

Gyrokinetic Simulations of Solar Wind Turbulence and its Dissipation: Importance of Nonlocal Effects on the Energy Cascade

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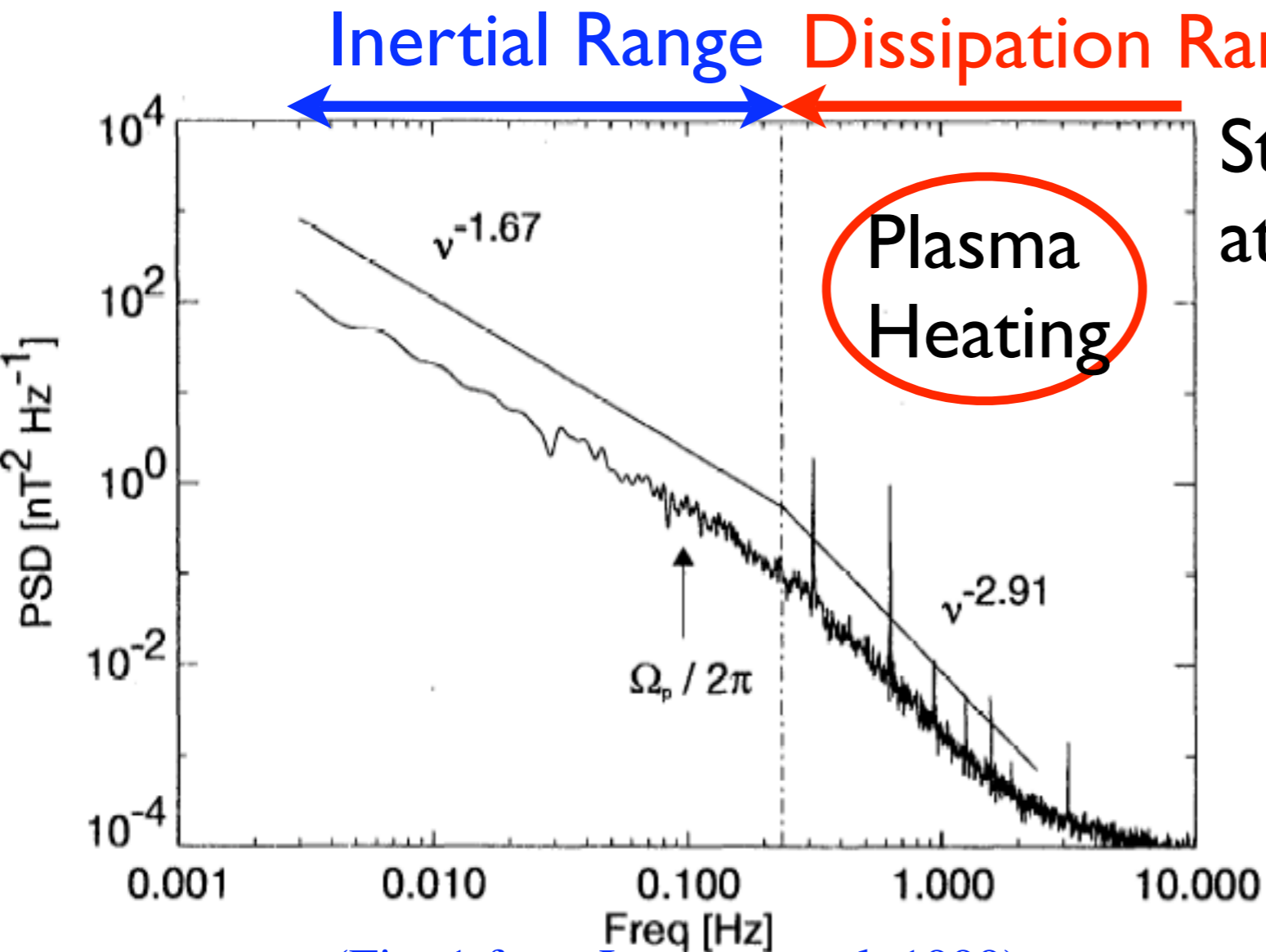
The Center for Multi-scale Plasma Dynamics

Outline

- Turbulent Spectra in the Dissipation Range of the Solar Wind
 - Questions about the Dissipation Range
- Theoretical Models of Kinetic Plasma Turbulence
- Gyrokinetic Simulations of Kinetic Turbulence
- Improved Cascade Model and Implications
- Answers to Questions about the Dissipation Range
- Conclusions

Solar Wind Magnetic Energy Spectrum

Early Observations



Steeper Dissipation Range attributed to:

- Proton Cyclotron Damping
(Goldstein, Roberts, & Fitch 1994, Leamon et al. 1998b, Gary 1999)
- Kinetic Alfvén Wave Damping
(Leamon et al. 1998a)
- Whistler Dispersion
(Stawicki, Gary, & Li 2001)

Physics underlying the dissipation range
is not well understood!

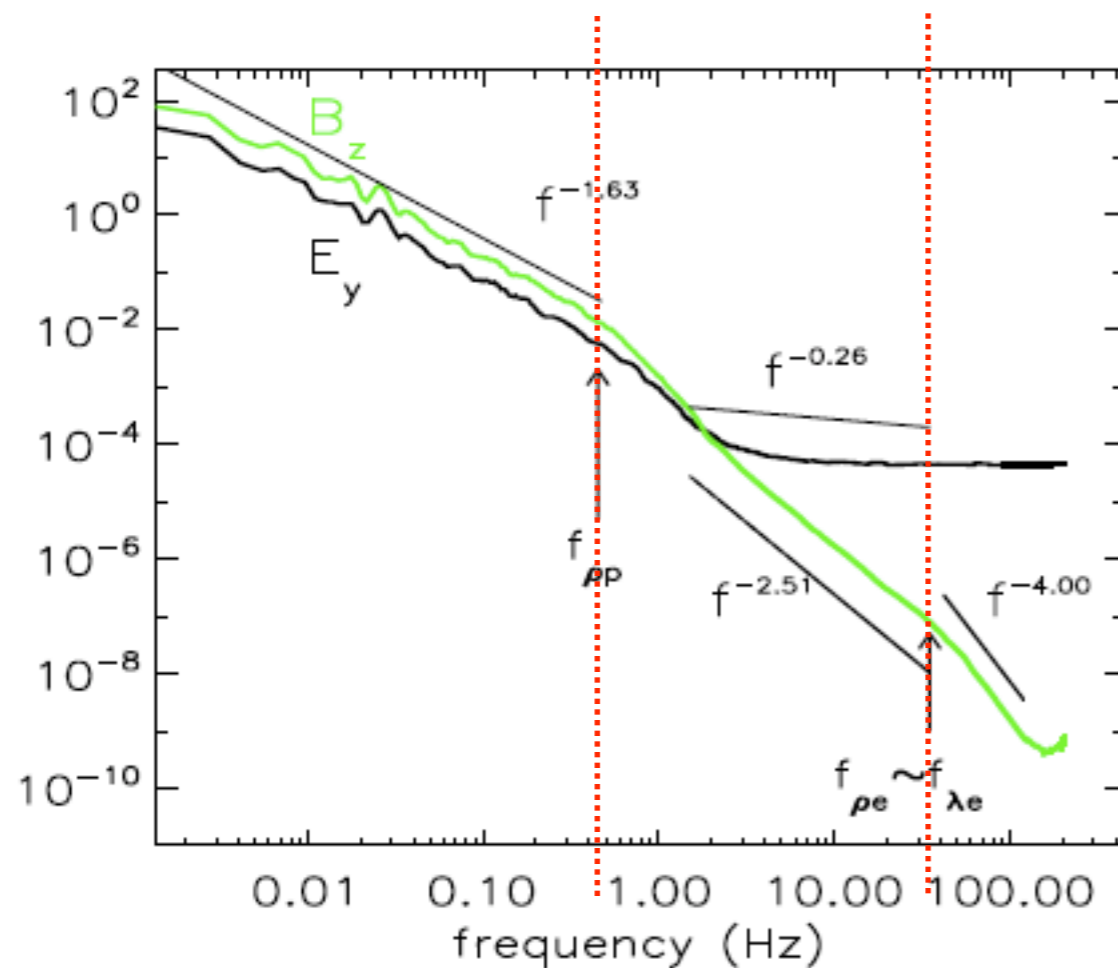
Turbulent plasma heating depends critically on this physics!

Dissipation Range Spectra

Recent Observations

Observations find nearly power-law behavior down to electron scales!

How do we interpret these observations?



1) Can theoretical models explain the observations?

2) What are the wave modes that comprise the dissipation range?
KAW or whistler?

3) Is this a **dissipation range** or a **dispersion range**?

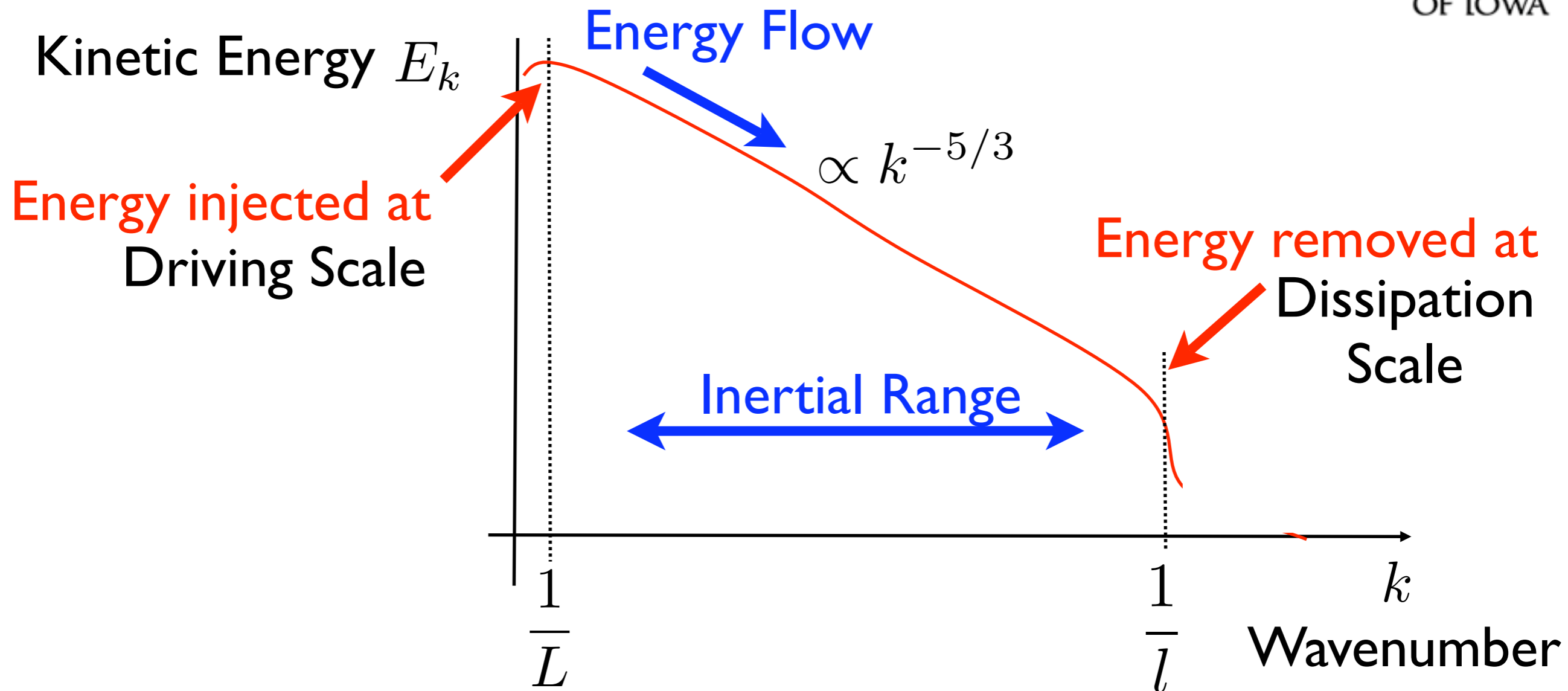
(Sahraoui, Goldstein, Robert, Khotyaintsev 2009, PRL)

(see also Kiyani et al. 2009, PRL and
Alexandrova et al. 2009 PRL)

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Kolmogorov Hydrodynamic Turbulence



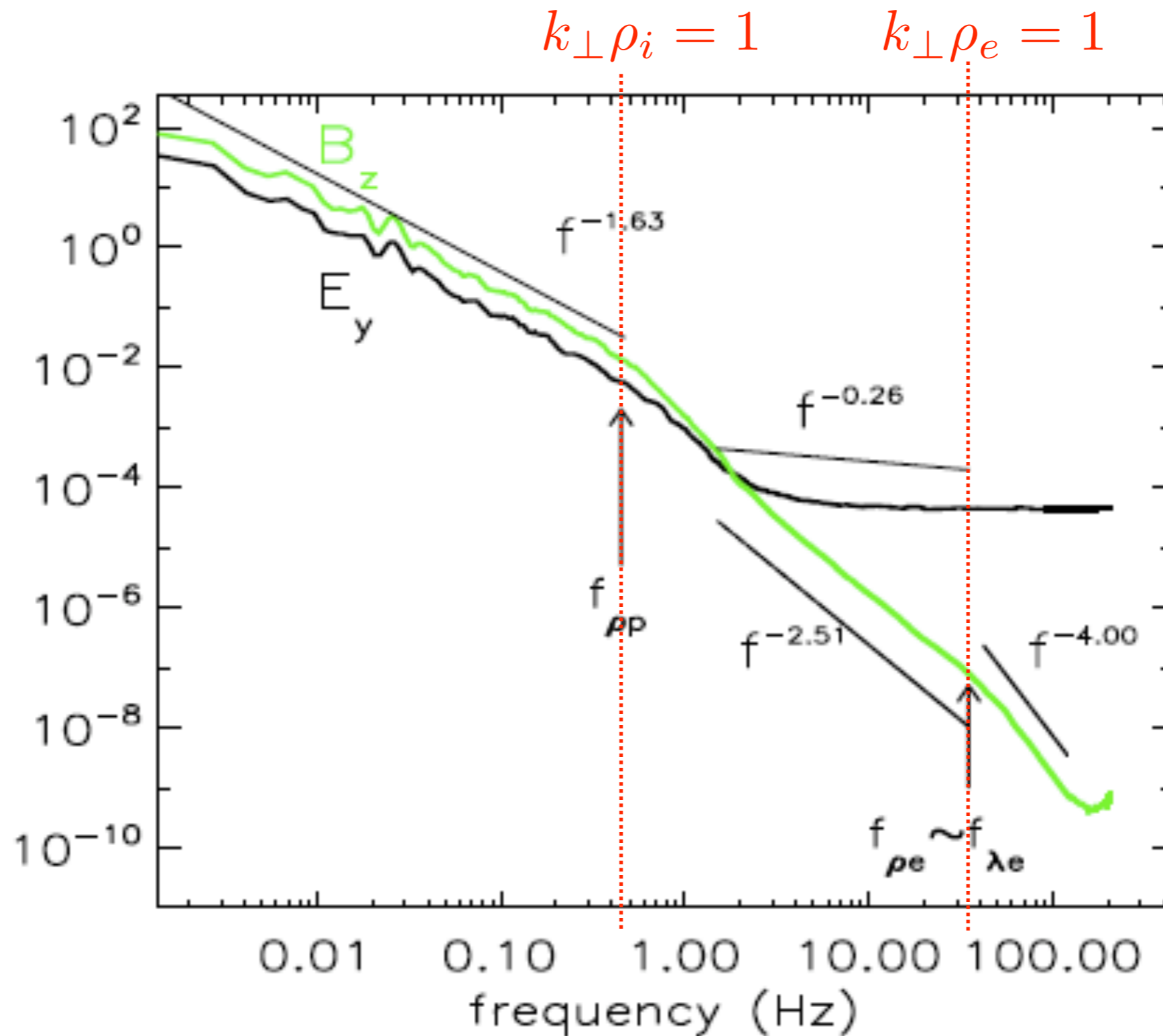
- Astrophysical turbulence develops an **Inertial Range**
- **Kolmogorov Hypothesis:** (Kolmogorov, 1941)
 - Energy transfer occurs locally in wavenumber
 - Energy cascade rate in inertial range is constant

Modifications for Kinetic Turbulence

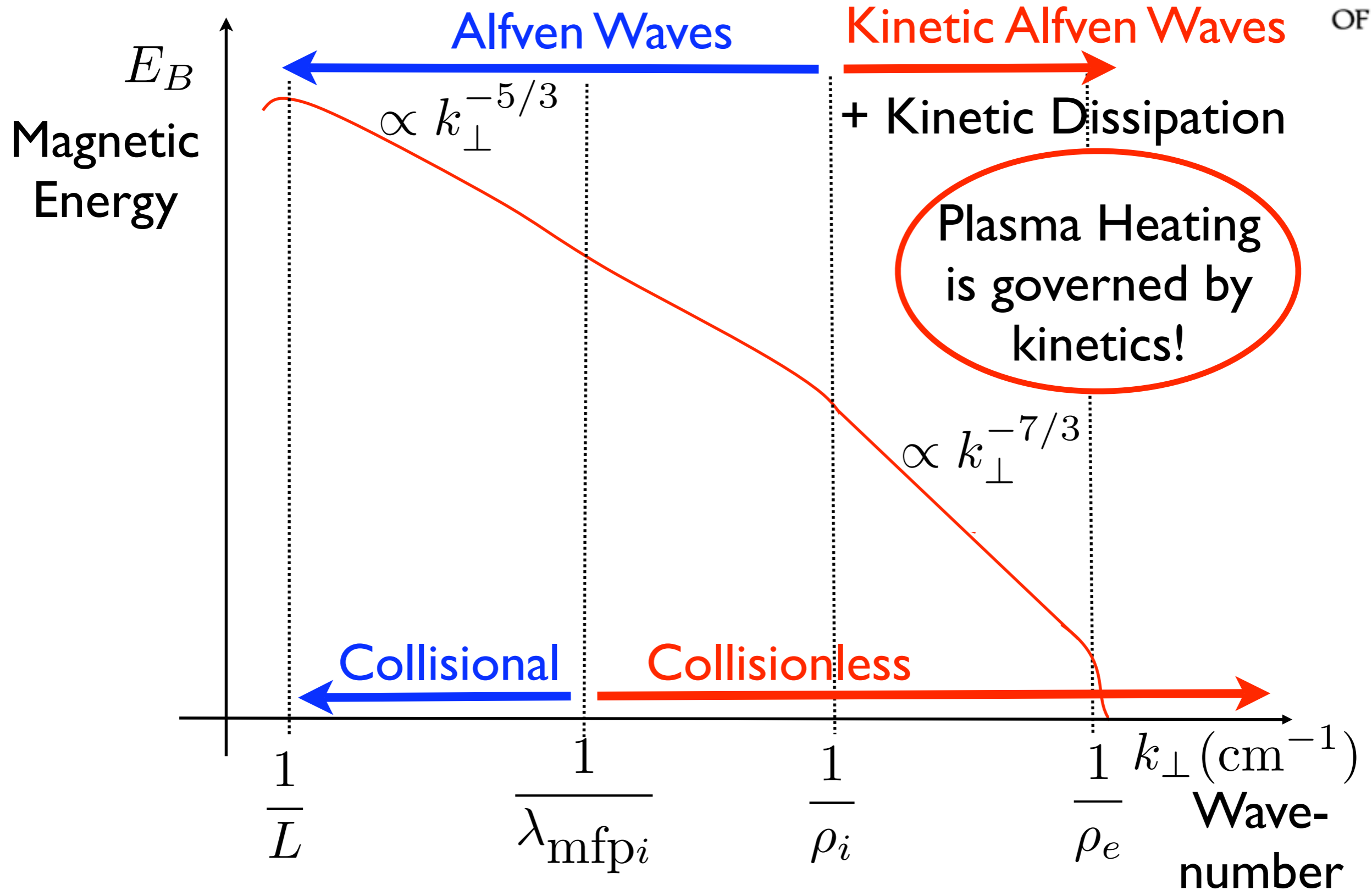
Turbulence Theory in Kinetic Plasmas must incorporate:

- 1) Inherent **anisotropy** of MHD turbulence
- 2) Weak collisionality at small scales
 - a) Transition to **Kinetic Alfvén Waves** at $k_{\perp} \rho_i \sim 1$
 - b) **Damping via Kinetic Mechanisms**, e.g. Landau damping

Solar Wind Observations



Model of Kinetic Turbulent Cascade



(Schekochihin, Cowley, Dorland, Hammett, Howes, Quataert, & Tatsuno, ApJS 182, 310, 2009)

Cascade Model for Kinetic Turbulence

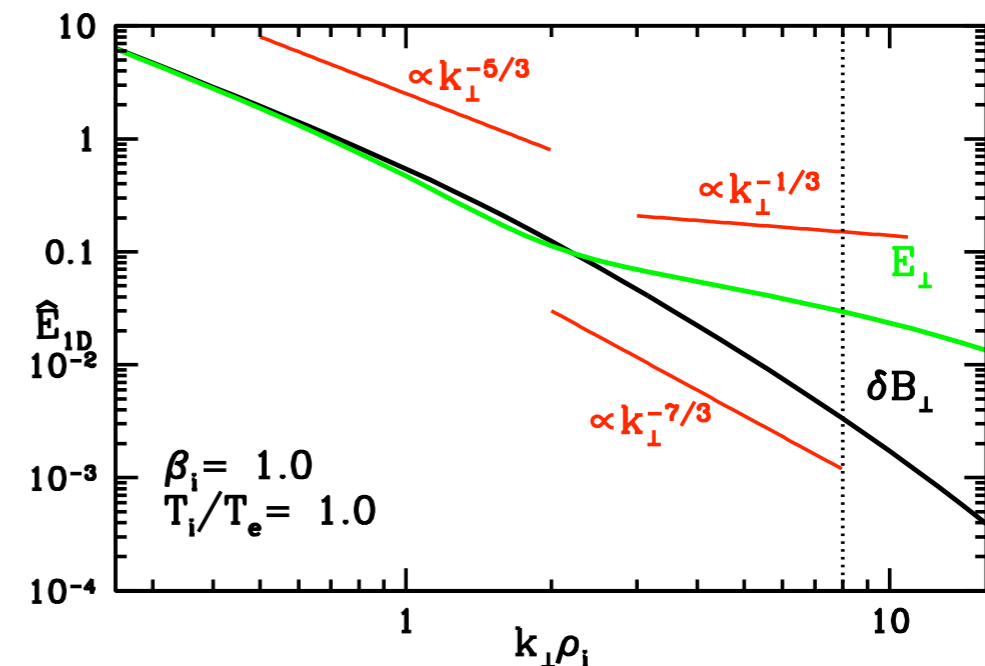
- Cascade Model based on three assumptions: (Howes et al., 2008b)
 1. **Kolmogorov Hypothesis**: Spectrally local nonlinear transfer
 2. **Critical Balance** of linear and nonlinear times
 3. Applicability of **linear kinetic damping** rates

$$\frac{\partial b_k^2}{\partial t} = - \underbrace{\frac{\partial}{\partial \ln k_{\perp}} \left(\frac{b_k^2 \omega_{nl}(k_{\perp})}{C_1 C_2} \right)}_{\text{Nonlinear Transfer}} + \underbrace{\mathcal{S}(k_{\perp})}_{\text{Source}} - \underbrace{2\gamma(k_{\perp}) b_k^2}_{\text{Dissipation}}$$

Predicted Magnetic Energy Spectrum



Source Dissipation

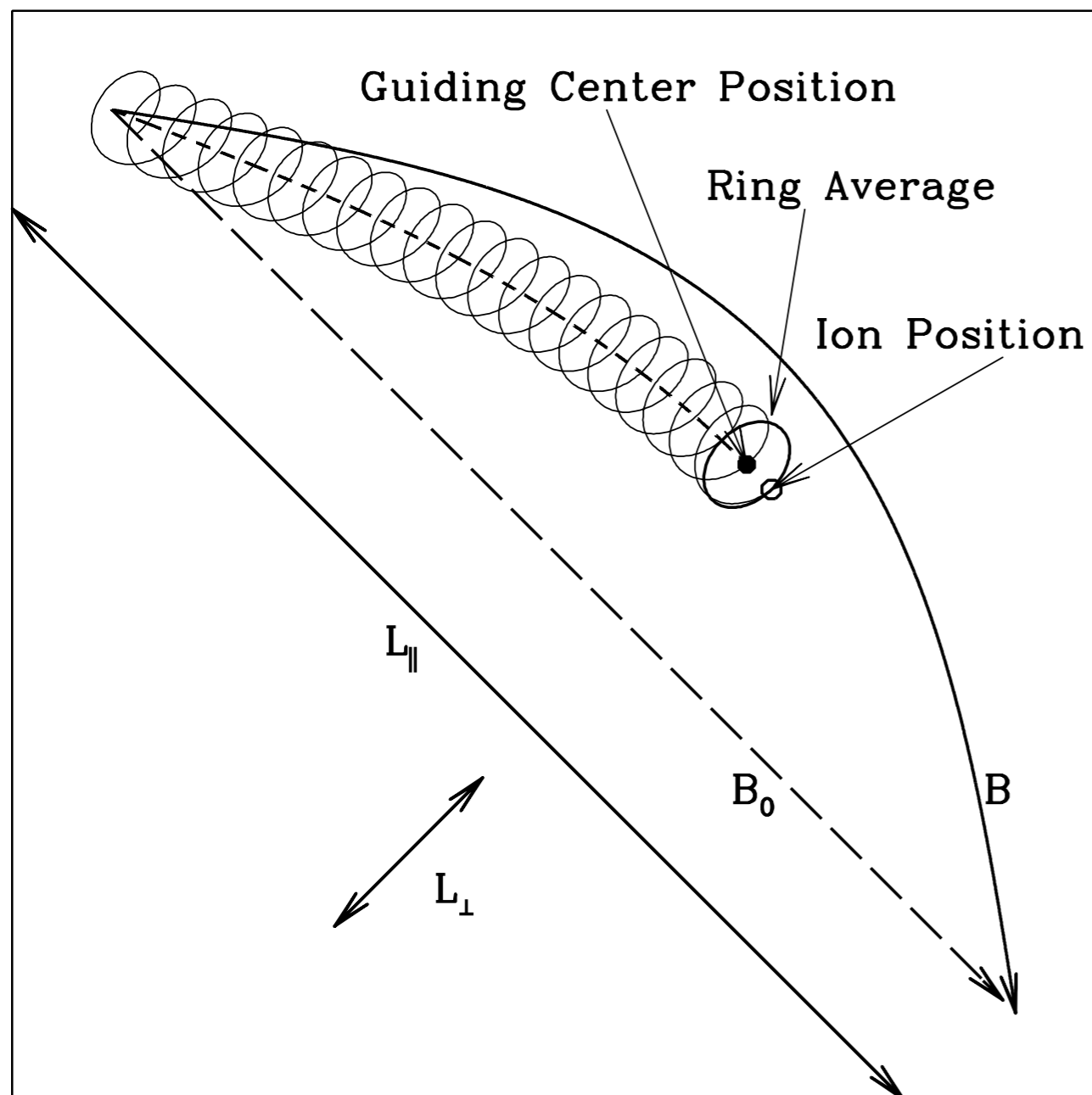


(see also Podesta, Borovsky, and Gary 2010)

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Gyrokinetics is kinetic theory averaged over the Larmor motion. (Rutherford & Frieman 1968; Taylor & Hastie 1968; Frieman & Chen 1982; Howes et al. 2006)



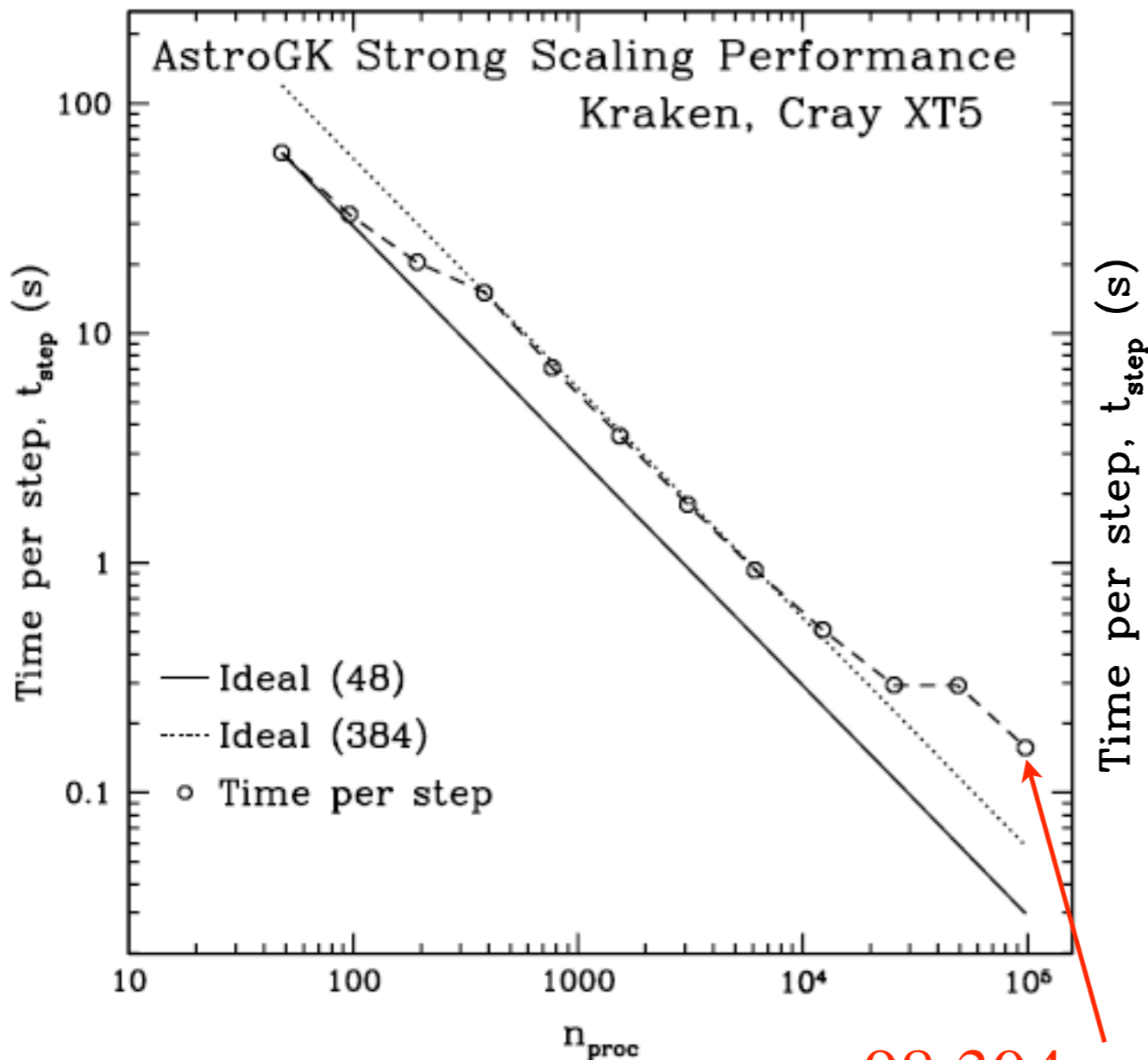
- Low-frequency limit eliminates fast cyclotron timescale $\omega \ll \Omega_i$
- Anisotropic $k_{||} \ll k_{\perp}$
- **Captures:** Finite Larmor radius, Landau resonance, and Collisions
- **Excludes:** Fast wave and cyclotron resonance

AstroGK Simulations

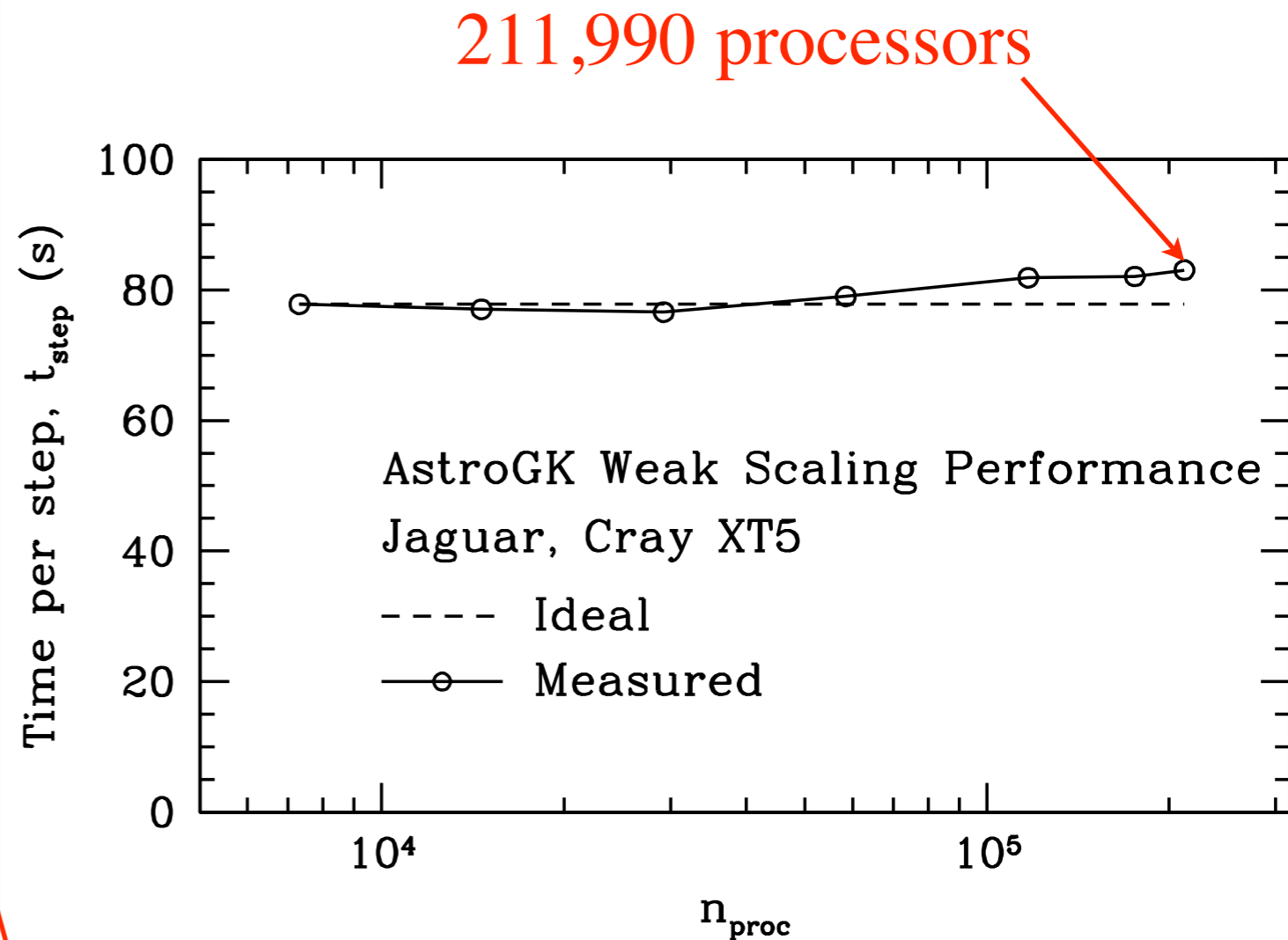
- 5-D Distribution Function, two species, fully electromagnetic
- Ion to Electron scale simulations require millions of CPU hours

The Astrophysical Gyrokinetics Code

(Numata, Howes, Tatsuno, Barnes, & Dorland, J. Comp. Phys., submitted 2010)



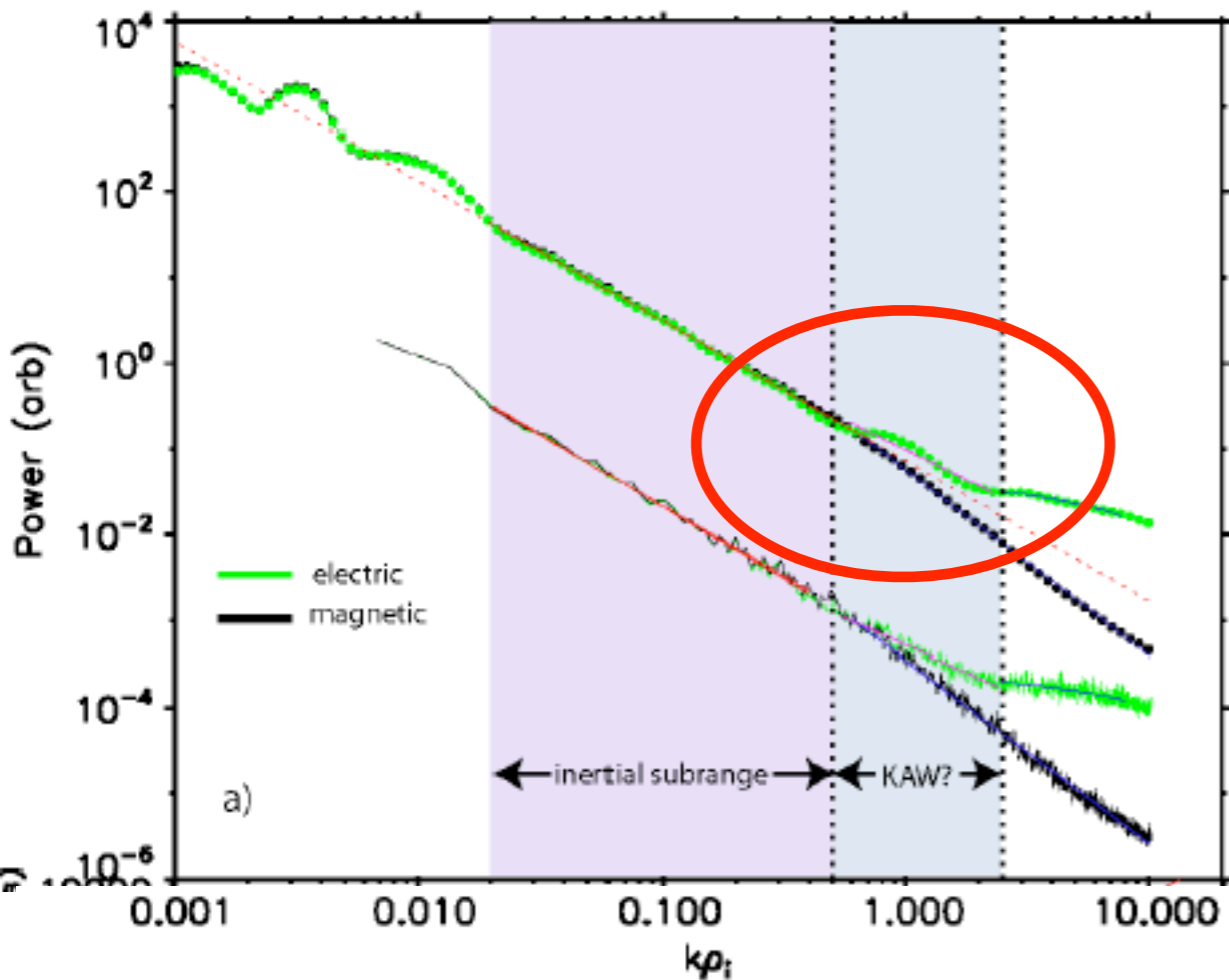
98,304 processors



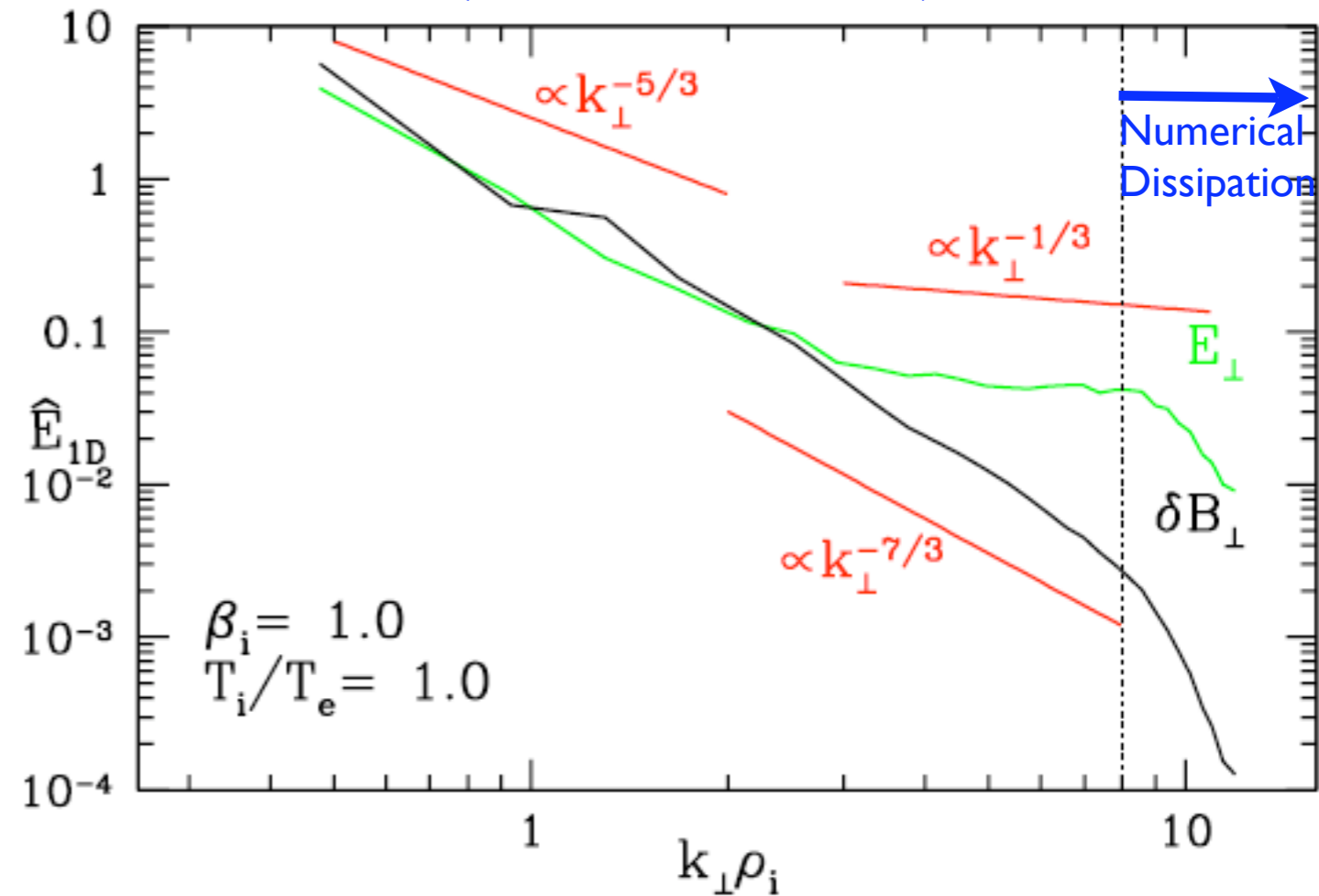
211,990 processors

Transition to Kinetic Alfven Wave Turbulence

(Fig. 3 from Bale et al. 2005)



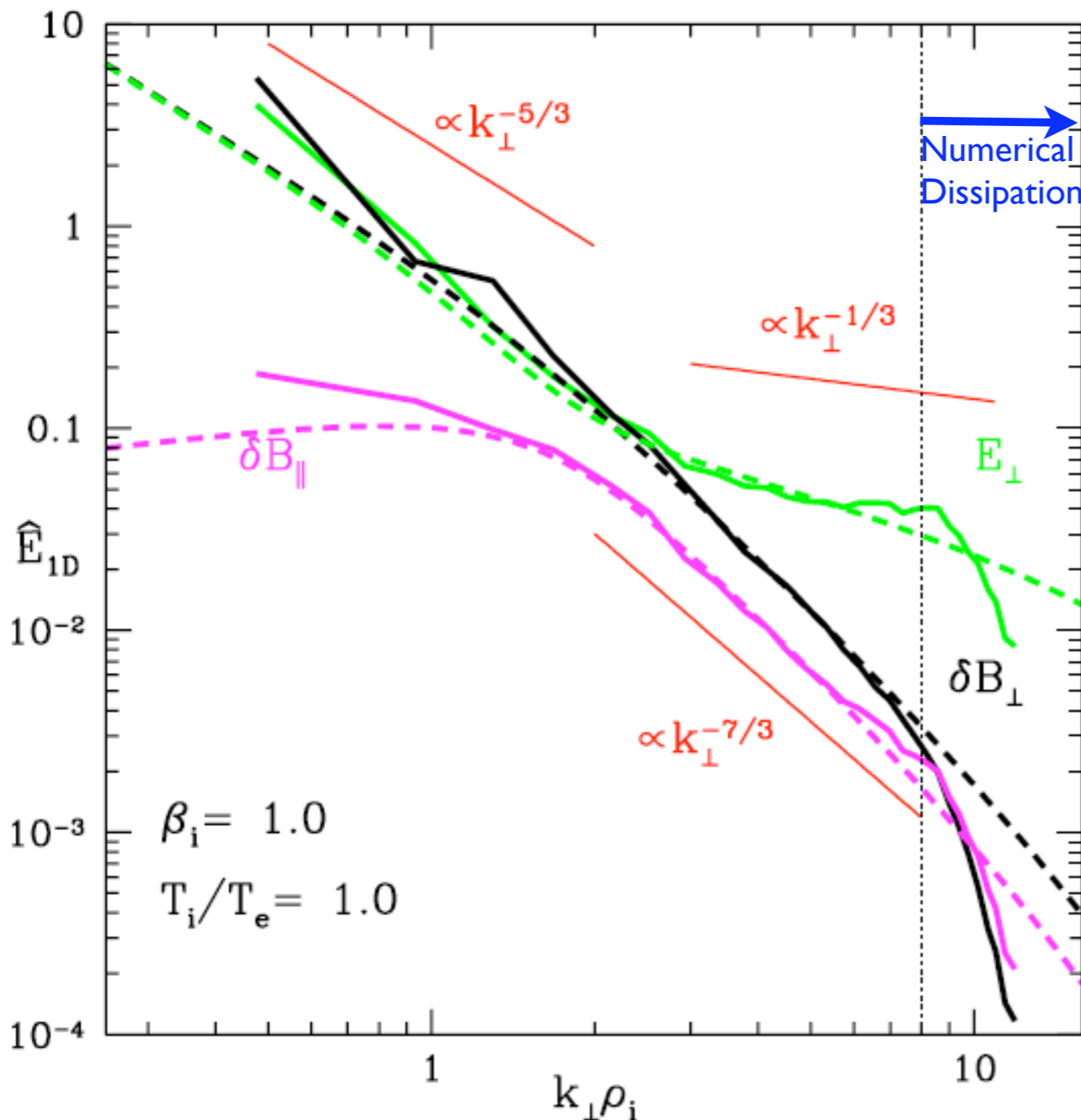
(Howes et al. 2008a)



Numerical results are strongly supportive of model of
transition from MHD Alfven Waves to Kinetic Alfven Waves

- Supports the hypothesis that frequencies remain low, $\omega \ll \Omega_i$

Comparison to Cascade Model



(Howes et al. 2008a, Howes et al. 2008b)

- One fitted parameter
- Cascade model shows excellent agreement with numerical simulation
- Kinetic damping is weak for these plasma parameters
- Need more strongly damped case to better test model

First-Principles Calculation of Turbulence

The use of **Gyrokinetics** enables direct numerical simulations of kinetic turbulence from first principles.

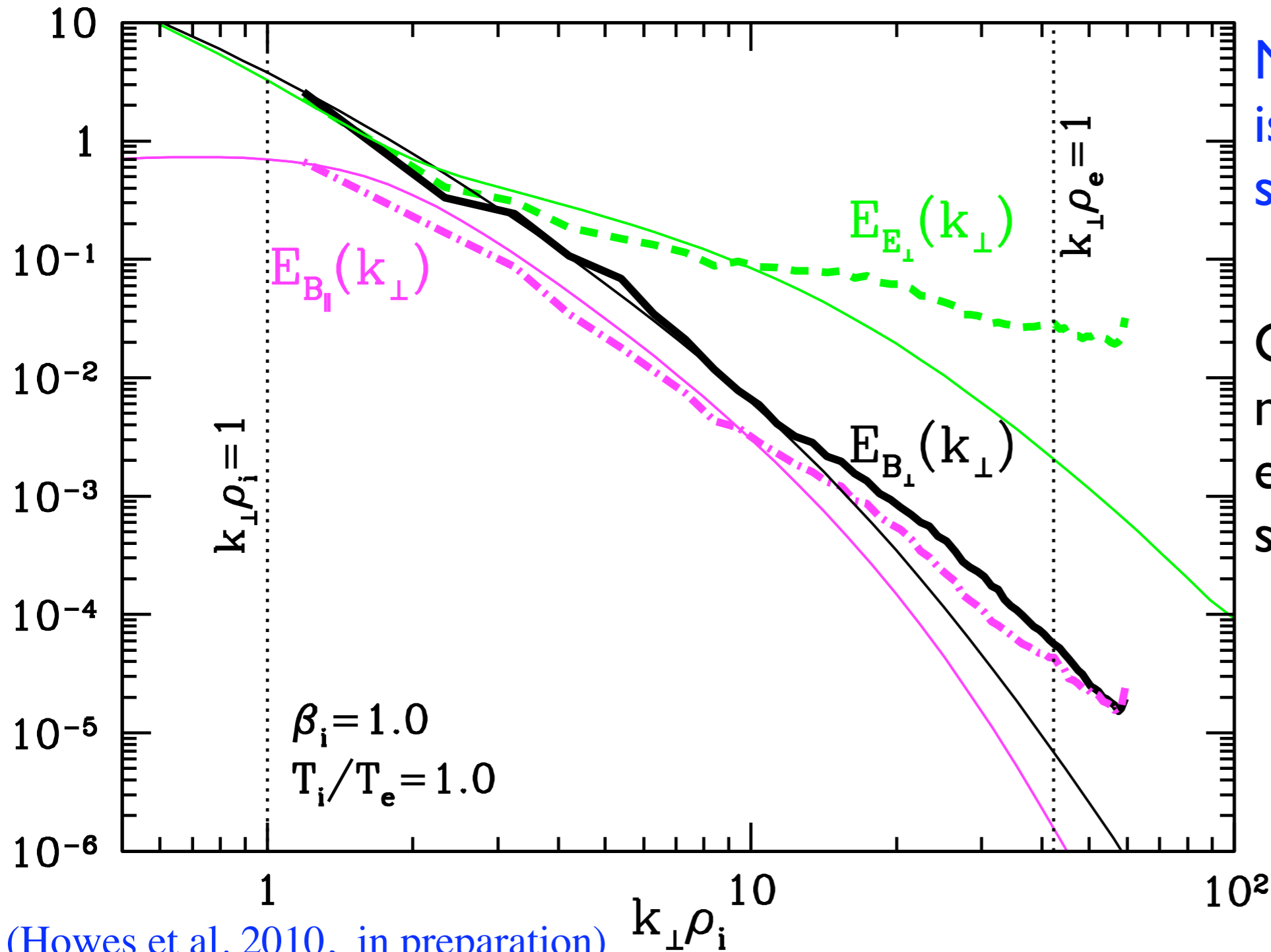
Using the nation's flagship supercomputing resources, today we can achieve turbulence simulations with the following properties:

- Three-dimensional:
Important because turbulent interactions are inherently 3-D
- Employ a physical mass ratio $m_i/m_e = 1836$
- Resolve from the scale of the ion to the electron Larmor radius
From $k_{\perp}\rho_i = 1$ to $k_{\perp}\rho_e = 1$
- Resolve sufficient kinetic damping to terminate the cascade without the need for artificial dissipation

These properties are necessary to make direct comparison to observations

Milestone Dissipation Range Simulation

Ion to Electron Scale Simulation



No artificial dissipation is needed to achieve a steady-state

Original Cascade model predicted an exponential drop off in spectra

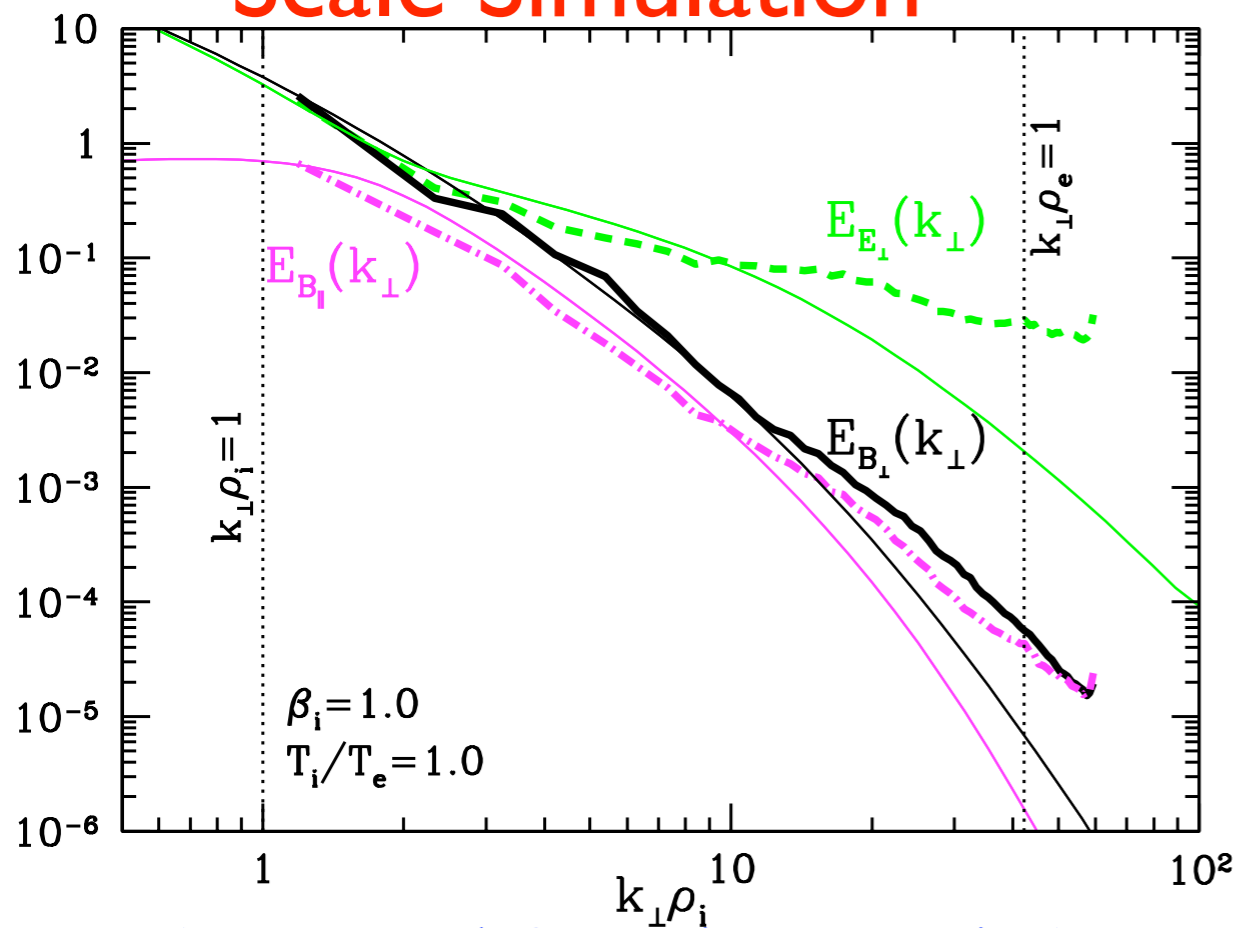
Key Finding:

Demonstrates that a Kinetic Alfvén Wave cascade can reach electron scales

(Howes et al. 2010, in preparation)

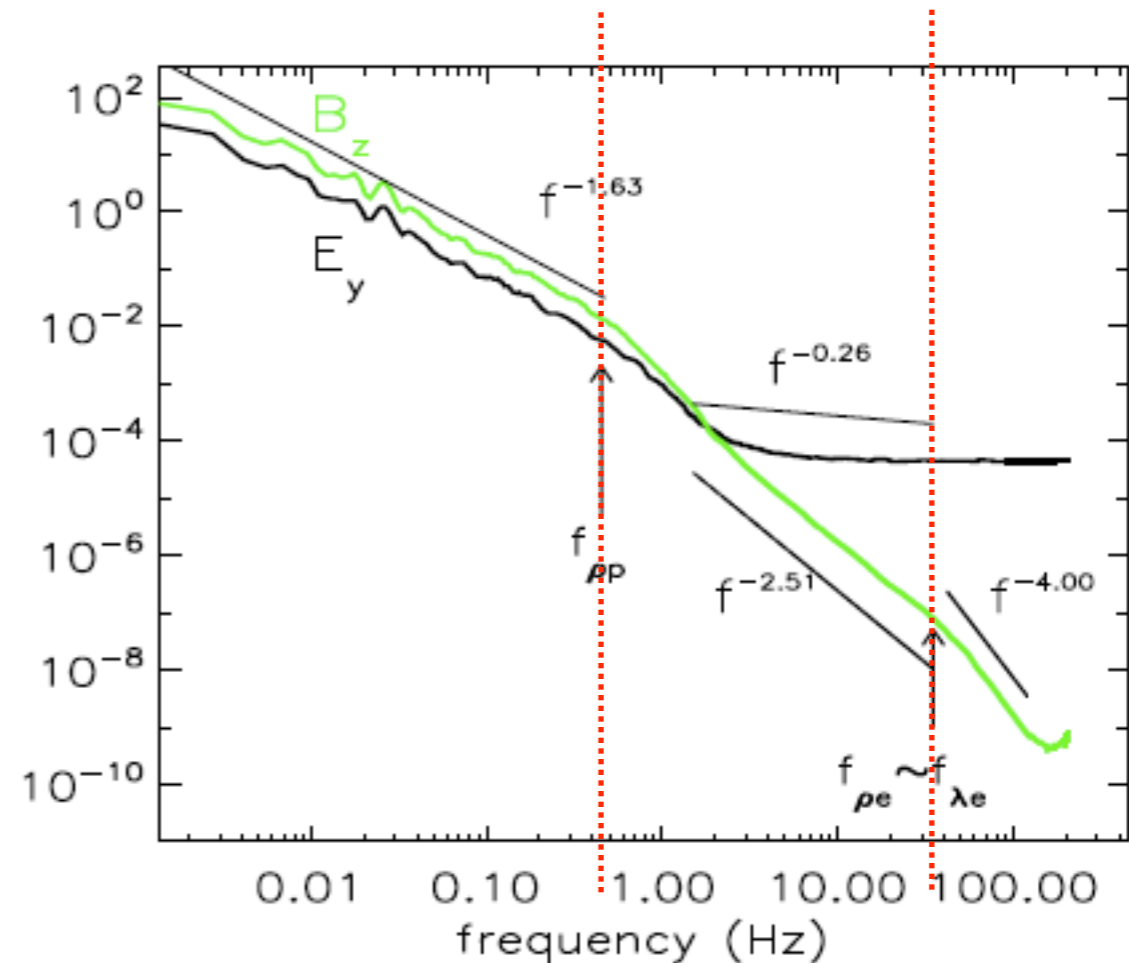
Direct Comparison to Observations

Ion to Electron Scale Simulation



(Howes et al. 2010, in preparation)

Recent Observations



(Sahraoui, Goldstein, Robert, Khotyaintsev 2009, PRL)

Both Numerical Simulations and Observations find nearly power-law behavior down to electron scales!

Cascade Models by Howes et al. 2008 and Podesta et al. 2010:

- Predicted an exponential roll off of spectrum
- Appear not to be in agreement with observations or simulations
- Simulations point to necessary refinements in cascade model!

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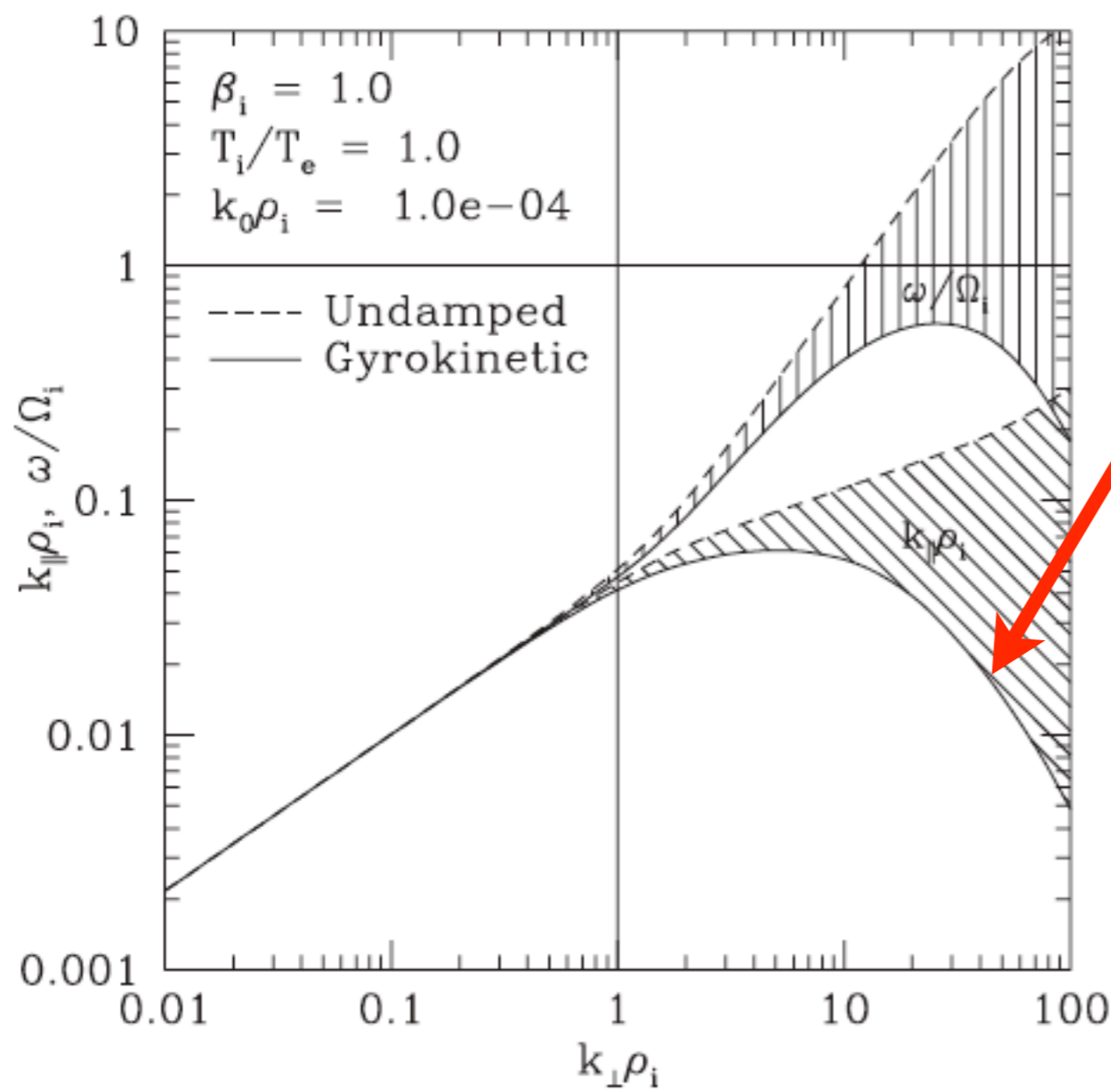
Improved Cascade Model

Problems with the original cascade model: (Howes et al., 2008b)

1. Apparent disagreement with simulations and observations

Dissipation range slopes do not fall off exponentially

2. Assuming critical balance even when turbulence is dissipating



Parallel wavenumber decreases as perpendicular wavenumber increases

This behavior does not seem to make physical sense!

Improved Cascade Model

Original Cascade Model: (Howes et al., 2008b)

1. **Kolmogorov Hypothesis**: Spectrally local nonlinear transfer
2. **Critical Balance** of linear and nonlinear times
3. Applicability of **linear kinetic damping** rates

Weakened Cascade Model:

1. Drop the assumption of Critical Balance

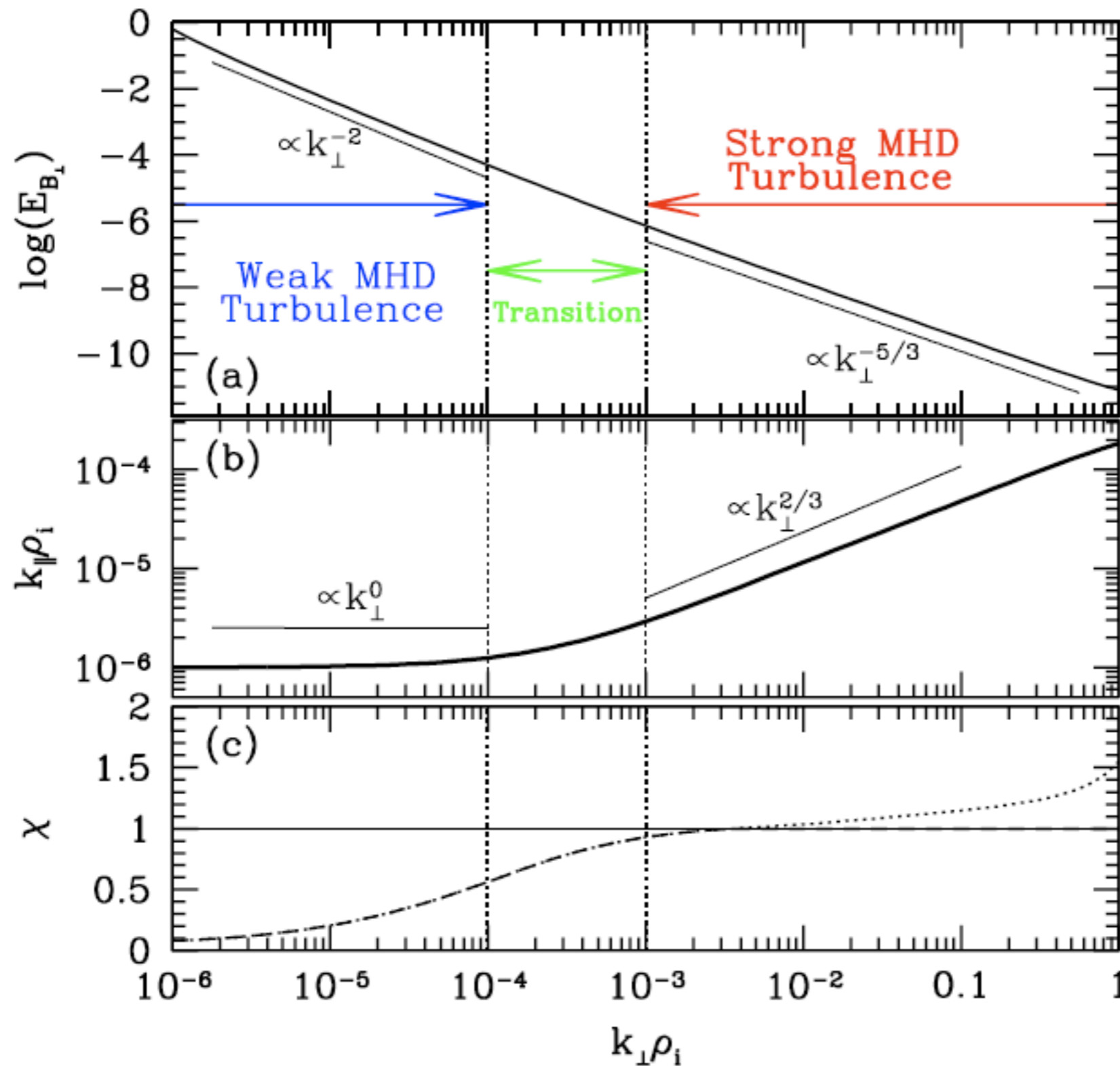
➡ **Model transition between weak and strong turbulence**

2. Drop Kolmogorov's Locality Hypothesis

➡ **Account for effect of nonlocal fluctuations on cascade**

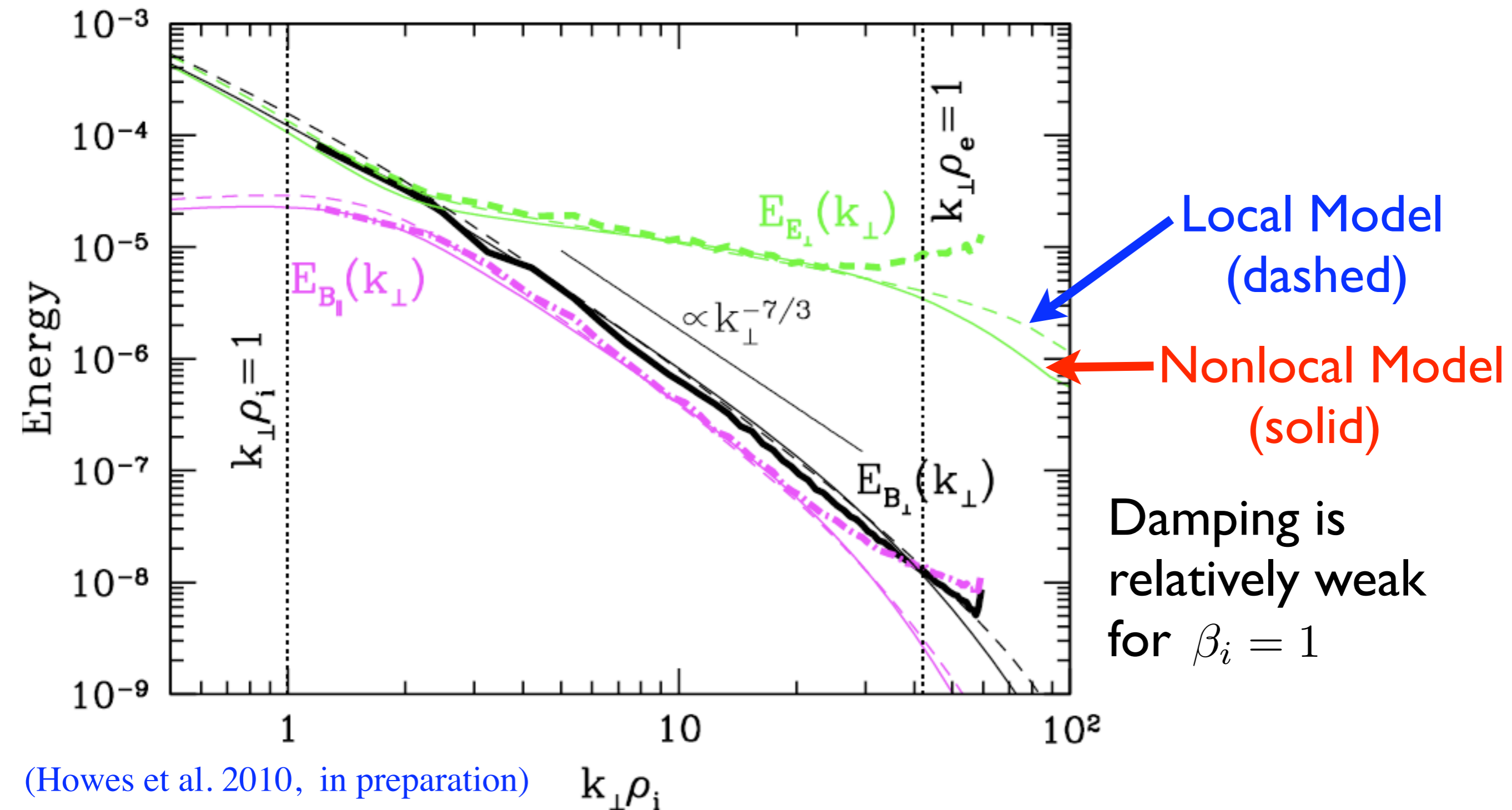
Weak to Strong Turbulence

Weakened Cascade Model



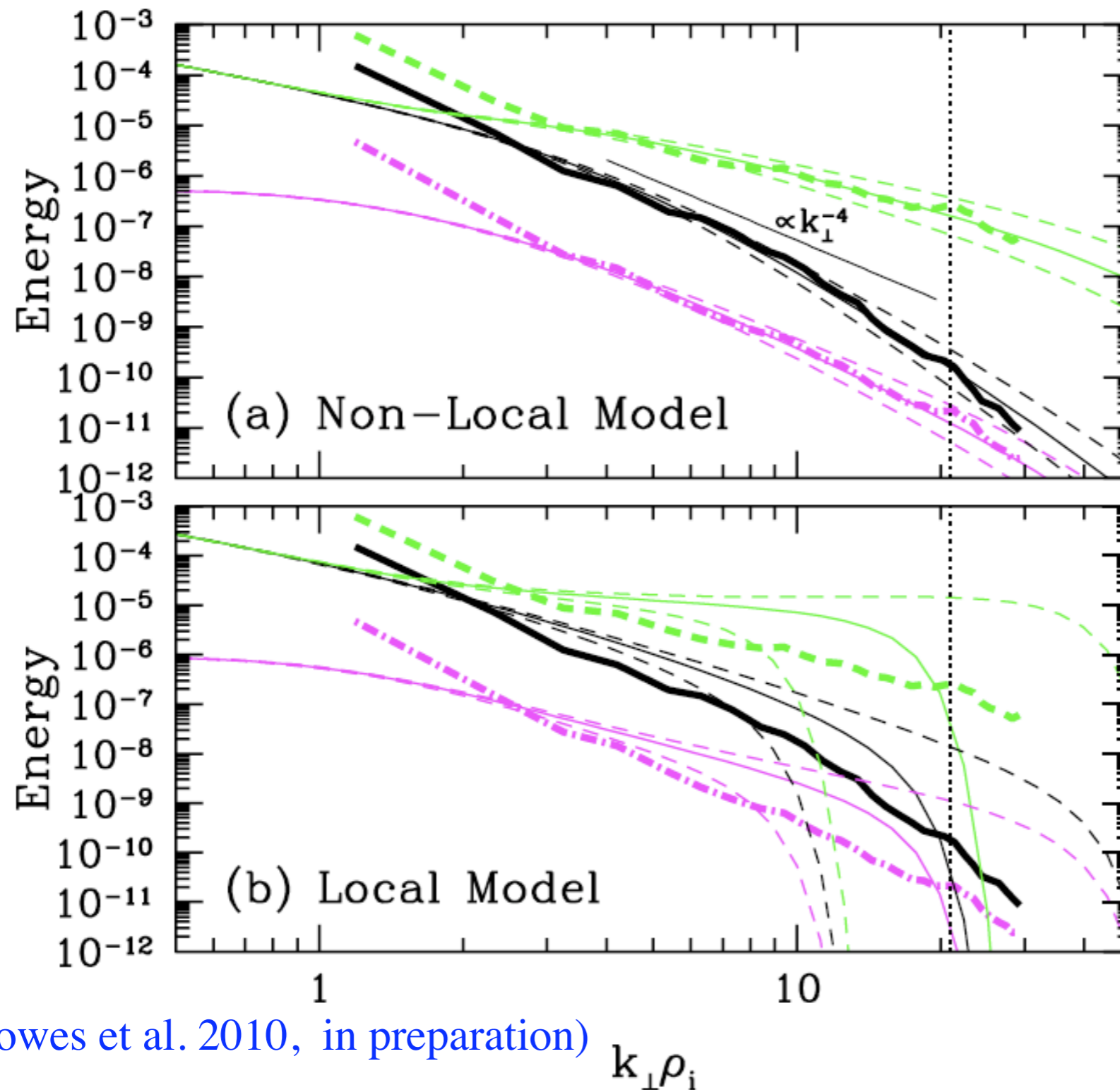
Nonlocal Interactions $\beta_i = 1$

For $\beta_i = 1$ plasma, *both* local and nonlocal models give similar results



Nonlocal Interactions $\beta_i = 0.01$

At $\beta_i = 0.01$, the kinetic damping is significantly stronger



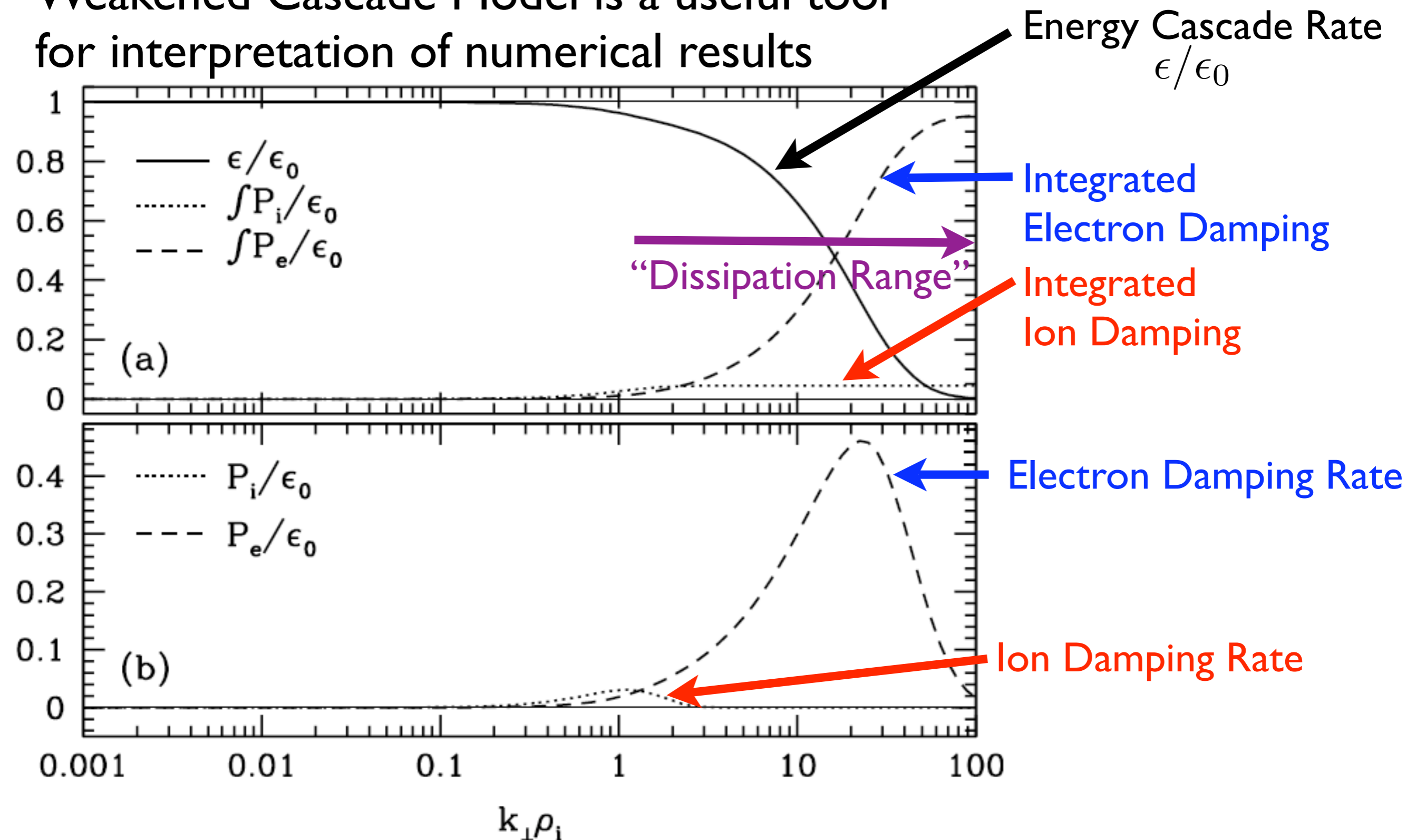
Nonlocal Model
fits spectra well

Local Model is
incapable of fitting
spectra

(Howes et al. 2010, in preparation)

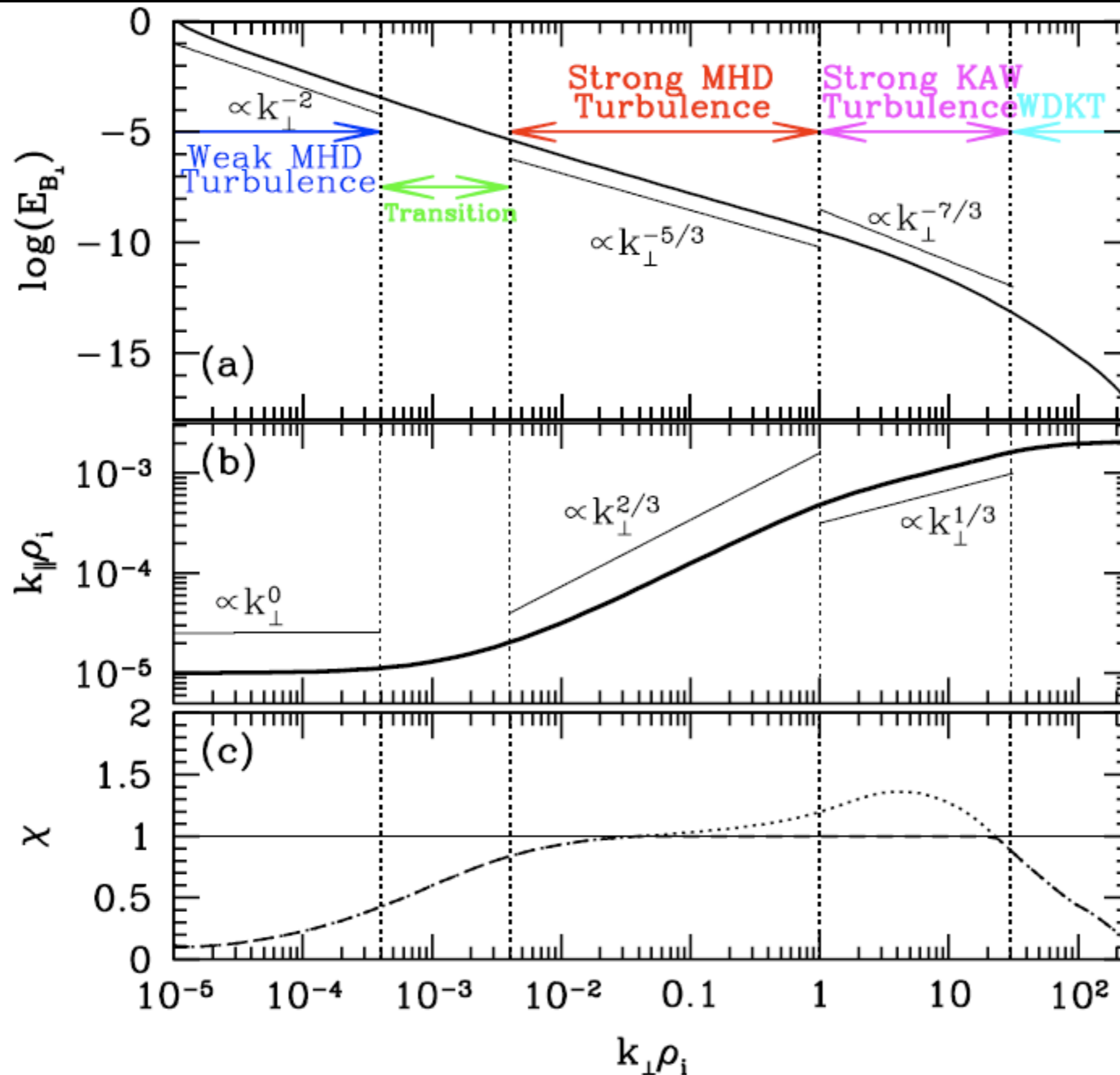
Energy Cascade and Kinetic Dissipation

Weakened Cascade Model is a useful tool for interpretation of numerical results



Future work will compare to ion/electron heating from simulations

Complete Kinetic Turbulence Spectrum



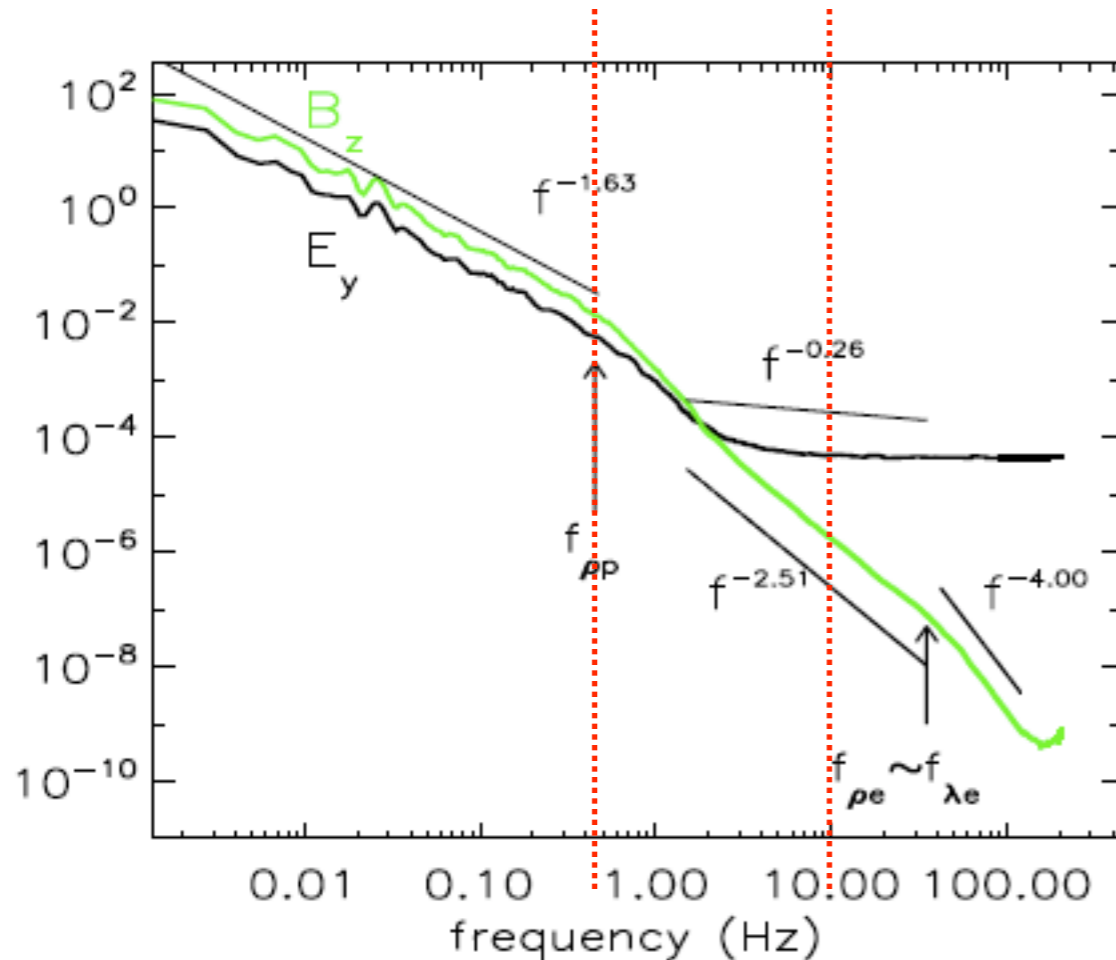
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Question #1

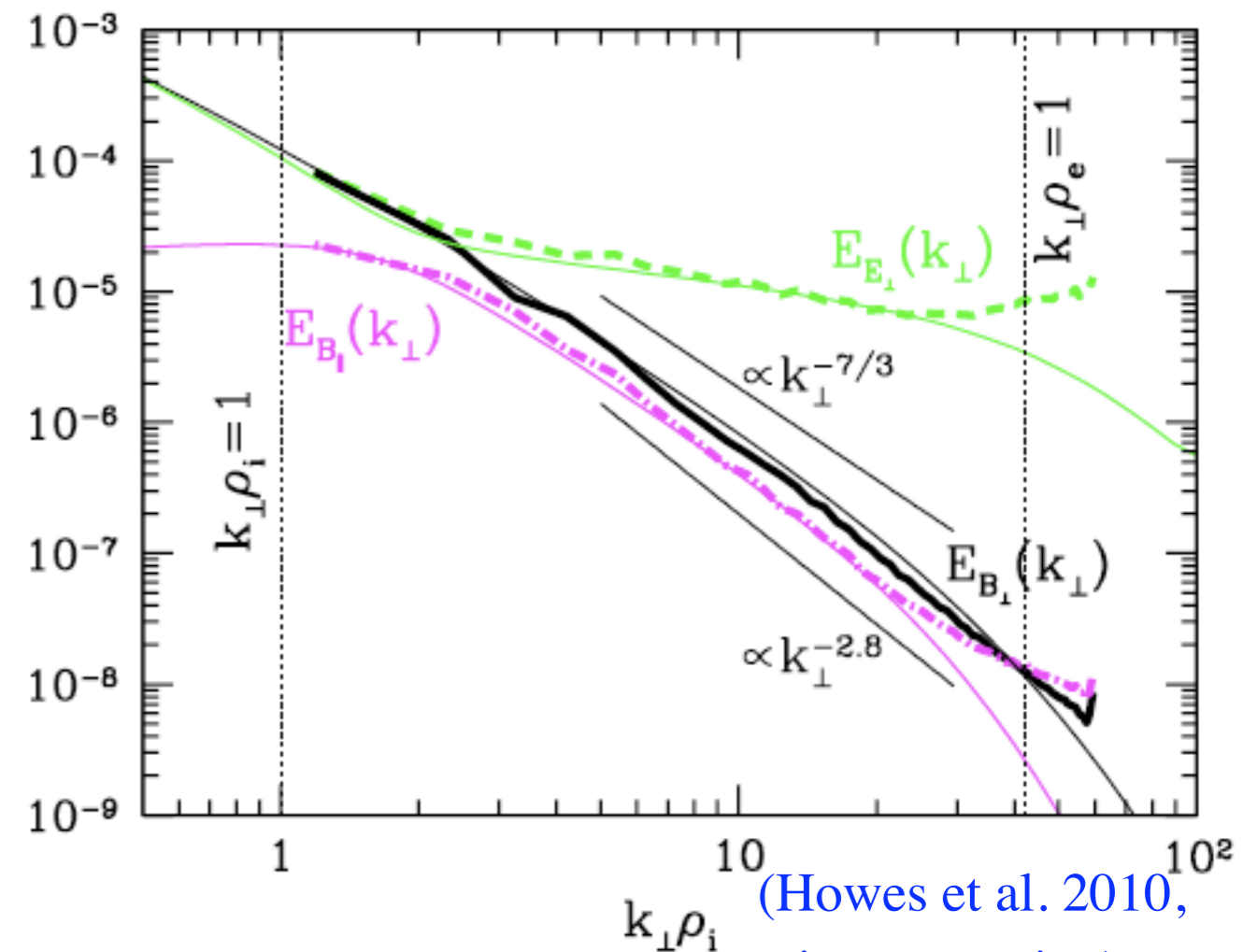
I) Can theoretical models explain the observations?

Solar Wind Observation



(Sahraoui, Goldstein, Robert, Khotyaintsev 2009, PRL)

AstroGK Simulation



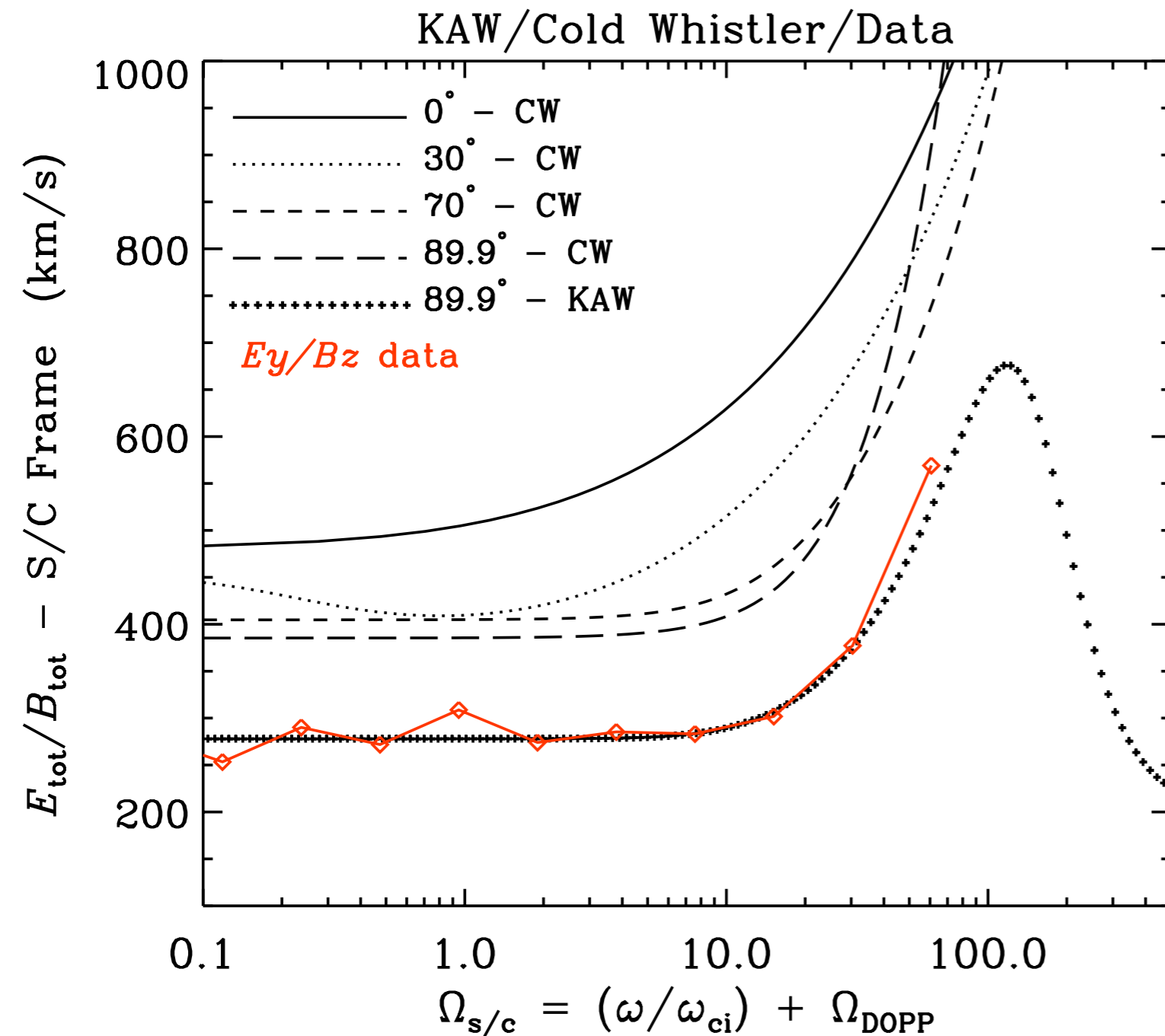
(Howes et al. 2010,
in preparation)

- **Cascade of Kinetic Alfvén Waves** reproduces observed behavior
- Interpretation using **Weakened Cascade Model** suggests importance of **nonlocal effects** on cascade on dissipation range spectrum
- **This evidence supports our theoretical model of kinetic turbulence**

Question #2

2) What are the wave modes that comprise the dissipation range?

KAW or whistler?



This excellent analysis demonstrates that the wave modes within the dissipation range are consistent only with **Kinetic Alfvén Waves**

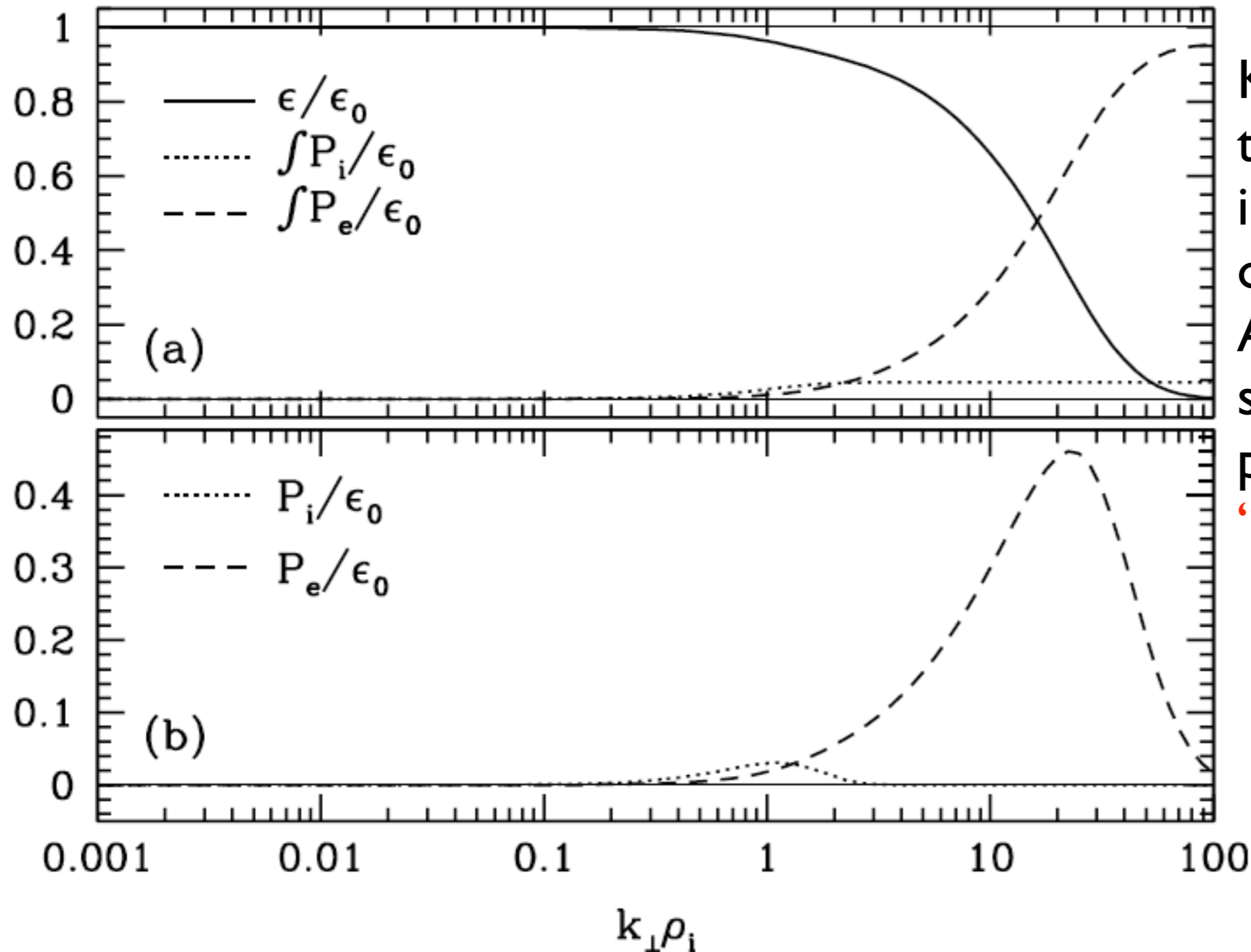
(Salem, Sundkvist, and Bale 2009, Solar Wind 12 Conference)

Question #3

3) Is this a “dissipation range” or a “dispersion range”?

Inertial Range

Dissipation Range



Kinetic damping, via the **Landau resonance**, is generally important over the entire Kinetic Alfvén Wave cascade, so this should properly be called the “**Dissipation Range**”

Conclusions

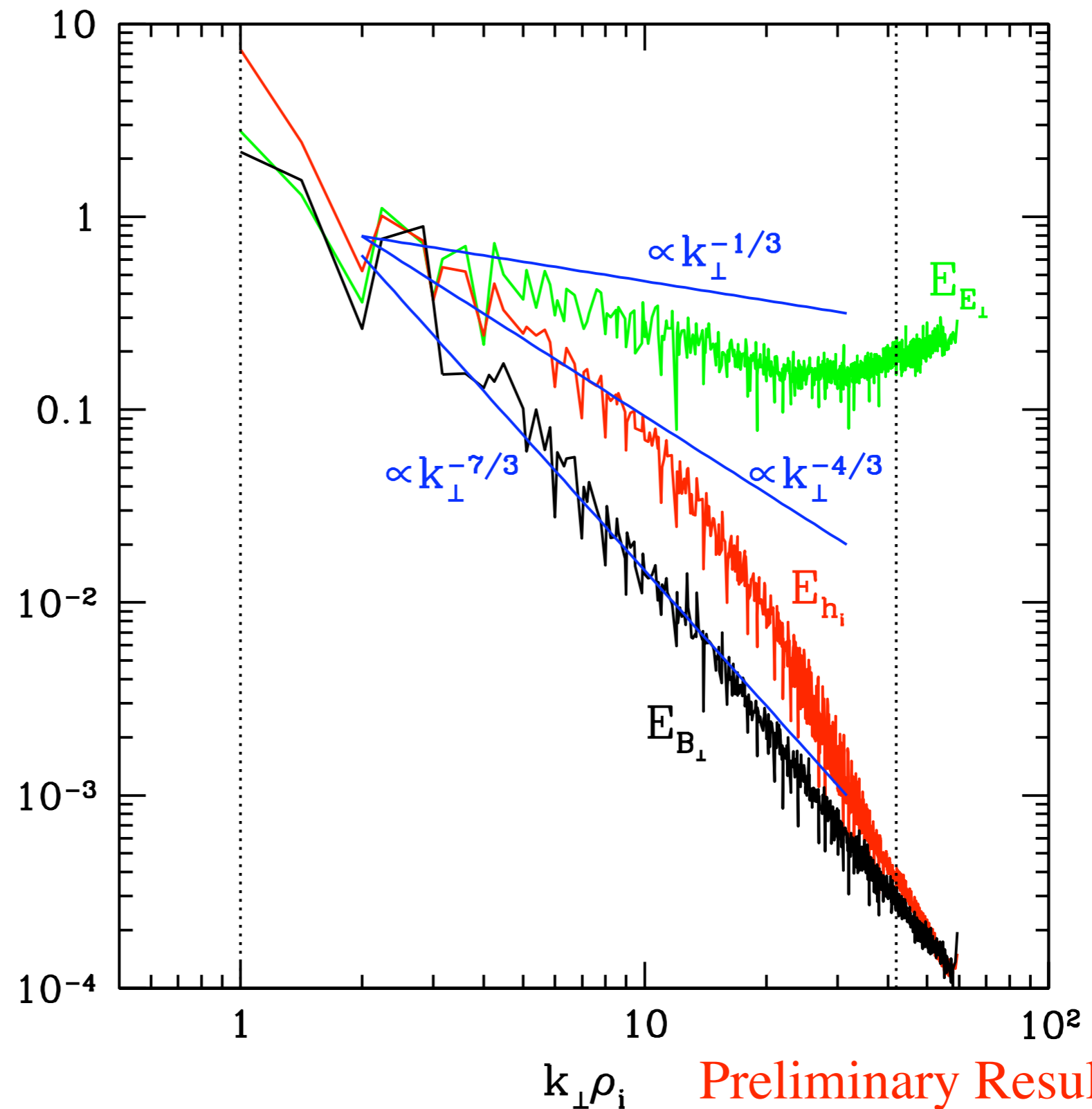
We have proposed a model of the **kinetic turbulent cascade**

- First-principles calculations of the kinetic turbulent cascade using **AstroGK** are an invaluable tool to test this model
- 3-D calculations with physical mass ratio enable direct comparison to observations
 - **Show that a Kinetic Alfvén Wave cascade can reach $k_{\perp} \rho_e \sim 1$**
- Ability to simulate a wide range of plasma parameters enables key testing of analytical models
 - **Identify the need to account for nonlocal effects on cascade**

We have just scratched the surface of questions we can answer with gyrokinetic simulations of kinetic turbulence

Current and Future Work

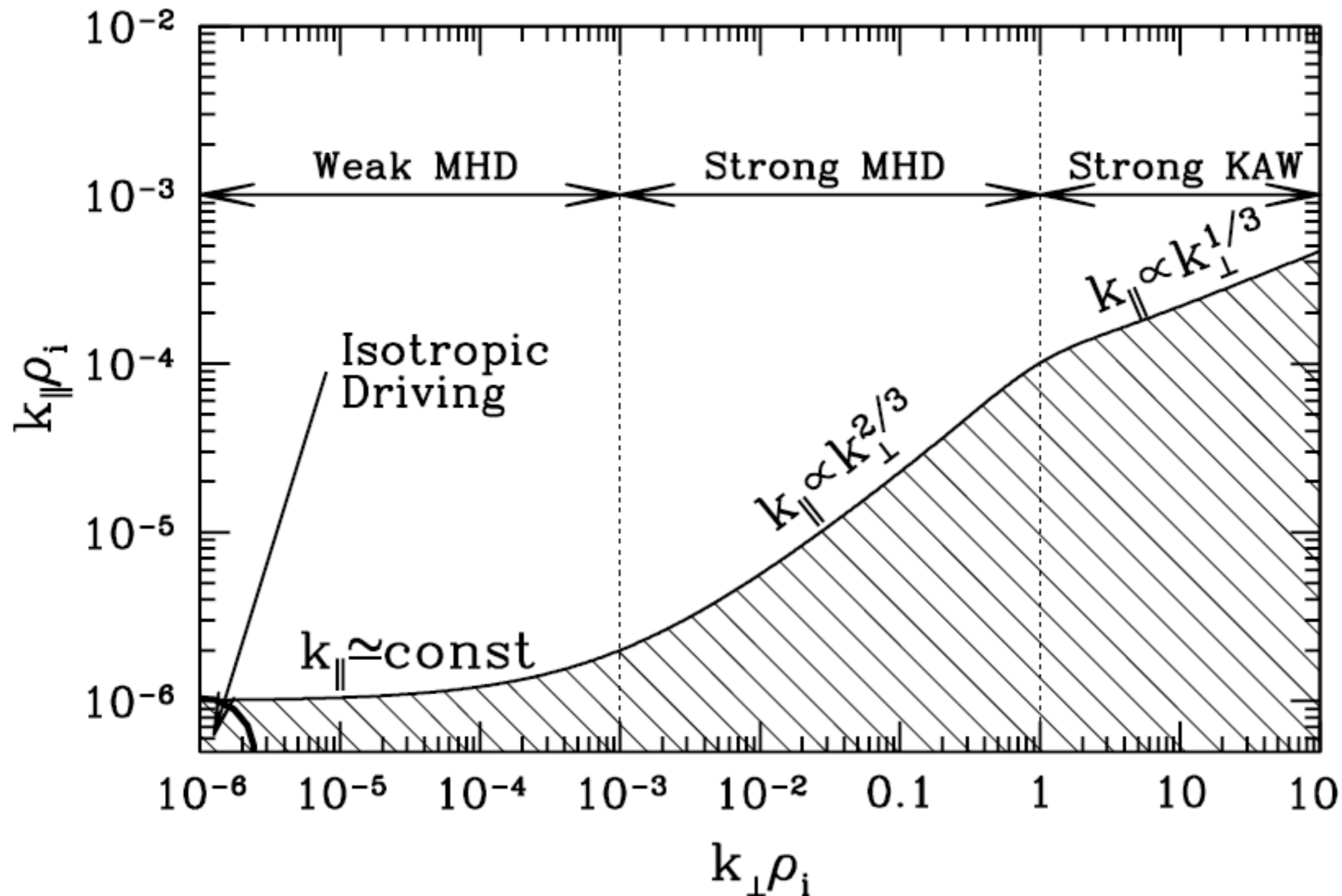
What is the role of entropy cascade in the turbulent dissipation and plasma heating?



Preliminary Results

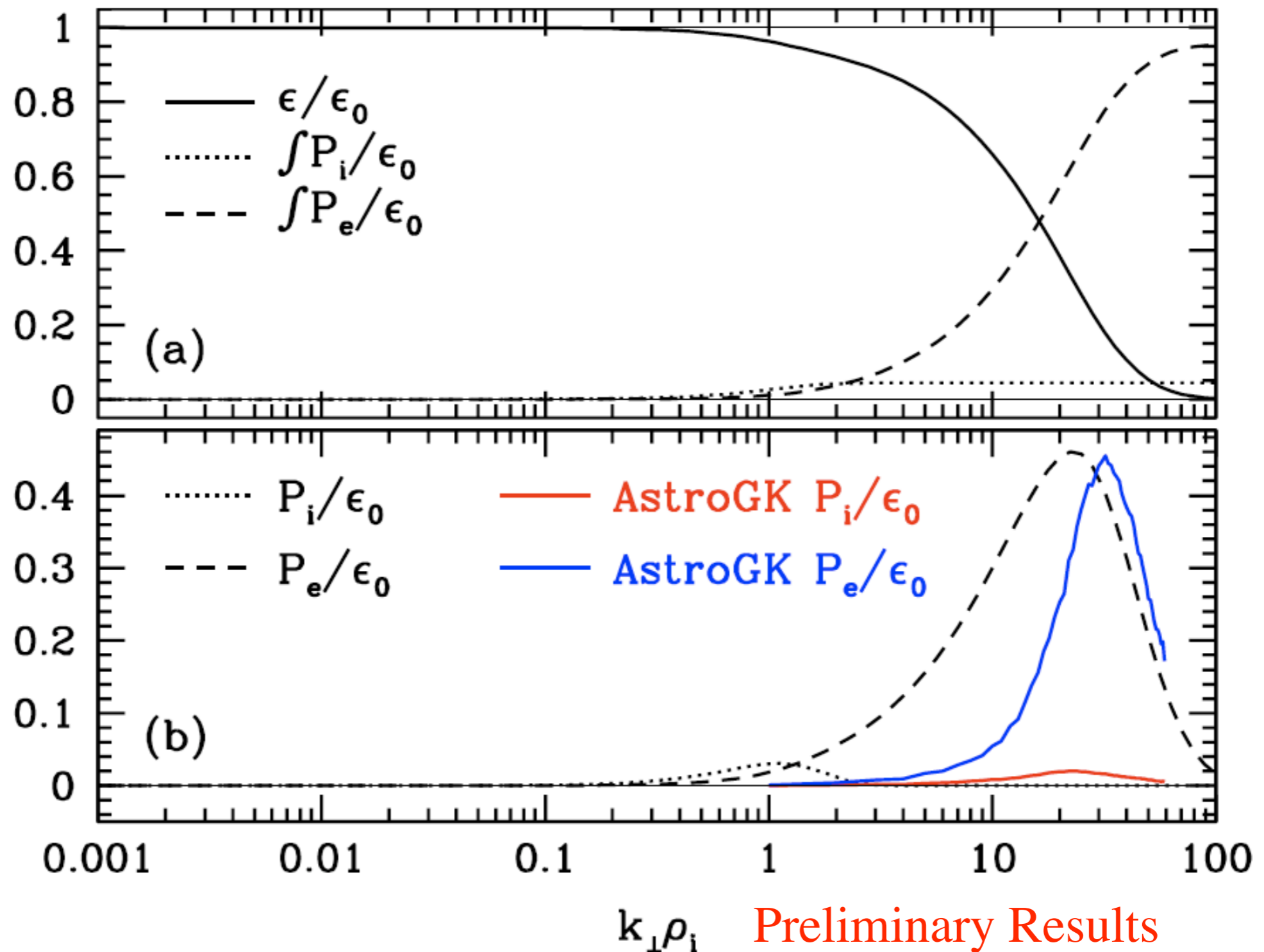
Current and Future Work

What anisotropy characterizes the distribution of power in wavevector space for KAW turbulence?



Current and Future Work

What determines the relative heating of ions, electrons, and minor ions?



THE END

Weakened Cascade Model

Magnetic Energy Continuity $\frac{\partial b_k^2}{\partial t} = -k_{\perp} \frac{\partial \epsilon}{\partial k_{\perp}} + S - 2\bar{\gamma} k_{\parallel} v_A b_k^2$

Nonlinear Energy Transfer **Linear Kinetic Damping**

Energy Cascade Rate $\epsilon = C_1^{-3/2} \omega_{nl} b_k^2$ **Nonlinear Frequency**

$$\omega_{nl}(k_{\perp}) = \int_{k_{\perp 0}}^{k_{\perp max}} d \ln k'_{\perp} \omega_{nl}^{(loc)}(k'_{\perp}) \times \left[\Theta(k_{\perp} - k'_{\perp}) + \frac{k_{\perp}^2}{k_{\perp}'^2} \Theta(k'_{\perp} - k_{\perp}) \right]$$

where $\omega_{nl}^{(loc)}(k'_{\perp}) = \chi(k'_{\perp}) k'_{\perp} b_k(k'_{\perp}) \bar{\omega}(k'_{\perp})$

Parallel Cascade $\frac{d \ln k_{\parallel}}{d \ln k_{\perp}} = \left[\frac{2/3 + (1/3)(k_{\perp} \rho_i)^2}{1 + (k_{\perp} \rho_i)^2} \right] \chi^2$

Nonlinearity Parameter $\chi = \frac{C_2 k_{\perp} b_k}{k_{\parallel} v_A}$

Nonlocal Effects on Cascade

