gradient scale lengths kept fixed

Rotation reverses direction across most of radius

