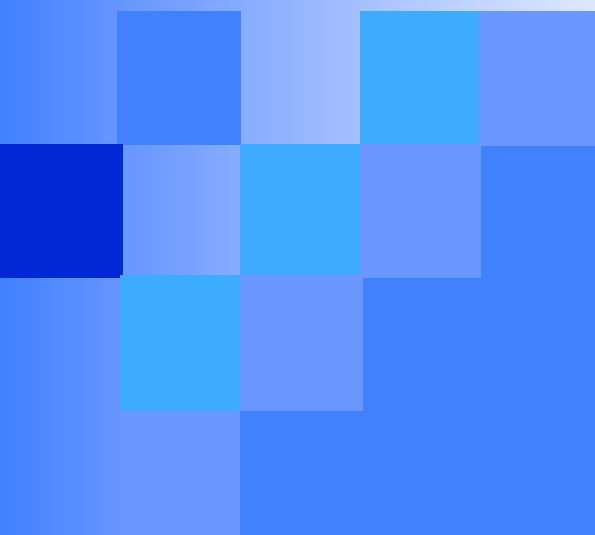


F. Jenko (on behalf of H. Doerk)

with T. Görler, D. Told, D. Hatch, M. Pueschel



Nature and role of microtearing turbulence in conventional tokamaks

Max-Planck-Institut für Plasmaphysik, Garching

Nth Vienna meeting ($N \gg 1$, thanks to Alex)
March 19th-30th, 2012

Outline of the talk

- **Historical context**
- Studies of the linear microtearing **instability**
 - expect existence in conventional tokamaks?
 - what are the critical plasma parameters?
- Nonlinear dynamics: microtearing **turbulence**
 - magnetic fluctuation amplitudes?
 - role of magnetic stochasticity?
- Microtearing modes in **finite β ITG/TEM** turbulence

A brief history of microtearing research

- **1973:** Small magnetic perturbations can lead to **stochastic fields** and **enhanced electron heat flux** ($\tilde{B}/B_0 \sim 10^{-4}$) [Stix]
- **1975-77:** Microtearing modes driven by ∇T_e are a **possible source** [Hazeltine et al., Drake et al., Chen et al.]
- **1980s and 1990s:** Linear theory of microtearing modes [Catto, Connor, Cowley, Drake, Hassam, Hastie, ...]
- **Since 2000:** Linear gyrokinetic simulations; microtearing modes responsible for electron heat flux in **spherical tokamaks**? [Kotschenreuther et al., Redi et al., Applegate et al.]
- **2007/08:** Microtearing modes also in **medium aspect ratio tokamaks**? [Applegate et al., Vermare et al., Told et al.]
- **2010/11:** **Nonlinear gyrokinetic microtearing** simulations point to relevance for NSTX [Guttenfelder] and AUG (as well as ITER) [this work]

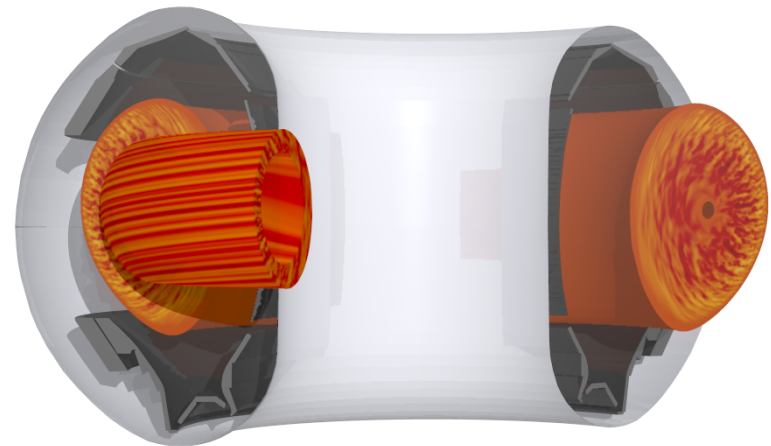
Scope of this work

Key questions:

- Properties of **linear microtearing instabilities**
- Electromagnetic **heat transport** due to microtearing turbulence
- **Nonlinear saturation** mechanism

Strategy:

- **Gyrokinetic simulations**
- Influence of **plasma parameters** (temperature gradient, beta, collisions...)
- Comparison to **analytical models**



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GENE

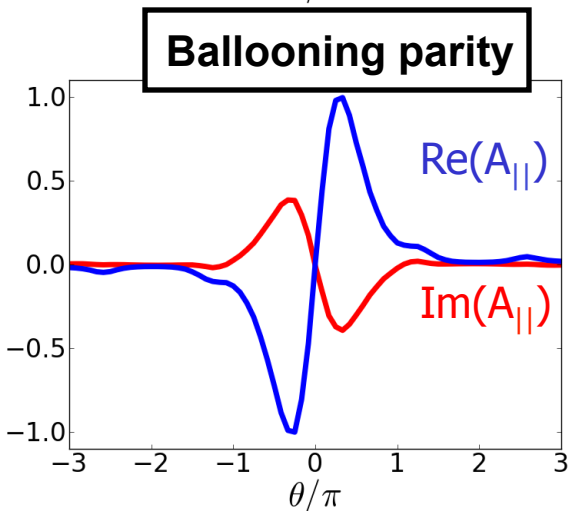
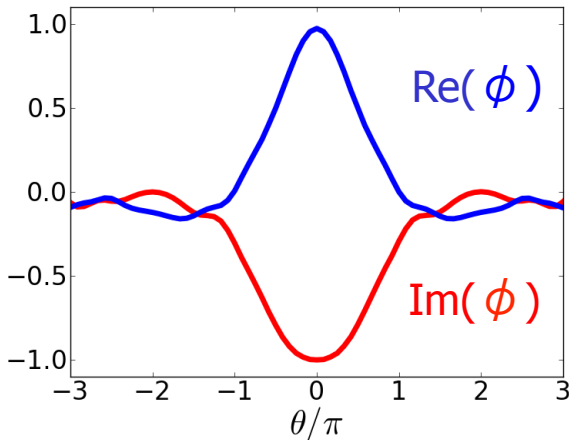
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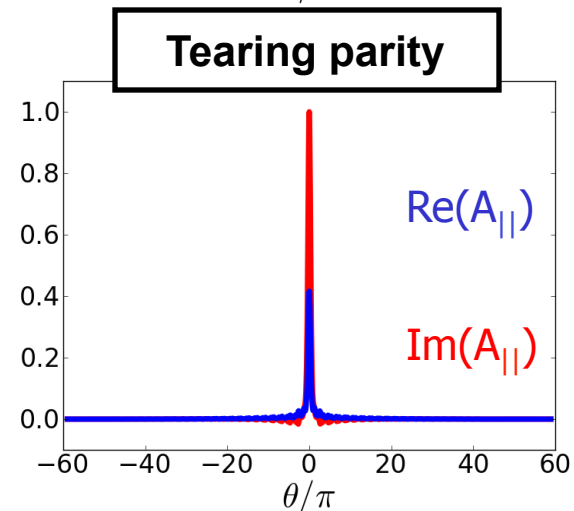
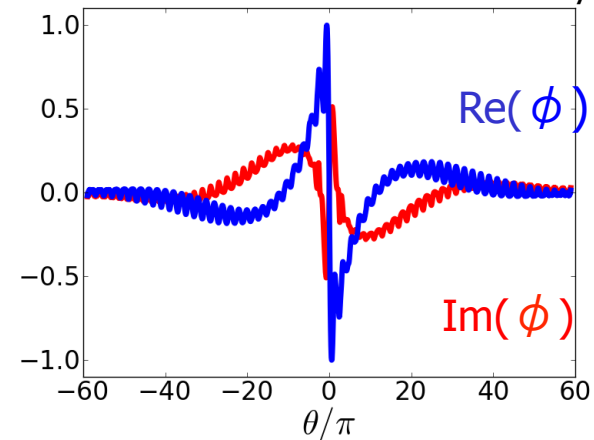
Studies of the linear microtearing instability

Comparison of microinstabilities in ballooning space

ITG modes (Ion Temperature Gradient)



Microtearing modes ($k_y=0.12$)

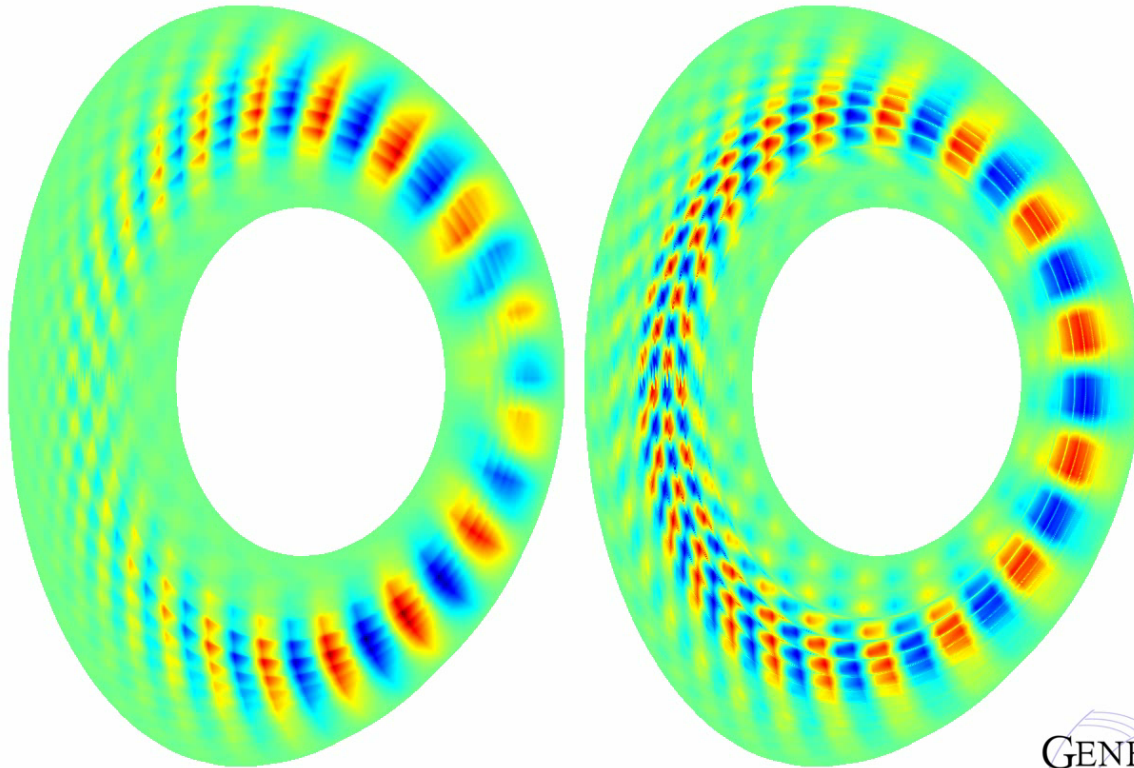


- Fields within $\sim 1-3$ poloidal turns

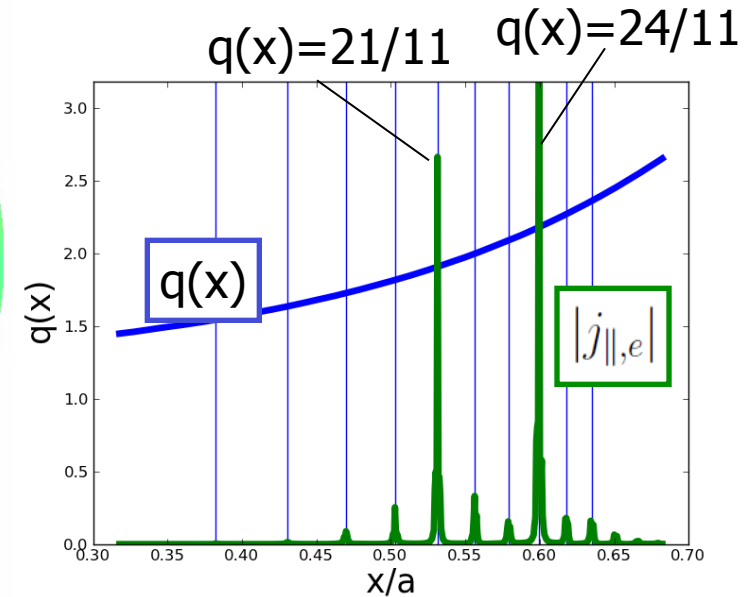
- ~ 60 poloidal turns in ϕ (not A_{\parallel})

Microtearing modes: intrinsic multiscale features in x-y

Global linear simulations of ASDEX Upgrade shot 26459

 \tilde{A}_{\parallel}
 $\tilde{\phi}$


Thanks to
 M. Dunne and D. Told (equilibrium)
 P. Schneider and M. Dunne (T_e, T_i, n)



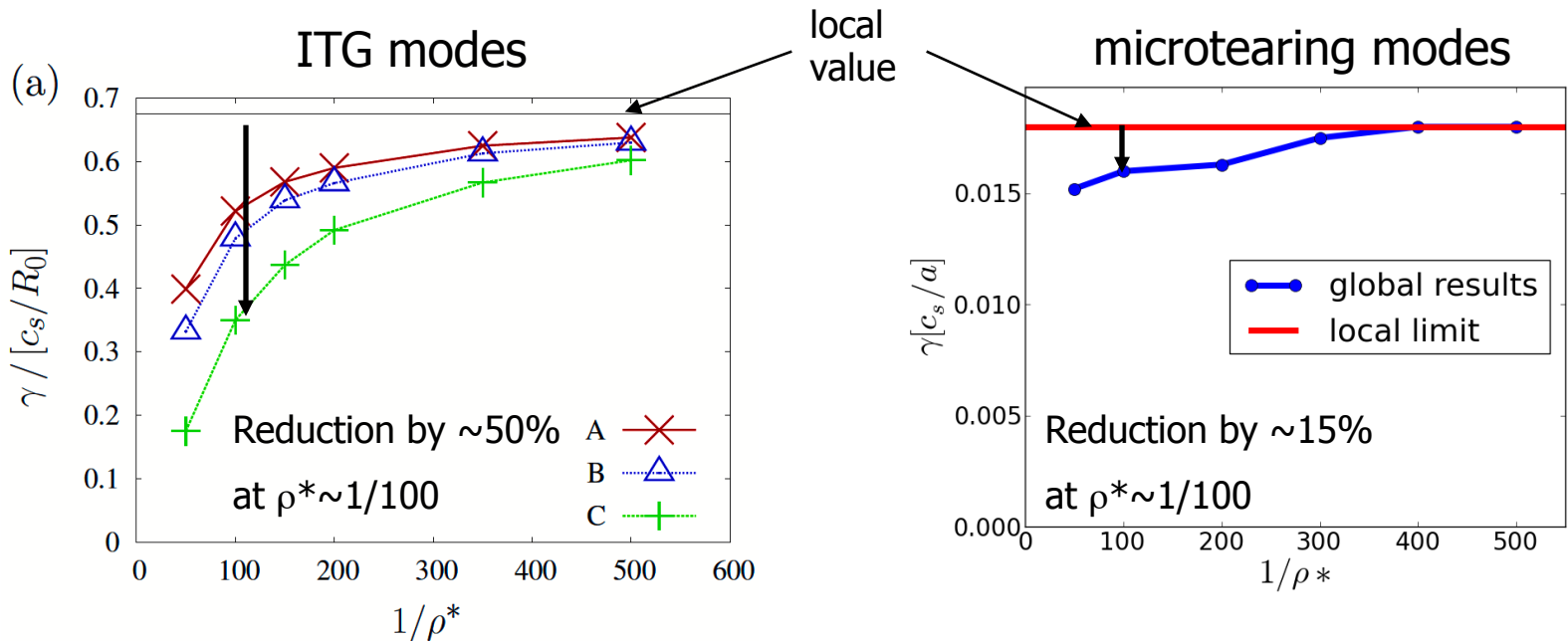
Radial grid spacing: $0.2\rho_i$

Microtearing instabilities expected to exist in AUG

System size effects

ρ^* scan ($\rho^* = \rho_i/a$):

- Even for $\rho^* \sim 1/100$, the local results are fairly accurate
- $k_y \rho_i = 0.12$ is $n=1$ for $\rho^* = 1/50$
- More toroidal mode numbers are microtearing unstable in larger tokamaks

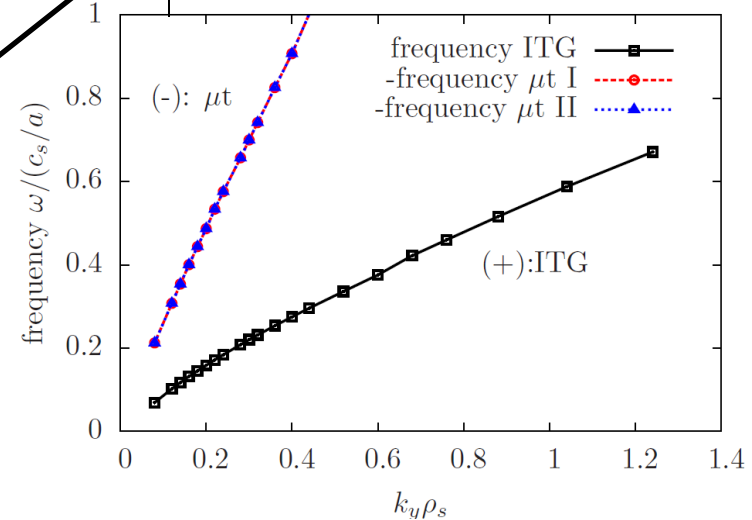
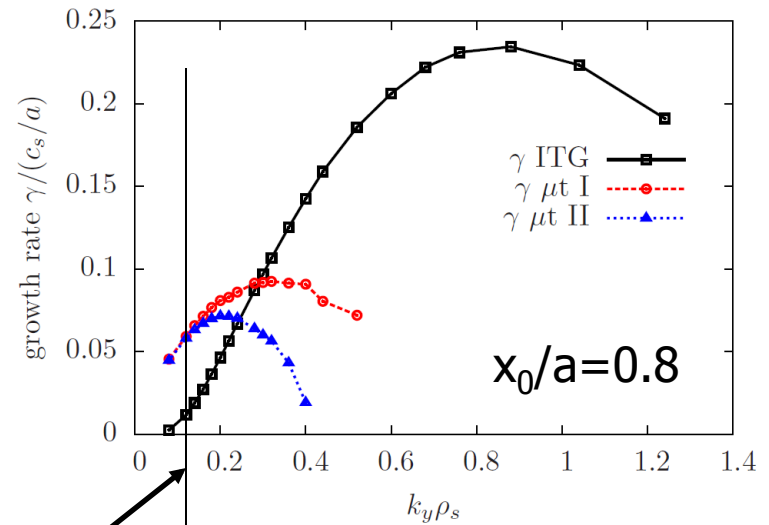
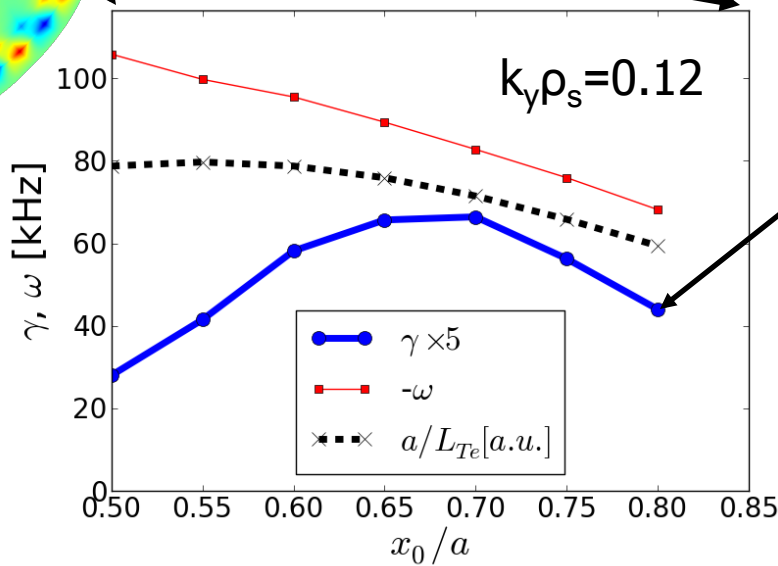
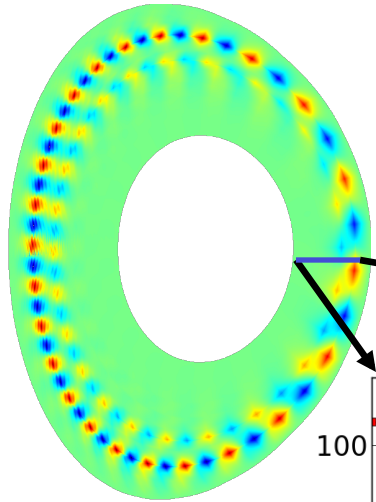


Linear modes are well described in the local approximation

Growth rate spectra involving microtearing modes

ASDEX Upgrade shot 26459

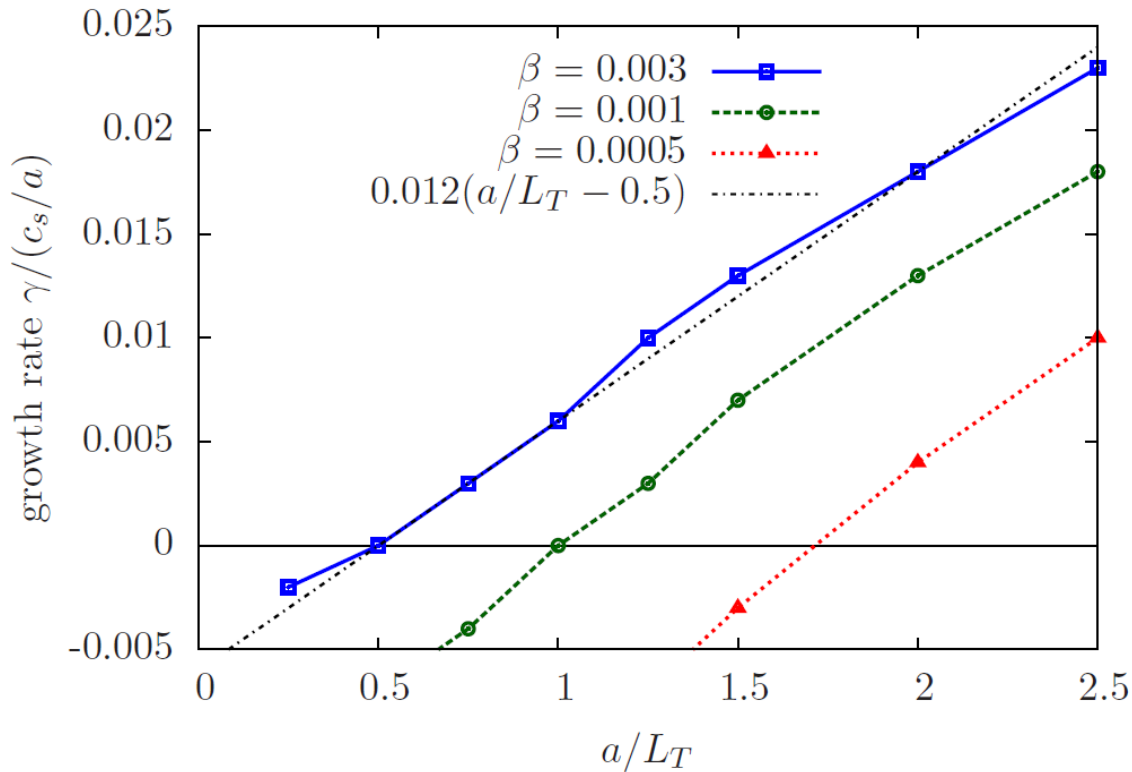
- Lower wavenumber part is dominated by **microtearing**
- Frequency jump at transition to **ITG**



Microtearing modes peak at low k_y , coexist with ITG

Influence of electron temperature gradient and β_e

$$k_y \rho_{\text{ref}} = 0.04$$



Experimental values:

- $\rho_{\text{tor}} = 0.8$
- $a/L_{Te} \sim 2.2$
- $\beta_e \sim 0.003$

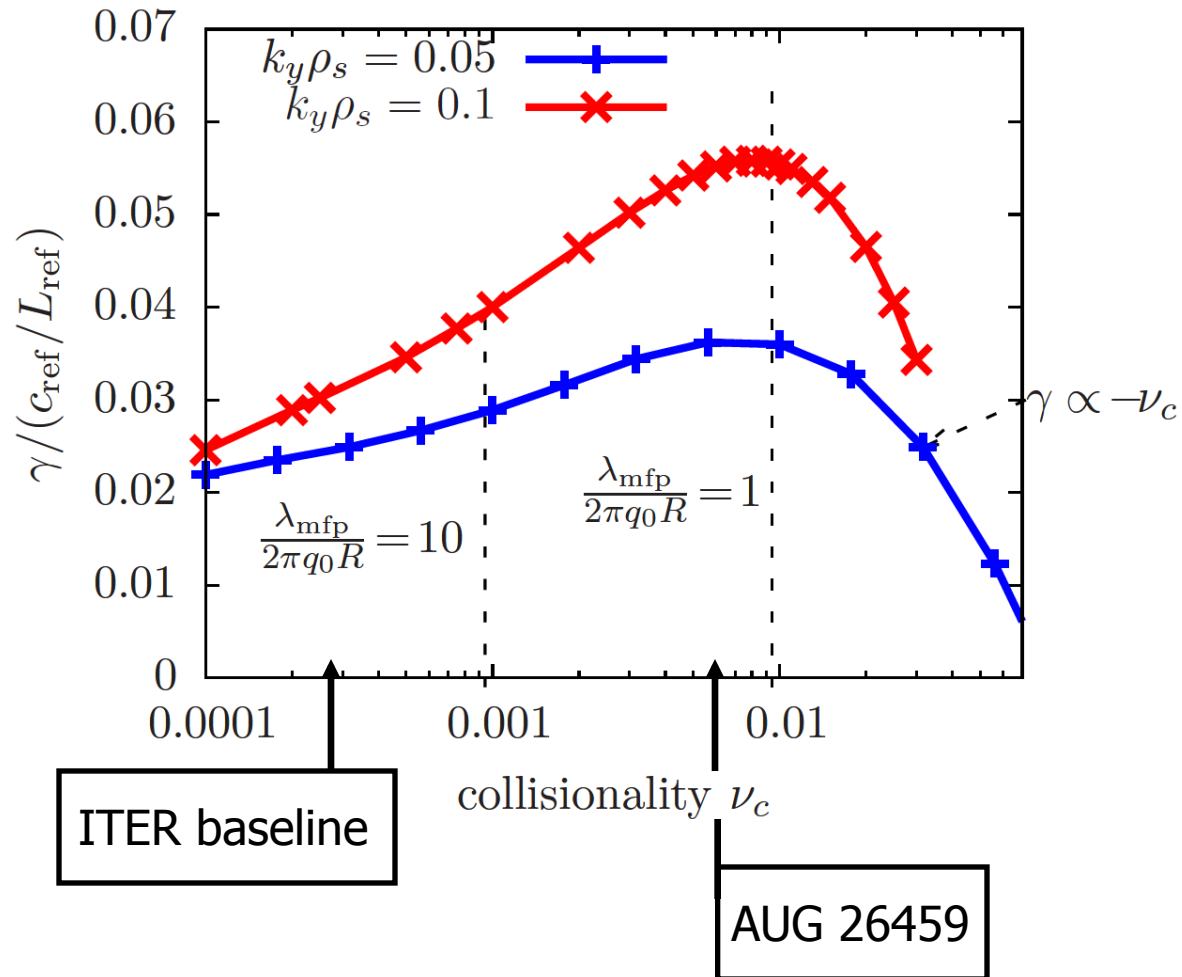
(well above threshold)

a/L_{Te} and β_e are critical plasma parameters

Collisional effects (a/L_{Te} and β kept constant)

- Including **collisions** is **important**
- Growth rate depends on collision frequency only **moderately** (in agreement, e.g., with Applegate '07)
- Experiments are in the **semicollisional** to **collisionless** regime

$$\nu_c \sim n T^{-3/2} \text{ but:}$$



Microtearing modes can also be present in hot (core) plasmas

Thin current layers at resonant flux surfaces

Model of Drake (1977):

$k_{\parallel} = 0$ at rational flux surface $x = 0$:
 Nonadiabatic electron response until $k_{\parallel} v_{te} \approx \omega$
 with $k_{\parallel} \approx k_y x / L_s$ and $L_s = qR / \hat{s}$

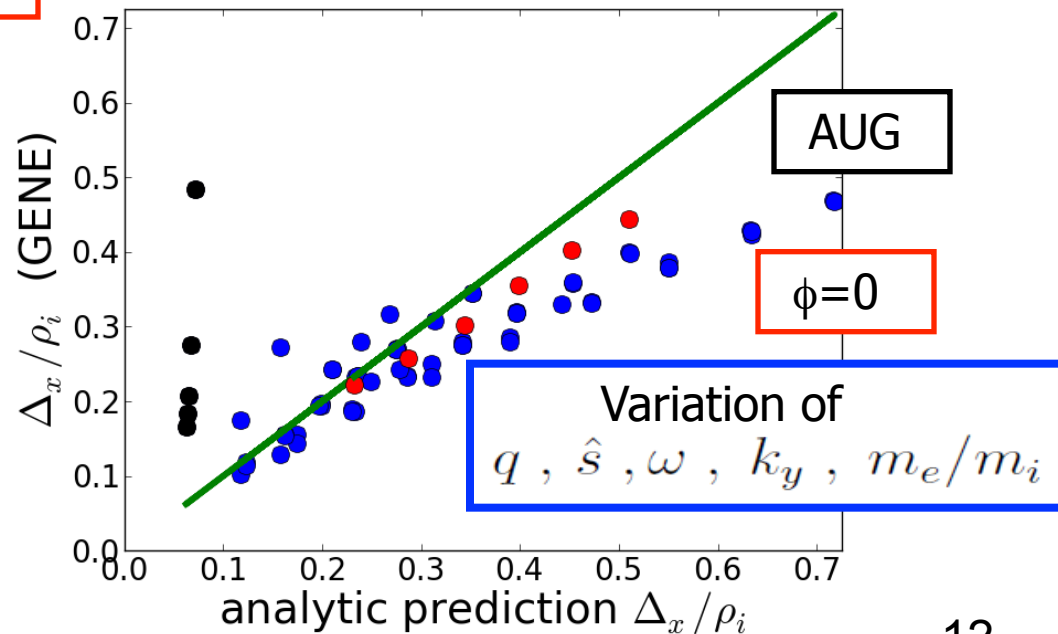
$$\Delta_x / \rho_i = \frac{q}{\hat{s}} \frac{k_y \rho_s}{\omega / (c_s / R)} \sqrt{\frac{m_e}{m_i}}$$


Current layers are thin:

$$\Delta_x = 0.1 - 0.5 \rho_i$$

(requires proper treatment)

GENE simulation results





Nonlinear dynamics: Microtearing turbulence

Turbulent electron heat flux spectrum

Simulation setup:

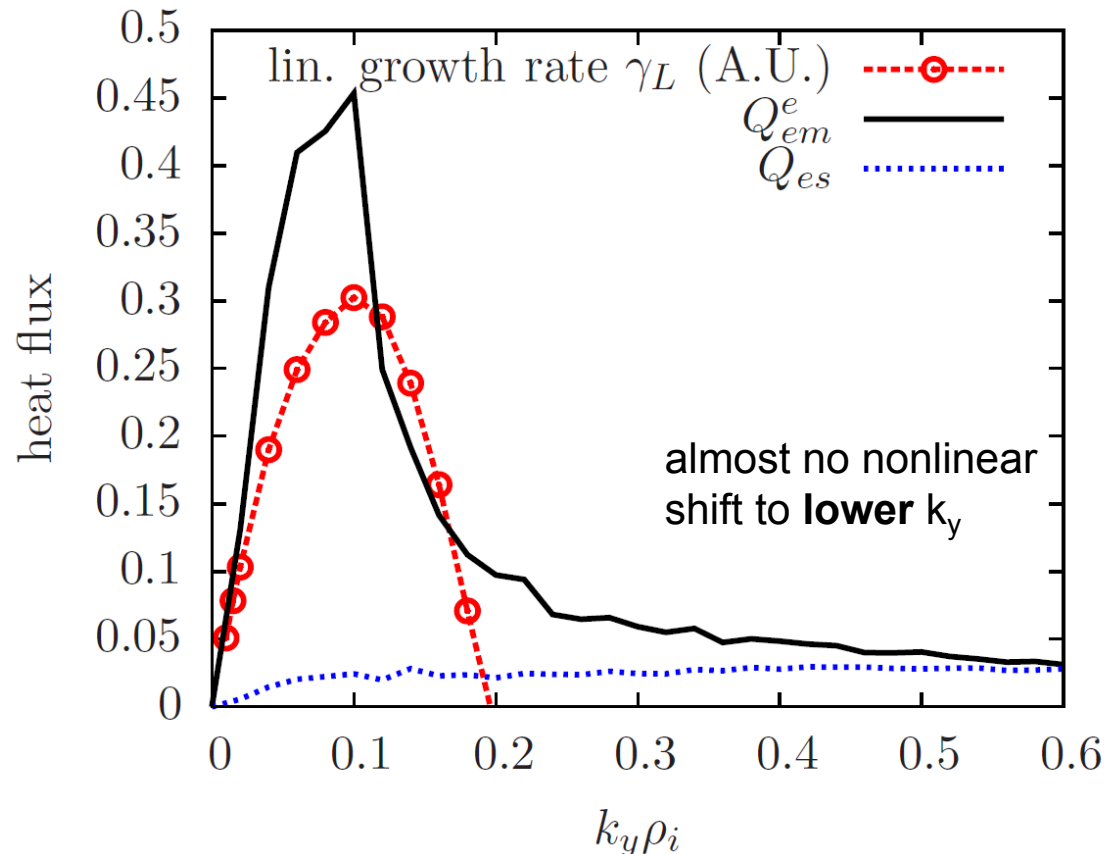
- **sufficiently large radial box**
- **high radial resolution**

384x64x24x32x16 grid points
in 5D ($x, y, z, v_{||}, \mu$) phase space

Very challenging simulations!

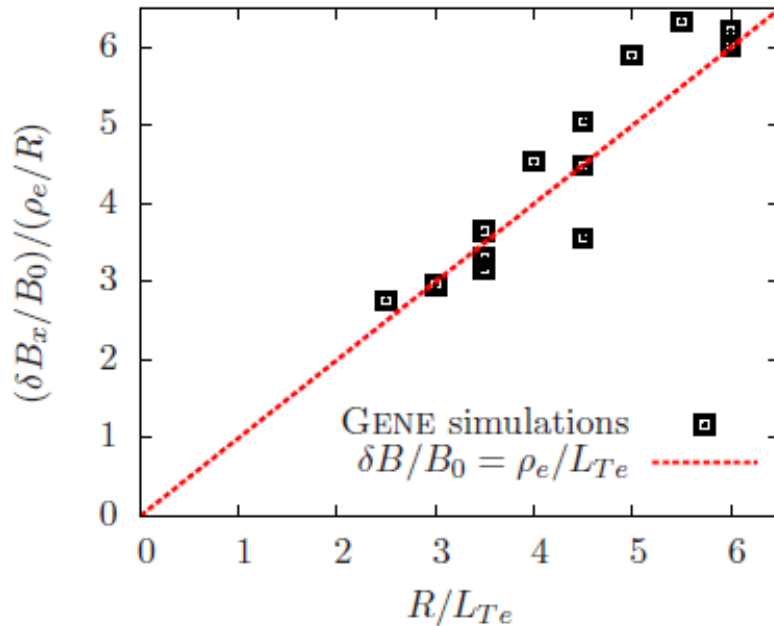
Simplifications:

- **ITG** drive switched **off**
(avoid multimode drive)
- **Circular** magnetic geometry
[Lapillonne et al. 2009]
(allow for maximum flexibility)



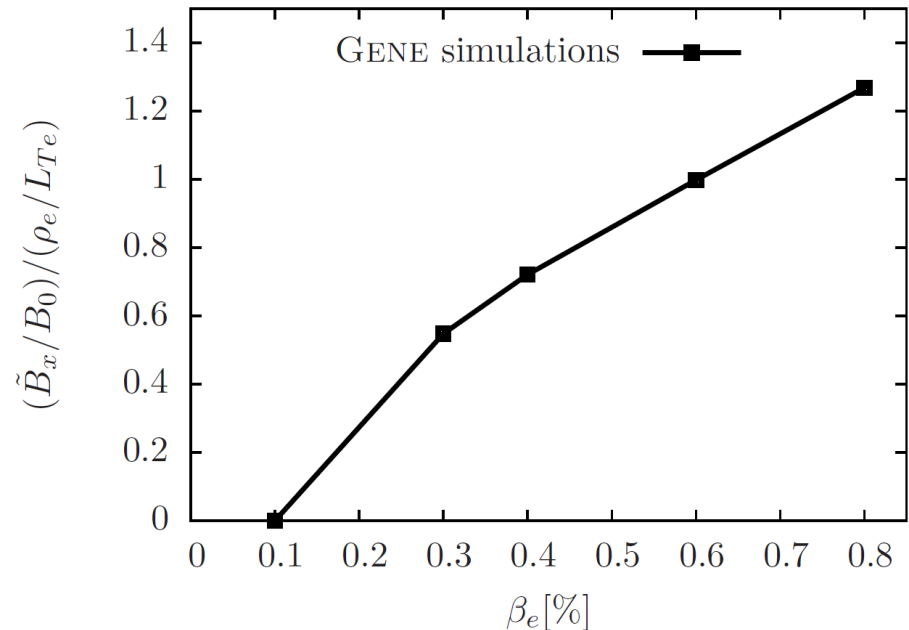
Heat flux is dominated by magnetic component at low k_y

Magnetic fluctuation level in microtearing turbulence



Model by Drake (1980):

$$\tilde{B}/B \sim \rho_e/L_{Te}$$

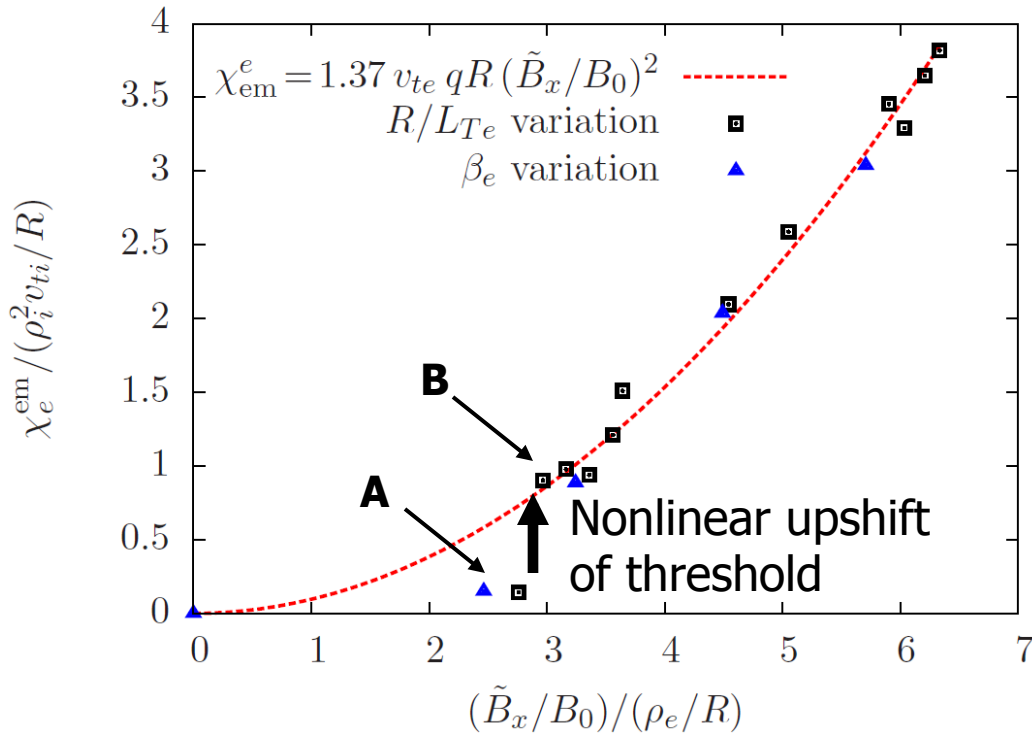


R/L_{Te} scaling

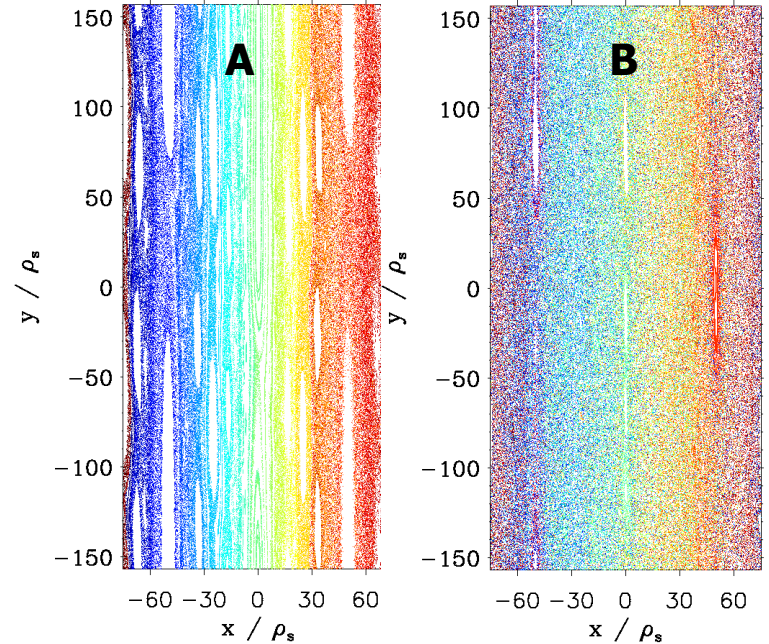
- “scatter”: reduction of the nominal **resolution** by a factor of 2-3 in various phase-space dimensions
- Other parameters like β are important

Drake's formula yields good estimate, but neglects β effects etc.

The role of magnetic stochasticity



Poincare plot of magnetic field lines



Diffusivity model:

$$\chi_e^{em} = v_{te} D_M$$

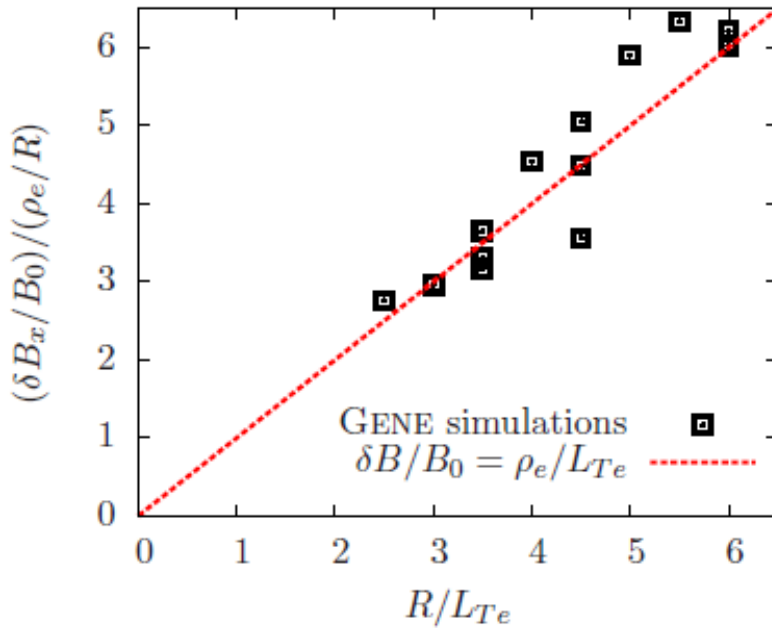
$$D_M = L_C (\delta B / B_0)^2 \text{ (quasilinear result)}$$

use $L_C = qR$ as correlation length

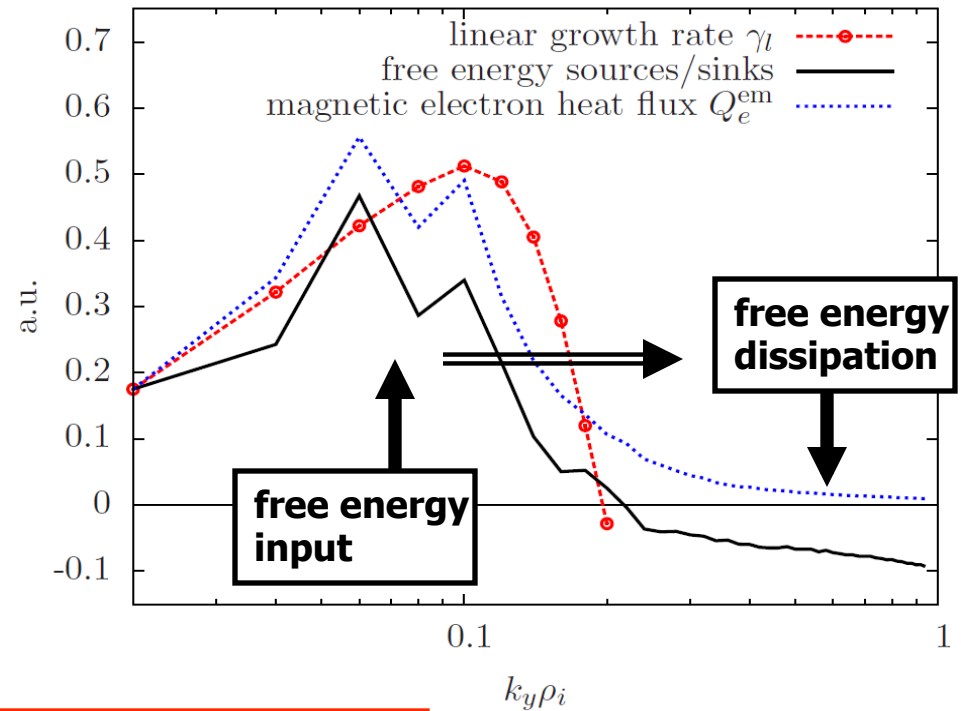
$$\chi_{em}^e = 1.37 q R v_{te} (\tilde{B}_x / B_0)^2$$

- **Effective threshold** for onset of strong transport (**A** to **B**)
- In stochastic cases (**B**), GENE results confirm test-particle models (e.g. Liewer 1985, collisionless case)

Nonlinear saturation of microtearing turbulence



Model by Drake (1980): $\tilde{B}/B \sim \rho_e/L_{Te}$
 (based on dissipation at **low** wavenumbers)



In the accessible parameter range:


Linear drive: $\gamma_l^{max} \propto (R/L_{Te})^2 v_{ti}/R$

Free energy sink at $k_{diss} \gtrsim 0.2/\rho_i$

Balance: $\gamma_{nl} = \chi_e^{em} k_{diss}^2$

$\tilde{B}/B \sim \rho_e/L_{Te}$

Free energy transfer to medium k_y , where it is dissipated

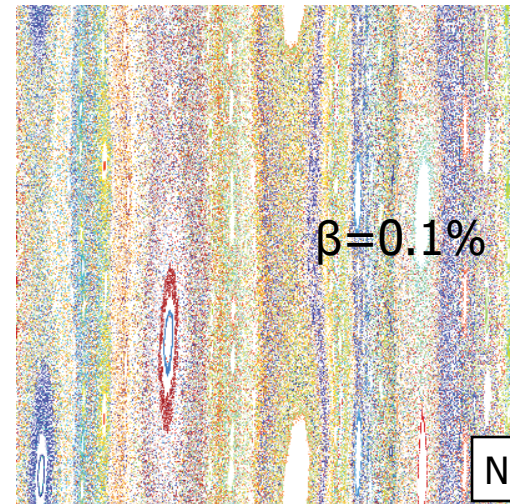
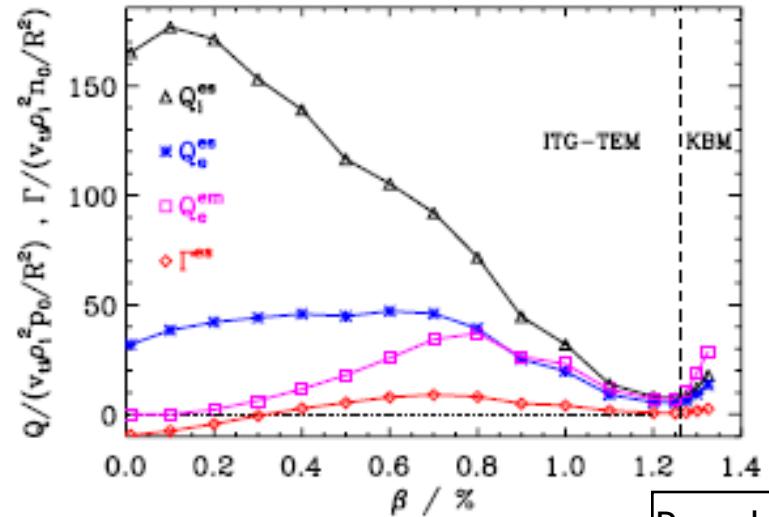
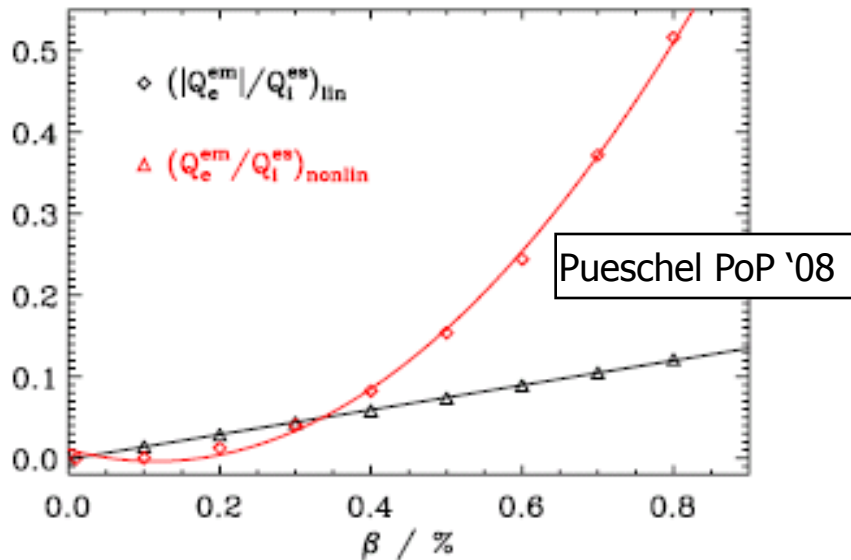


Microtearing modes in finite β ITG/TEM turbulence

Detailed account: D. Hatch (next week)

Motivation: Recent results on finite β turbulence

- Electron **magnetic heat transport** can approach (or even surpass) electrostatic transport as β **increases** [Candy PoP '05, Pueschel PoP '08]
- Magnetic transport violates quasilinear theory – β^2 -**scaling**



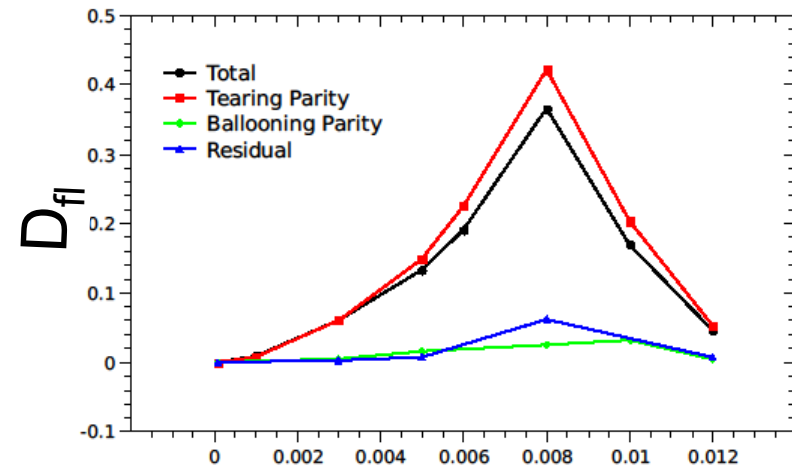
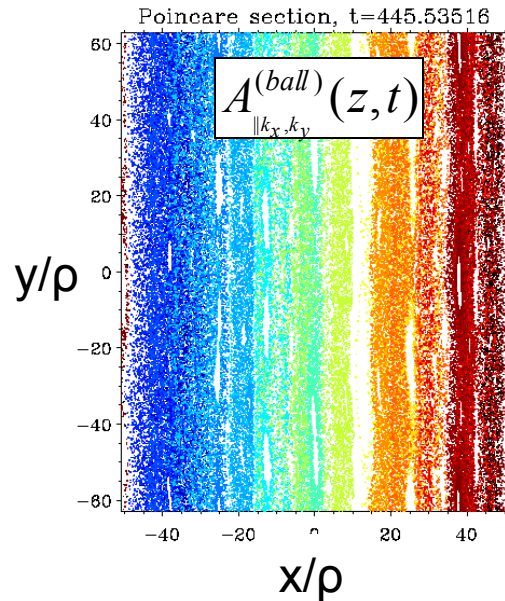
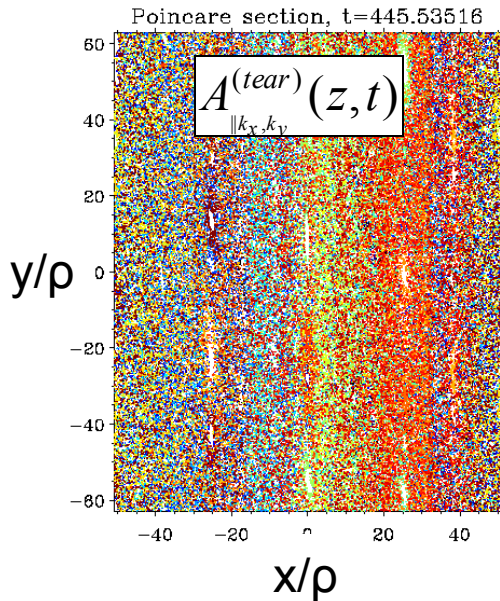
- Observation of near-ubiquitous **magnetic stochasticity** – even at low values of β (Nevins PRL '11, Wang PoP '11)

...needs an explanation!

What is the reason for magnetic stochasticity?

- **Ballooning parity** modes: no reconnection/stochastic fields
- **Tearing parity** modes allow reconnection
- Analysis: **POD** of vector potential (optimal basis) : $A_{\parallel k}(z, t) = \sum_n A_{\parallel k}^{(n)}(z) h_k^{(n)}(t)$
- First two modes (ballooning and tearing) plus residual modes

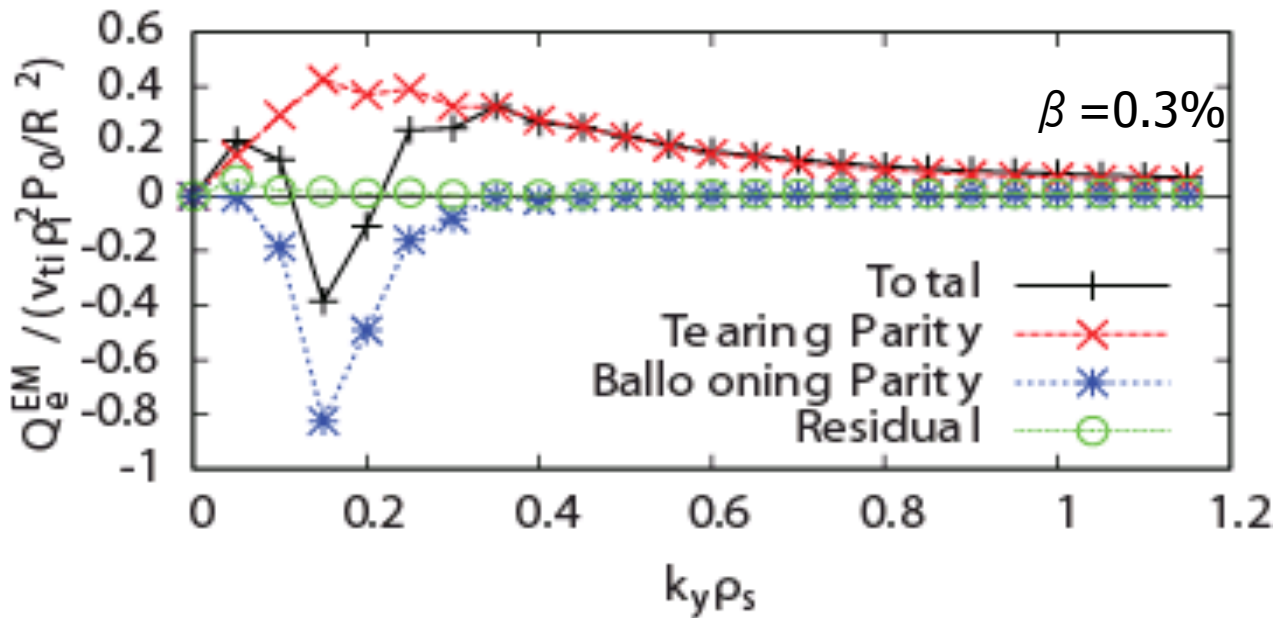
$$A_{\parallel k_x, k_y}(z, t) = A_{\parallel k_x, k_y}^{(ball)}(z, t) + A_{\parallel k_x, k_y}^{(tear)}(z, t) + A_{\parallel k_x, k_y}^{(res)}(z, t)$$



$$D_{fl} = \lim_{l \rightarrow \infty} \frac{1}{l} \left\langle \left| r_i(l) - r_i(0) \right|^2 \right\rangle^\beta$$

Tearing component causes stochasticity

Magnetic transport – superposition of ITG and tearing



$$Q_e^{EM(tot)} = q_{\parallel} B_x$$

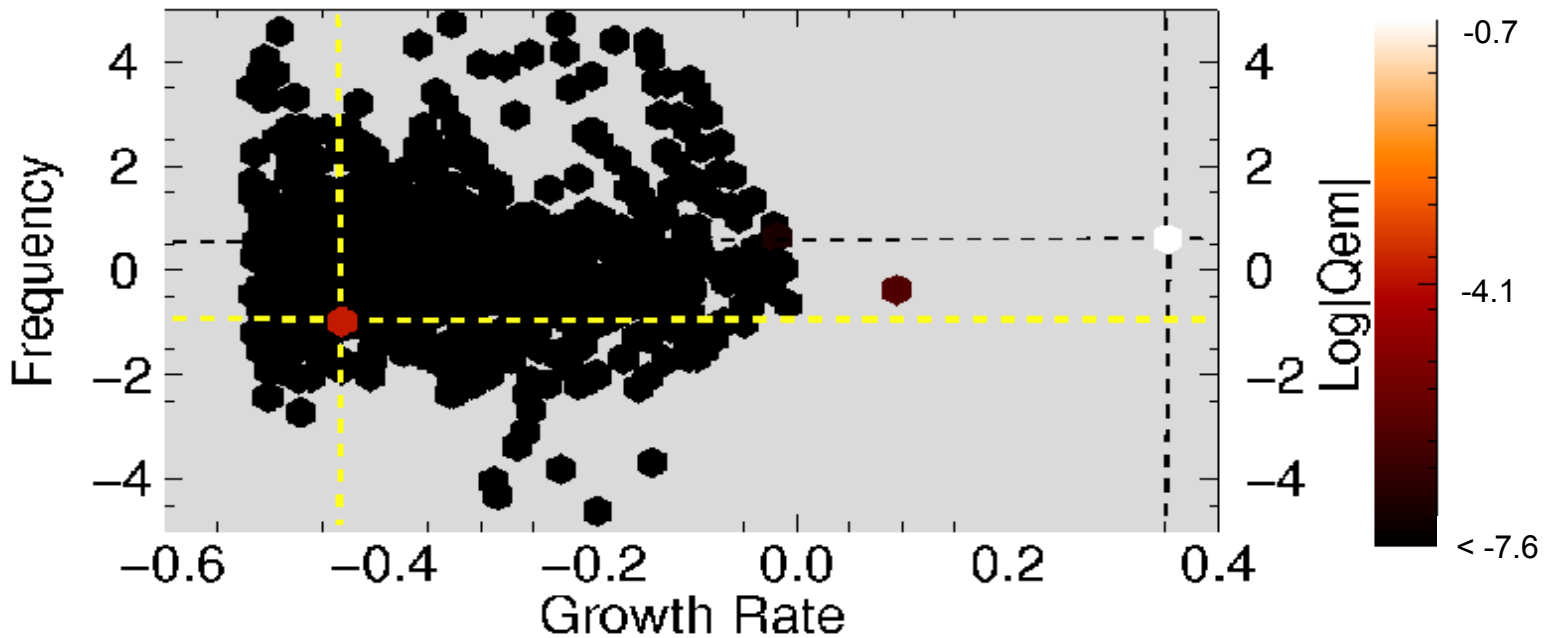
$$Q_e^{EM(tear)} = q_{\parallel} B_x^{(tear)}$$

$$Q_e^{EM(ball)} = q_{\parallel} B_x^{(ball)}$$

$$Q_e^{EM(res)} = q_{\parallel} B_x^{(res)}$$

- **'Dip'** in the magnetic **electron heat transport spectrum** (at k_y of electrostatic transport peak) [Candy, PoP'05]
- **Understood** by attributing transport to **ballooning** and **tearing** part of \mathbf{B}_x
- Stochastic transport **explains** $Q_e^{EM} \sim \beta^2$

Origin of the magnetic electron heat transport



- GENE eigenmode solver finds **all gyrokinetic eigenmodes**
- Nonlinearly evolved **distribution function projected onto 1000 orthogonalized linear eigenmodes** ($k_y \rho_i = 0.2, k_x \rho_i = 0$)
- **One damped eigenmode dominates**, has properties of **microtearing mode**
- Analysis of nonlinear transfer: **Excitation via coupling to zonal modes**

Nonlinearly excited microtearing mode causes electromagnetic heat flux



Summary & Outlook

Summary & Outlook

- **Microtearing modes** are expected to exist in **ASDEX Upgrade** and presumably also in other medium-aspect-ratio tokamaks like **ITER**
- Nonlinear gyrokinetic simulations establish microtearing as a **candidate** to explain **electron heat flux** in outer core of tokamaks: $\chi_{em}^e \sim 1 \text{m}^2/\text{s}$
- **Magnetic stochasticity** occurs when magnetic fluctuation amplitude exceeds an **effective threshold (in analogy to Dimits shift)**
- **Nonlinearly excited (linearly stable) microtearing modes** cause stochasticity and magnetic electron heat flux in **ITG turbulence**
- Underway: Microtearing + ITG turbulence (preliminary result: coexistence!)

References:

- Doerk et al., PRL 2011
- Doerk et al., APS / PoP
- Hatch et al., submitted