Intrinsic rotation in tokamaks

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Can we get large rotational shear for free?

- Recent observations of significant rotation (~Mach 0.3) without external momentum injection
- Important because difficult to force rotation in large, dense plasmas (i.e. ITER)
- Complicated dependence on experimental parameters



Parra et al., PRL (2012).

What sets intrinsic rotation?

• For small rotation and rotation gradient:

$$\Pi = -C\Omega_{\zeta} - \nu\Omega_{\zeta}'$$

• Steady-state with no momentum input ($\Pi = 0$):

$$\Omega_{\zeta} \propto \exp\left(-\int_{0}^{\psi} d\psi' \frac{C(\psi')}{\nu(\psi')}\right)$$

- No rotation sign reversal possible, so can't explain many experimental observations
- Must consider additional physics (not just diffusion + pinch)

Symmetry property



Symmetry property

No momentum transport from local turbulence in up-down symmetric system!



Symmetry of gyrokinetic system

 Lowest-order GK system invariant under updown reflection about equatorial plane, reversal of rotation and rotation gradient, and

$$v_{\parallel} \to -v_{\parallel}, \ \theta \to -\theta, \ \partial_{\psi} \to -\partial_{\psi}$$

$$\delta f \to -\delta f, \ \Phi \to -\Phi$$

 Consequently, momentum flux must vanish for zero rotation and rotation gradient (relevant for intrinsic rotation):

$$\Pi(\Omega_{\zeta}, \Omega_{\zeta}') = -\Pi(-\Omega_{\zeta}, -\Omega_{\zeta}') \Rightarrow \Pi(0, 0) = 0$$

Parra, Barnes, and Peeters, Phys. Plasmas (2011).

Symmetry in GS2 simulations

 $\mathsf{P}(heta,k_\psi,k_lpha,v_\parallel,\mu)$ is the time-averaged momentum flux integrand

Symmetry constraints:

 $\mathsf{P}(v_{\parallel},\theta) = -\mathsf{P}(-v_{\parallel},-\theta) \qquad \mathsf{P}(k_{\psi},k_{\alpha}) = -\mathsf{P}(-k_{\psi},k_{\alpha})$





$$\gamma_E = -0.2 (v_{
m th}/R)$$



Symmetry breaking: mean flow



Mean flow generation



Mean flow generation



Mean flow generation



How big is the effect?

$$u \propto \Delta r \frac{d \ln p}{dr} v_t \sim \frac{\rho_{\theta}}{L_p} v_t$$

$$\Pi = -mRn\nu \frac{du}{dr}, \ Q = -n\chi \frac{dT}{dr}$$

$$\nu \sim \chi \sim q R \frac{\rho_{\theta}^2}{v_t}$$

$$\frac{\Pi}{Q} \frac{Q_{GB}}{\Pi_{GB}} \sim m v_t \frac{du}{dr} \left(\frac{dT}{dr}\right)^{-1} \sim \rho_i \left(\frac{dT}{dr}\right)^{-1} \frac{d}{dr} \left(\frac{q}{\epsilon} \frac{dT}{dr}\right)$$

Symmetry breaking: up-down asymmetry



Symmetry breaking: orbit length



Symmetry breaking: orbit length



Symmetry breaking: radial profile variation



Higher-order GK equation (B_{θ} /B small)

(R,E,µ) variables

Parra, Barnes, and Catto, Nucl. Fusion (2011).

$$\begin{split} \frac{dg_s}{dt} + \mathbf{v}_{\parallel} \cdot \nabla \left(g_s - Z_s e \left\langle \Phi \right\rangle \frac{\partial F_{0s}}{\partial E} \right) + \left\langle \mathbf{v}_E^{\perp} \right\rangle \cdot \nabla F_{0s} \\ &+ \left(\mathbf{v}_{Cs} + \mathbf{v}_{Ms} + \left\langle \mathbf{v}_E^{\perp} \right\rangle \right) \cdot \nabla_{\perp} \left(g_s - Z_s e \left\langle \Phi \right\rangle \frac{\partial F_{0s}}{\partial E} \right) \\ &= - \mathbf{v}_{Ms} \cdot \nabla \theta \frac{\partial}{\partial \theta} \left(g_s - Z_s e \left\langle \Phi \right\rangle \frac{\partial F_{0s}}{\partial E} \right) - \left\langle \mathbf{v}_E^{\parallel} \right\rangle \cdot \nabla_{\perp} g_s \\ &- \left\langle \mathbf{v}_E^{\perp} \right\rangle \cdot \nabla \theta \frac{\partial g_s}{\partial \theta} - \left\langle \mathbf{v}_E^{\parallel} \right\rangle \cdot \nabla F_{0s} - \left\langle \mathbf{v}_E^{\perp} \right\rangle \cdot \nabla F_{1s} \\ &+ Z_s e \left(\mathbf{v}_{\parallel} \cdot \nabla \left\langle \Phi \right\rangle + \mathbf{v}_{Ms} \cdot \nabla_{\perp} \left\langle \Phi \right\rangle \right) \left(\frac{\partial g_s}{\partial E} + \frac{\partial F_{1s}}{\partial E} \right) \\ &- \mathbf{v}_E^{nc} \cdot \nabla_{\perp} g_s + Z_s e \mathbf{v}_{\parallel} \cdot \nabla \Phi^{nc} \frac{\partial g_s}{\partial E} + \psi \text{-profile variation} \end{split}$$

Symmetry breaking



Symmetry breaking



Symmetry breaking

