Internal Transport Barriers at JE

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Joint European Torus: JET

JET is the largest fusion device in the world capable of operating with mixture of Deuterium/Tritium.

- Tokamak: plasma confinement in a torus by means of magnetic field



Plasma volume: 80m² Magnetic field: 2.0-3.4T (4T) Current: 1.0-3.5 (4.5MA)



Tokamak confinement

- In order to achieve fusion conditions one aims to obtain the best possible (energy) confinement with high temperatures and pressures in the core.
- This requires large gradients in the pressure (temperature)





Transport in Tokamak plasmas

- At high gradients, turbulent modes may grow in Tokamak plasmas that dominate the transport processes.
 - Thus transport is not diffusive and these turbulent modes will set a critical gradient yielding so-called stiff (ion) temperature profiles.





Internal Transport Barriers

- But the profile stiffness can be broken (locally) yielding socalled internal transport barriers (ITBs)
 - ITBs provide an opportunity to improve the Tokamak confinement
 - Studying ITBs may improve our understanding of turbulence







- How do we make internal transport barriers at JET?
 - What triggers the growth of strong ion ITBs?
 - How important is rotational shear?
- Compare JET ITBs with those in JT-60U
 - What about electron and particle transport?
- Some general remarks about ITB and their use in Tokamaks
- Give references to more detailed descriptions



How to make an ITB at JET

- Empirical recipe to form strong internal ion transport barriers
 - Optimised q-profiles with low or negative magnetic shear (q'/q)
 - Similar recipe used in various Tokamaks (JT-60U, DIII-D, ...)





How to make an ITB at JET

Empirical recipe to form strong ion internal transport barriers

- Focus here on q-profiles with central negative magnetic shear (q'/q)
 - The q profile develops in time (current diffusion, impact of bootstrap current)
- Use significant Neutral Beam Injection (NBI) heating (rotation?)





How to make an ITB at JET

A sudden appearance/growth of a strong ITB can be seen
 The growth of the ITB could eventually be limited by the onset of disruptive kink modes (infernal modes)





Triggering ITBs in JET plasmas

- A trigger mechanism starts the ITBs growth
 - The physics mechanism is however not well understood
 - Different for plasmas negative¹ or low² central magnetic shear
 - For the latter it is related to the appearance of a rational q_{min}



[1] E. JOFFRIN, Nucl. Fusion 43 (2003) 1167[2] E. JOFFRIN, et al., Nucl. Fusion 42 (2002) 235



Triggering ITBs in JET plasmas

- A trigger mechanism starts the ITBs growth
 - The physics mechanism is however not always understood
 - The formation of ion ITBs in JET are usually not predicted from theory based transport/turbulence models^{1,2}
 - Such triggers are also found to act in JT-60U, DIIID, ...

[1] Y.F. BARANOV, et al., Plasma Phys. Control. Fusion 46 (2004) 1181.
[2] T. TALA, T, et al., Nucl. Fusion 46 (2006) 548.



ITBs and plasma rotation at JET

- How important is the NBI ingredient? \rightarrow rotation?
 - Can we make strong ITBs without fast plasma rotation?
- Experiments on ITBs at JET were carried out, where the plasma rotation was changed by:
 - Replacing the NBI by ICRH ion heating^{1,2}
 - Not easy to keep the heat flux unchanged
 - Applying larger toroidal field ripples^{2,3}
 - Change rotation independent from heat flux

N.C. HAWKES, et al., Contribution to the 32nd EPS Conference (Warsaw) 2008.
 P.C. DE VRIES, et al., Nucl. Fusion 49 (2009) 075007.
 P.C. DE VRIES, et al., Plasma Phys. Control. Fusion 50 (2008) 065008.



TF ripple and plasma rotation

JET has the unique capability to alter its toroidal field ripple.

- This has a significant effect on the torque on the plasma¹ but less on the heat deposition by NBI and ICRH
- In combination with momentum transport effects (pinch) a higher TF ripple yields a lower rotation and smaller rotation gradients².



[1] P.C. DE VRIES, et al., Nucl. Fusion 48 (2008) 035007.
[2] P.C. DE VRIES, et al., Plasma Phys. Control. Fusion 50 (2010) 065004



ITBs and plasma rotation

Increasing the TF ripple amplitude results in a reduction of the rotational shear:

- has a detrimental effect on the growth of the ITB¹.
- But an ITB triggering event is still visible!²



[1] P.C. DE VRIES, et al., Plasma Phys. Control. Fusion 50 (2008) 065008.

[2] P.C. DE VRIES, et al., Nucl. Fusion 49 (2009) 075007.



ITBs and plasma rotation

- At JET, replacing NBI by more ICRH and hence reducing the rotation, had similar effects^{1,2}
 - Detrimental effect on the growth of the ITB.
 - But an ITB triggering event is still visible!
- JT-60U and DIII-D were able to do such experiments using balanced NBI^{4,5}

N.C. HAWKES, et al., Contribution to the 32nd EPS Conference (Warsaw) 2008.
 P.C. DE VRIES,, et al., Nucl. Fusion 49 (2009) 075007.
 P.C. DE VRIES, et al., Plasma Phys. Control. Fusion 50 (2008) 065008.
 Y. SAKAMOTO, et al., Nucl. Fusion 41 (2001) 865
 M.W, SHAFER, M.W, et al., Phys. Rev. Lett. 103 (2009) 075004.



Example from DIII-D

Similar observations have also been made at DIII-D¹.



[1] M.W, SHAFER, M.W, et al., Phys. Rev. Lett. 103 (2009) 075004.



Interpretation

- At the time the transport barrier forms/triggers:
 - for high TF ripple or a larger ICRH fractions: $\omega_{ExB} \sim 1-2 \cdot 10^4 [s^{-1}]$ almost one order of magnitude below the ITG growth rate γ_{ITG}
 - for low TF ripple and high NBI fractions: $\omega_{ExB} \sim 6.10^4 [s^{-1}]$ of the order of ITG growth rate γ_{ITG}
- Detailed modelling with the GYRO¹ code showed
 - That γ_{ITG} =6-7 10⁴ s⁻¹ without rotational shear (high TF ripple)
 - For low TF ripple and rotational shear:
 - At time of triggering: γ_{ITG} =1.5 10⁴ s⁻¹
 - During the growth phase ITG modes are fully stabilized.
- The JET plasmas are sub-critical to suppress the turbulence, but the trigger of a seed barrier pushes it over the threshold
 - Reversed shear/minimum q not needed for ITG stabilization

[1] J. CANDY, and R.E. WALTZ, Phys. Rev. Lett. (2003) 045001



ITB growth in JET

- The ITB will enhance the gradient in toroidal rotation
 - Thus the ITB itself may be able further increase $\omega_{\text{ExB}}/\gamma_{\text{ITG}}$.
 - That γ_{ITG} =6-7 10⁴ s⁻¹ without rotational shear (high TF ripple)



Using ITBs in an AT scenario

- Beside being of interest for the understanding of the underlying transport processes, ITBs could be used in so-called Advance Tokamak (AT) scenarios.
 - AT scenario's aim to have a fully non-inductive current drive
- Requires steady-state, stable and wide ITBs



- ITB triggering similar between JET and JT-60U
 - Example: RS q_{min}=3 ITBs for JT-60U and JET (same TF ripple, NBI)
 - Stronger ion ITBs formed in plasmas with more rotational shear at the time they are triggered.



[1] P.C. DE VRIES, et al., Plasma Phys. Control. Fusion 51 (2009) 124050.



- The dynamic behaviour is not identical¹
 - Example: RS q_{min} =3 ITBs for JT-60U and JET (same TF ripple, NBI)
 - It is not always easy to capture the dynamic behaviour of ITBs
 - Difference in density profile \rightarrow different j_{BS} \rightarrow different q profile develop
 - Impact of ELMs/pedestal



[1] P.C. DE VRIES, et al., Plasma Phys. Control. Fusion 51 (2009) 124050.



- Beside the ion transport channel also the transport of particles and electron transport channel and that of particles has to be considered.
 - It is known that negative magnetic shear is beneficial for the formation transport barriers in the electron channel¹.



[1] Y.F. BARANOV, et al., Plasma Phys. Control. Fusion 46 (2004) 1181.



- Beside the ion transport channel also the transport of particles and electron transport channel and that of particles has to be considered¹.
 - The slightly larger negative shear in JT-60U affected the particle transport → more peaked density profiles^{1,2}





[1] X. LITAUDON, et al., IAEA Fusion Energy Conference (2010) Deajon, Korea
[2] C BOURDELLE, et al., PHYS. PLASMAS 14 (2007) 112501



- Beside the ion transport channel also the transport of particles and electron transport channel and that of particles has to be considered.
 - The slightly larger negative shear in JT-60U affected the particle transport → more peaked density profiles → larger bootstrap current →





[1] X. LITAUDON, et al., IAEA Fusion Energy Conference (2010) Deajon, Korea



- Beside the ion transport channel also the transport of particles and electron transport channel and that of particles has to be considered.
 - The slightly larger negative shear in JT-60U affected the particle transport → more peaked density profiles → affects ITG turbulence

[1] X. LITAUDON, et al., IAEA Fusion Energy Conference (2010) Deajon, Korea



Summary & Conclusions (1)

- In JET rotational shear is not sufficient to trigger the growth of ion ITBs and a special trigger mechanism is at play, related to the appearance of a rational q_{min}.
- But the further growth of the ITB is affected by the level of rotational shear at the time of triggering. Larger ion ITBs grow in plasmas with sufficient rotational shear.
- A Very similar picture is found in other devices.
 - Same physics!
- Can ion ITBs form without the need of a trigger mechanism, but just because of sufficient rotational shear?
 - In JET basic rotational shear is usually not enough (need trigger)
 - May well be possible in devices with very large Mach numbers



Summary & Conclusions (2)

- An ITB is a local reduction in transport or turbulence and this can be achieved in many different ways
 - Sharapov shift, strongly peaked density profiles, fast particles,...
 - Transport channels for electrons/particles and ions differ
 - Hence these results do not necessarily apply to all ITBs one has to be careful comparing ITBs but basically:
 - Rotational shear seems to be important for the formation of ion ITBs, while magnetic shear affects electron heat and particle transport.
- When developing (or modelling) an ITB Tokamak scenario the complete picture of all transport channels/turbulence and the interplay between them needs to be considered.





Rotational shear and TF ripple

The rotational shear or shearing rate ω_{ExB} has been calculated under the assumption of neo-classical poloidal rotation.



[1] P.C. DE VRIES, et al., Nucl. Fusion 49 (2009) 075007.



ITBs and plasma rotation

- Increasing the ICRH fraction and reducing the rotation
 - Even in plasmas with little torque/rotation ITB triggers were found
 - But the growth of ITB was limited in these plasmas



[1] N.C. HAWKES, et al., Contribution to the 32nd EPS Conference (Warsaw) 2008
[2] P.C. DE VRIES, et al., Nucl. Fusion 49 (2009) 075007.



Turbulence and Profile Stiffness

- Turbulence plays a dominant role in Tokamak transport
 - Ion temperature gradient driven modes \rightarrow 'stiff' temperature profiles
 - Profile gradient is determined by ITG stability, not by the heat flux and neo-classical diffusivity.





Turbulence and Profile Stiffness

- However, detailed experiments at JET have shown that the stiffness is reduced in plasmas with more rotation¹.
 - Rotation is thought to reduce the (ITG) turbulence growth rate
 - Here the profile stiffness is found to be affected
 - How do ion ITBs fit into this picture?





TF ripple experiments and Pinch

- TF ripple affects the toroidal rotation
 - When the TF ripple is increased, the toroidal rotation profile is affected but torque flux (ρ =0.5) is not
 - This suggests momentum transport is altered.





TF ripple experiments and Pinch

- Increasing TF ripple reduces the rotation/momentum in the outer part of the plasma, yielding a smaller effect of pinch
 - Hence, less peaked rotation profiles and larger effective Prandtl nr.



[1] P.C. de Vries, et al. PPCF 52 (2010) 065004.



Magnitude of the Pinch

What is the magnitude of the momentum pinch to explain these observations?

$$\frac{\nabla\Omega}{\Gamma_{\phi}} = -\frac{V_p}{\chi_{\phi}}\frac{\Omega}{\Gamma_{\phi}} - \frac{1}{\chi_{\phi}}$$

Averaged over discharges $\chi_{\phi} = 1.5(m^2 / s)$ $\chi_i = 1.2(m^2 / s)$ $V_p = 1.4(m / s)$



[1] P.C. de Vries, et al. PPCF 52 (2010) 065004.



Estimated pinch from Database

Assuming that the momentum diffusivity is equal to the ion heat diffusivity one could estimate the magnitude of the pinch for all entries in the database.

$$\begin{split} \Gamma_{\phi} &= -\chi_{\phi}^{eff} \nabla \Omega \\ \Gamma_{\phi} &= -\chi_{\phi} \nabla \Omega - V_{p} \Omega \\ & \downarrow \\ \frac{V_{p}}{\chi_{i}} \approx (1 - P_{r}^{eff}) \frac{\nabla \Omega}{\Omega} \end{split}$$

- Scales with $R/L_n^{1,2}$ - For H-modes: $2 < RV_p / \chi < 10$

[1] P.C. de Vries, et al. PPCF **52** (2010) 065004.
[2] A.G. Peeters, Phys. Rev. Lett. **98** (2007) 265003.

