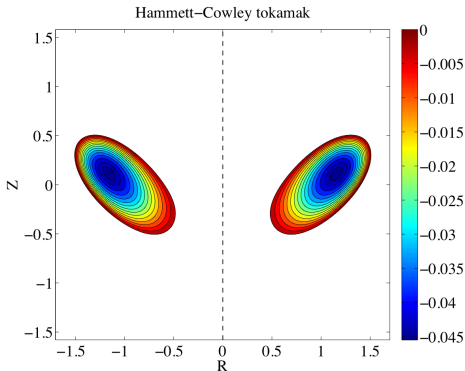


# A fast high-order solver for the Grad-Shafranov equation



Andras Pataki, CIMS

with L. Greengard (CIMS), A. Cerfon (MIT), J. Freidberg (MIT)

# GRAD-SHAFRANOV EQUATION AS POISSON'S EQUATION

$$R \frac{\partial}{\partial R} \left( \frac{1}{R} \frac{\partial \Psi}{\partial R} \right) + \frac{\partial^2 \Psi}{\partial Z^2} = -\mu_0 R^2 \frac{dp}{d\Psi} - F \frac{dF}{d\Psi}$$

$$\Psi = Cst \quad \text{on plasma boundary}$$

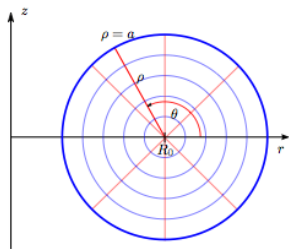
- ▶ Change the unknown function:  $\Psi = \sqrt{RU}$

$$\frac{\partial^2 U}{\partial R^2} + \frac{\partial^2 U}{\partial Z^2} = \frac{3}{4} \frac{U}{R^2} - \mu_0 R \frac{dp}{dU} - \frac{1}{2R} \frac{dF^2}{dU}$$

$$\sqrt{RU} = Cst \quad \text{on plasma boundary}$$

- ▶ The left-hand side is the **2-D Laplacian**
- ▶ We want the first and second derivatives of  $U$  with very good accuracy  $\Rightarrow$  **Spectral methods**

# FAST POISSON SOLVER ON A CIRCLE



$$\begin{aligned}\Delta u(\rho, \theta) &= f(\rho, \theta) \\ u(a, \theta) &= g(\theta)\end{aligned}$$

- ▶ Write data and solution as Fourier series (use FFT)

$$u(\rho, \theta) = \sum \hat{u}_n(\rho)e^{in\theta}, \quad f(\rho, \theta) = \sum \hat{f}_n(\rho)e^{in\theta}, \quad g(\theta) = \sum \hat{g}_n e^{in\theta}$$

- ▶ Plugging into Poisson's equation, we get mode-by-mode ODE

$$\begin{aligned}\hat{u}_n''(\rho) + \frac{1}{\rho}\hat{u}_n'(\rho) - \frac{n^2}{\rho^2}\hat{u}_n(\rho) &= \hat{f}_n(\rho) \\ \hat{u}_n(a) &= \hat{g}_n\end{aligned}$$

## FAST POISSON SOLVER ON A CIRCLE

- ▶ Using Green's functions to construct particular solution (that does not satisfy B.C.)

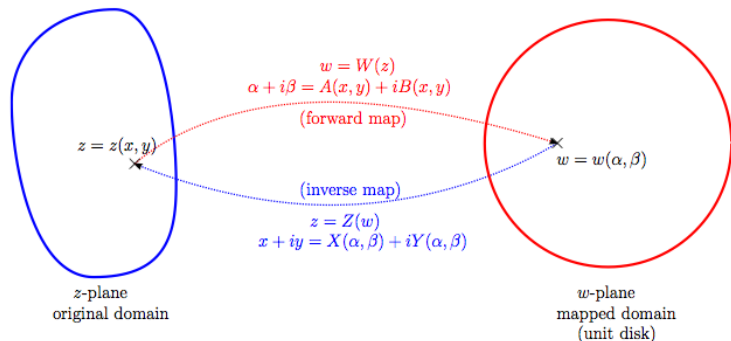
$$\begin{aligned}\hat{u}_n^{part}(\rho) &= \int_0^\infty G_n(\rho, s)f(s)ds \\ &= -\frac{1}{2n} \left[ \rho^{-n} \int_0^\rho s^{n+1}f(s)ds + \rho^n \int_\rho^\infty s^{-n+1}f(s)ds \right]\end{aligned}$$

- ▶ Piecewise Chebyshev / Legendre grid crucial for high order
- ▶ To match the B.C., add homogeneous correction:

$$\hat{u}_n(\rho) = \hat{u}_n^{hom}(\rho) + \hat{u}_n^{part}(\rho) \quad \text{with } \hat{u}_n^{hom}(\rho) = c_n \left(\frac{\rho}{a}\right)^n$$

- ▶ Boundary condition yields **explicit condition**  $c_n = \hat{g}_n - \hat{u}_n^{part}$
- ▶ Major advantage of Green's function method: **differentiation**

# CONFORMAL MAPPING

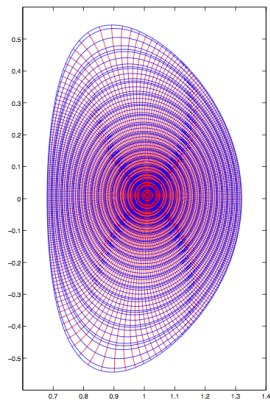


$$\Delta u(x, y) = f(x, y) \quad \rightarrow \quad \Delta v(\alpha, \beta) = f(X(\alpha, \beta), Y(\alpha, \beta)) \cdot \left| \frac{dZ}{dw} \right|^2$$

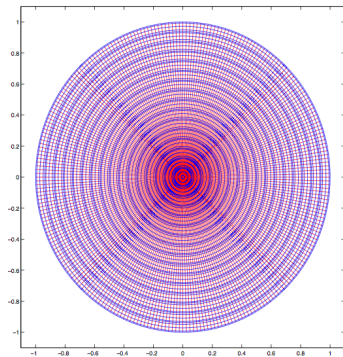
$$u|_{\partial\Omega} = g(x, y) \quad \rightarrow \quad v|_{\partial\Omega} = g(X(\alpha, \beta), Y(\alpha, \beta))$$

- ▶  $W$  computed with the [Kerzman-Stein](#) integral equation
- ▶ Problem with [crowding](#) for large aspect ratio

# CONFORMAL MAPPING

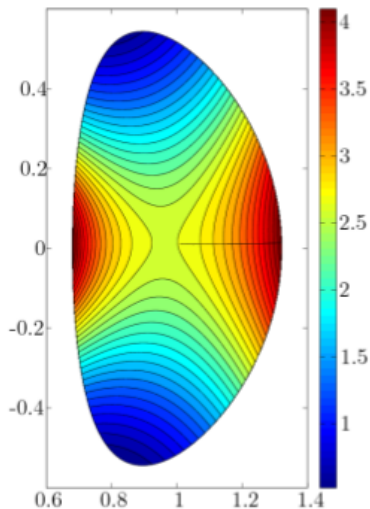


z-plane  
original domain



w-plane  
mapped domain  
(unit circle)

# CONFORMAL MAPPING

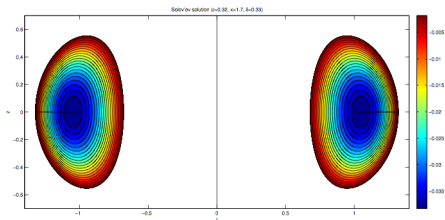


ITER crowding factor: 5

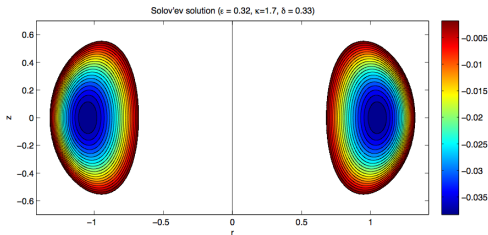
NSTX crowding factor: **20**

# ITER – RESULTS

## Numerical solution



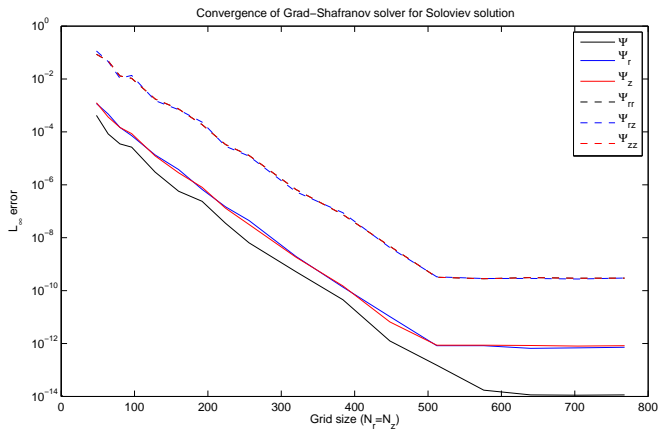
## Exact solution



$$\epsilon = 0.32, \kappa = 1.7, \delta = 0.33$$

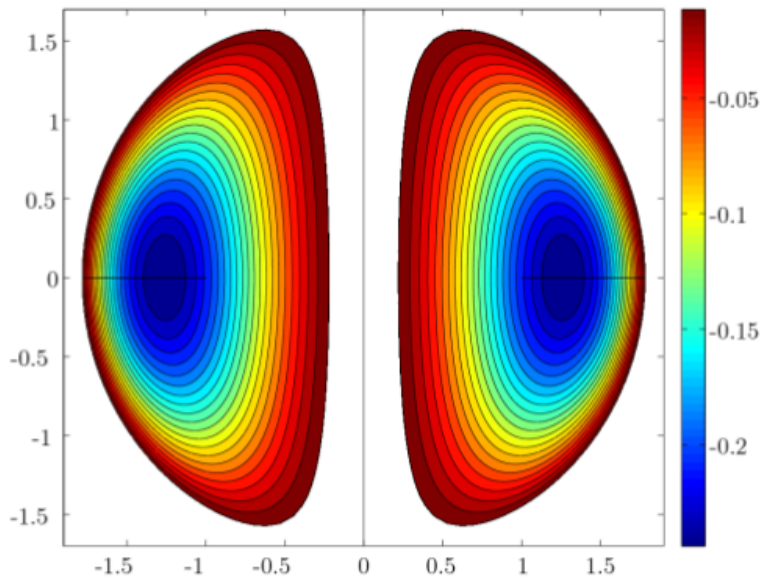


# ITER – CONVERGENCE

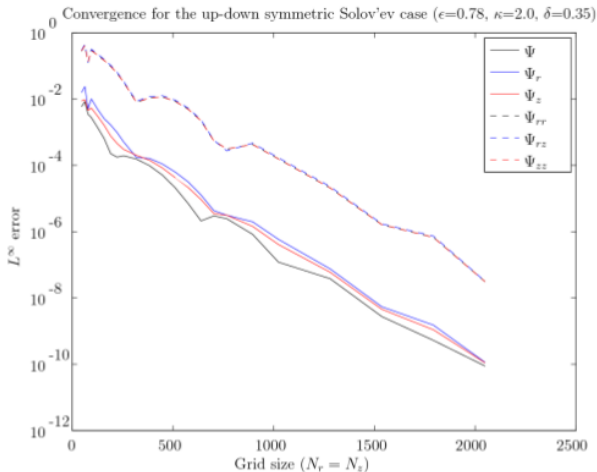


- ▶ Exponential convergence – solution almost to machine accuracy
- ▶ Good accuracy for derivatives (error due to numerical derivative of boundary conditions)

# NSTX – RESULTS



# NSTX – CONVERGENCE



Results are not as good  
for STs

Due to:

- ▶ Oversampling of 3 compared to ITER
- ▶ Domain boundary requires slightly more points to be resolved