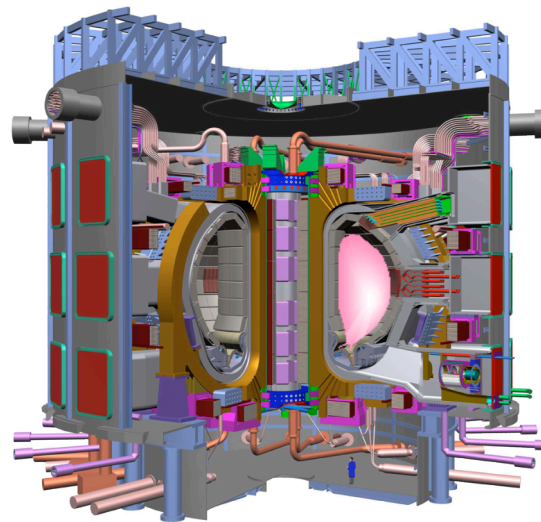


Shape Matters -- Building a Better Bottle?

Steve Cowley *Culham, Imperial*

Bill Dorland, K³, Omar Hurricane and Pierre Gourdain,



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Causes Without Results

“I am afraid I rather give myself away when I explain. Results without causes are much more impressive.”

Sherlock Holmes in The Stock-Broker's Clerk, Arthur Conan Doyle.

Better Bottle?

We have a magnetic configuration that will take us to burning plasmas in ITER. This will probably be the configuration of the first generation of fusion reactors.

Can we improve this? What would that mean?

Better confinement? τ_E

Higher pressure -- lower B? β

More reliable? stability.

Less heat flux? Divertor.

Has every configuration been tried? In 2D?



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Requirements For Fusion.

$$\text{Fusion Power} \propto n_D n_T T^2 \propto \beta^2 B^4$$

$$10 \text{keV} < T < 20 \text{keV}.$$

Rough criterion for ignition.

$$nT\tau_E > 3 \times 10^{15} \text{cm}^3 \text{keV s}$$

Physics limits the achievable values of these quantities.

n : Density “Greenwald” limit.

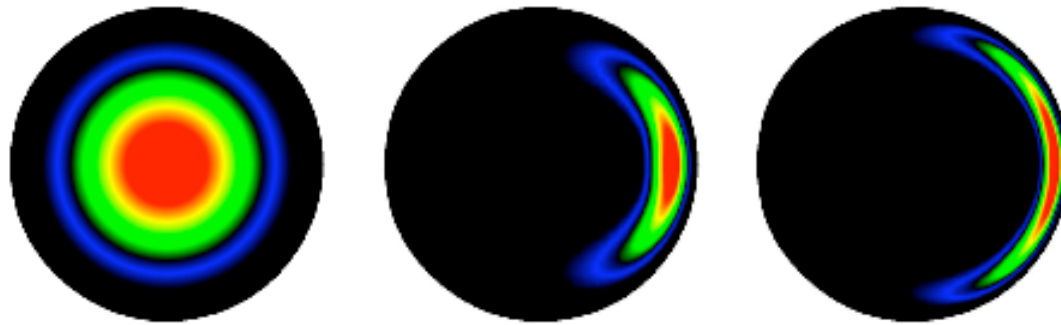
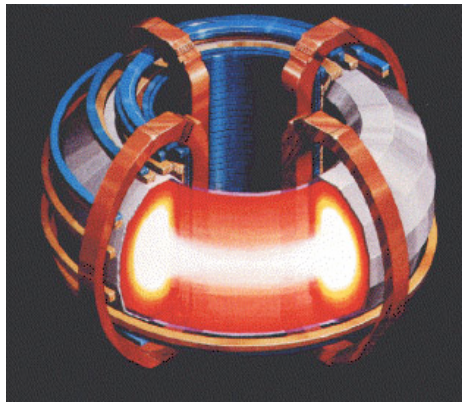
$$nT = \beta \frac{B^2}{8\pi} \quad \text{: Beta Limit.} \quad \beta = \beta_N \frac{I}{aB}$$

τ_E : Turbulence.

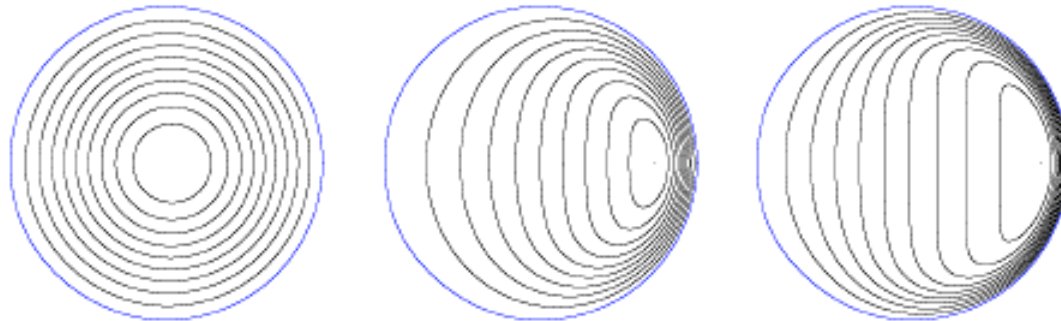


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Raising Beta



$$\mathbf{J} \times \mathbf{B} = \nabla p$$



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Radial Force Balance

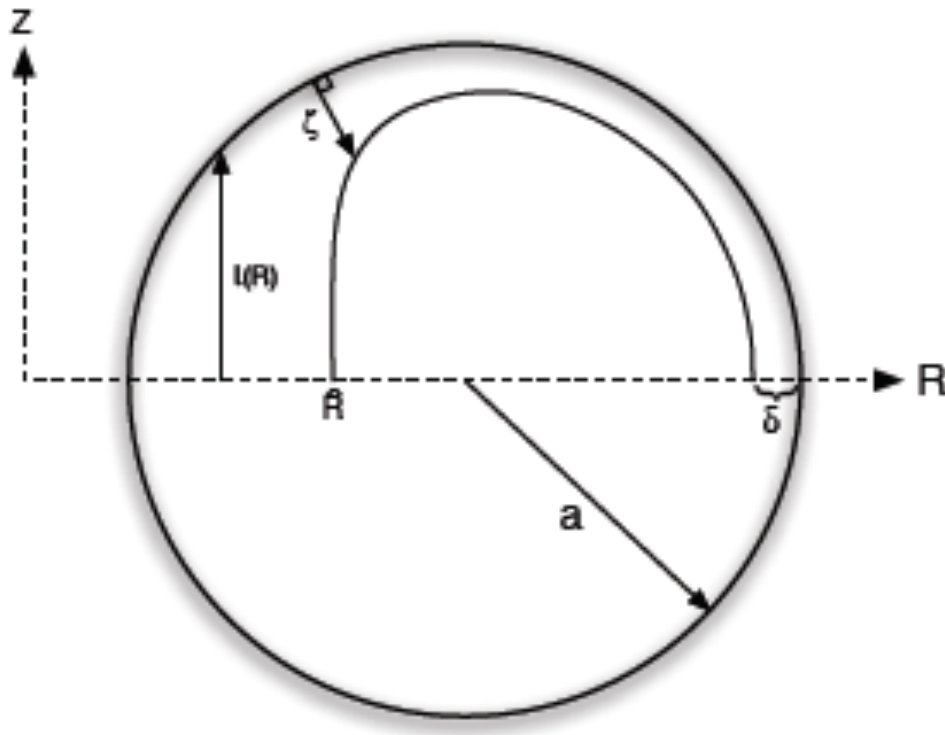
$$\nabla\psi \cdot [\mathbf{J} \times \mathbf{B}] = \nabla p$$

$$\left[R \frac{\partial}{\partial R} \left(\frac{1}{R} \frac{\partial}{\partial R} \right) + \frac{\partial^2}{\partial Z^2} \right] \psi = -\mu_0 R^2 \frac{dp}{d\psi} - F \frac{dF}{d\psi}$$

Grad-Shafranov Equation.

where

$$\mathbf{B} = \frac{\nabla\psi \times \mathbf{e}_T}{R} + \frac{F(\psi)}{R} \mathbf{e}_T$$



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The Small Parameter

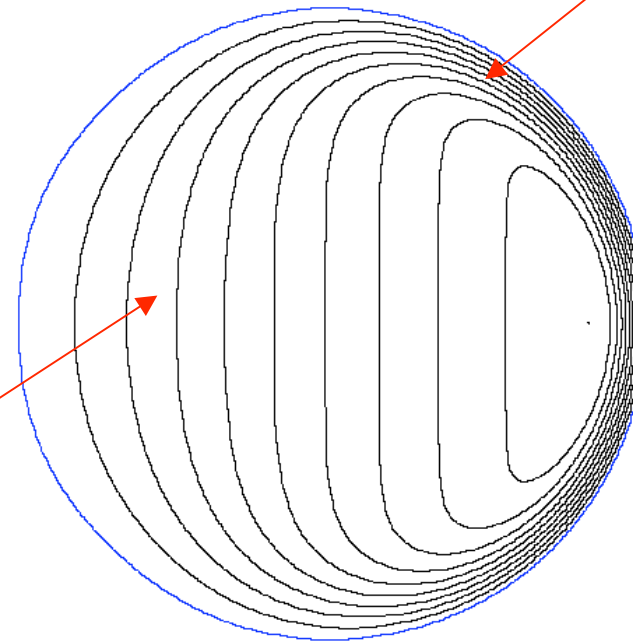
$$\frac{\epsilon}{q^2} = \frac{a}{q^2 R} \ll 1$$

$$\mathbf{B}_p \ll \mathbf{B}_T = F/R$$

Surfaces shift and
Squash against edge.

CORE

Boundary Layer



$$\underbrace{\nabla p}_{\mathcal{O}(\beta)} + \underbrace{\frac{\nabla B_T^2}{2}}_{\mathcal{O}(1)} + \underbrace{\frac{\nabla B_P^2}{2}}_{\mathcal{O}(\epsilon^2/q^2)} + \underbrace{\frac{B_T^2 \nabla R}{R}}_{\mathcal{O}(\epsilon)} + \underbrace{B_P \cdot \nabla B_P}_{\mathcal{O}(\epsilon^2/q^2)} = 0$$



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CORE

$$\left[R \frac{\partial}{\partial R} \left(\frac{1}{R} \frac{\partial}{\partial R} \right) + \frac{\partial^2}{\partial Z^2} \right] \psi = -\mu_0 R^2 \frac{dp}{d\psi} - F \frac{dF}{d\psi}$$

small

$\mathcal{O}(\epsilon)$

$$\mu_0 R^2 \frac{dp}{d\psi} = -F \frac{dF}{d\psi}$$

Toroidal Field
Confinement

$$R = R(\psi) \text{ or } \psi = \psi(R)$$

Straight vertical flux surfaces.



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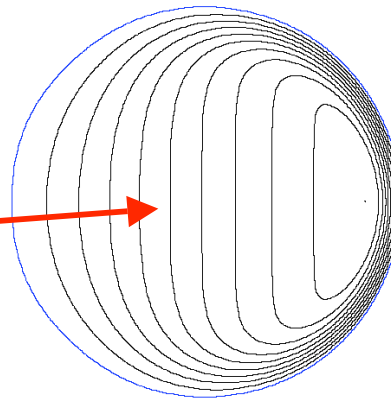
CORE

$$F(\hat{R}) = \sqrt{2 \left(C - \mu_0 \int_{R_{\min}}^{\hat{R}} \hat{R}'^2 \frac{dp}{d\hat{R}'} d\hat{R}' \right)}$$

C = constant & p increases and F decreases towards the axis

$$R = R(\psi) \text{ or } \psi = \psi(R)$$

Straight vertical
flux surfaces in
core



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Boundary Layer -- BL

Gradients are large perpendicular to wall ξ = distance to wall.

$$\frac{\partial^2 \psi}{\partial \xi^2} = -\mu_0 (R^2 - \hat{R}^2(\psi)) \frac{dp}{d\psi}$$

Width of Boundary Layer is small and Poloidal Field is strong

$$\left(\frac{\partial \psi}{\partial \xi} \right)^2 = -2\mu_0 \int_R^{\hat{R}} (R^2 - \hat{R}''^2) \frac{dp}{d\psi} \frac{\partial \psi}{\partial \xi} d\xi$$

Poloidal field pressure forces balance the residual force from Lack of cancellation of pressure and toroidal field forces.

$$|\mathbf{B}_p| \sim \sqrt{\epsilon p} \ll |\mathbf{B}_T|$$

Boundary Layer --- BL

$$\xi(R, \hat{R}) = \int_{R_{\min}}^{\hat{R}} \frac{d\hat{R}' \frac{\partial \psi}{\partial \hat{R}'}}{\sqrt{-2\mu_0 \int_R^{\hat{R}'} d\hat{R}'' \frac{dp}{d\hat{R}''} (R^2 - \hat{R}''^2)}}$$

$$\delta = a \sqrt{\frac{\epsilon}{q^2 \beta}}$$

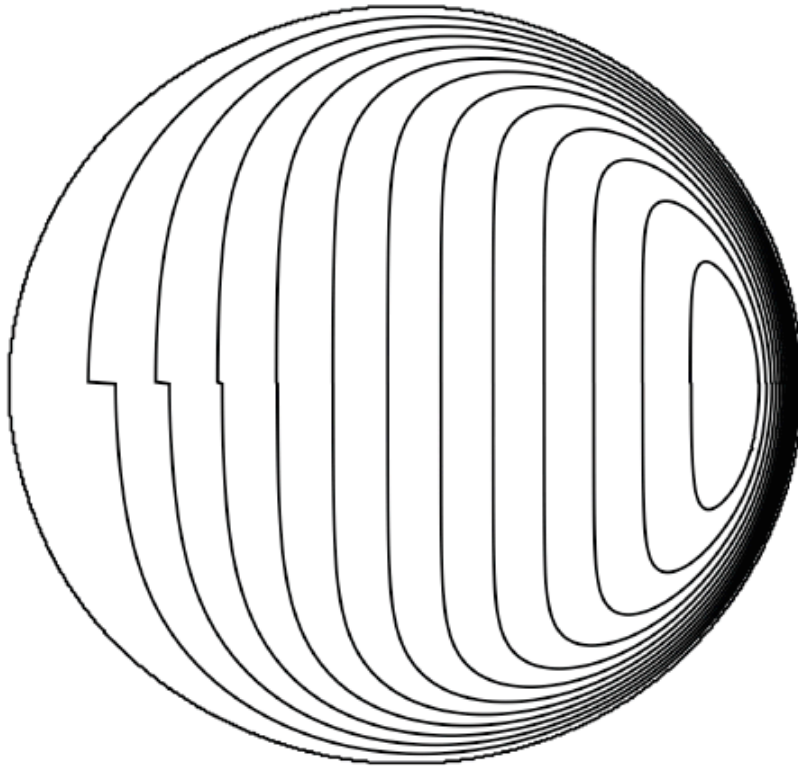
Boundary layer width.
Expansion works if $\delta < a$

Poloidal field increases outwards
in Boundary Layer.



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Comparison



Agreement gets better as we increase beta

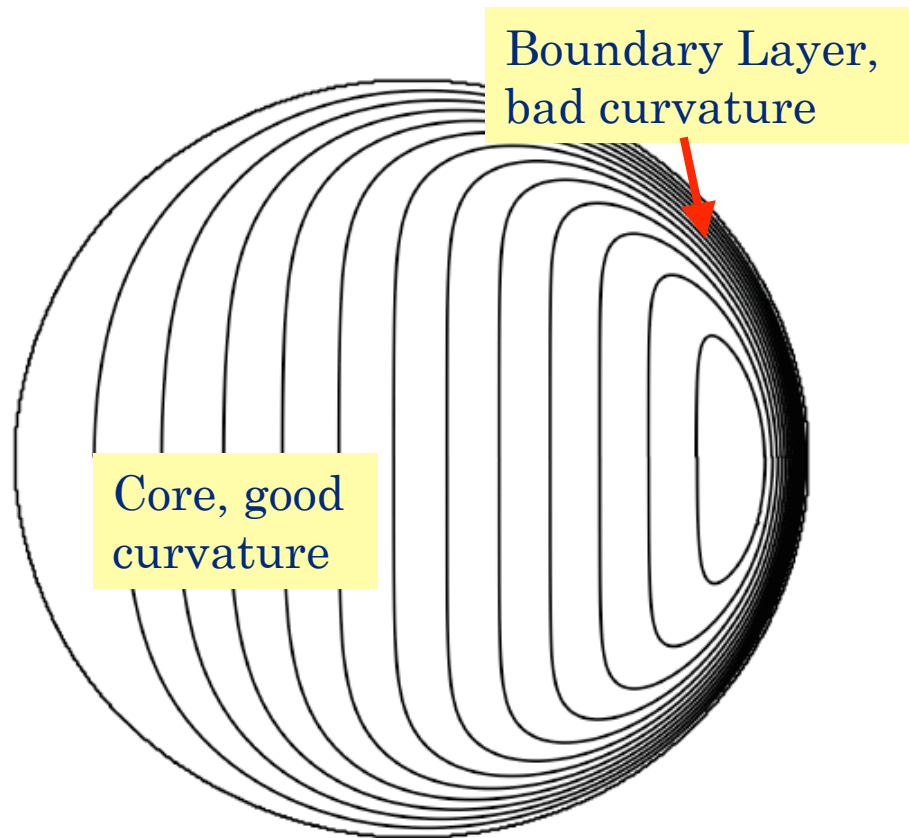
FIG. 3: Comparison of a equilibrium solution computed in CUBE (top) and the same solution calculated using the analytic theory (bottom).



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Good properties.

1. Good Average Curvature



Bad field line curvature in the boundary layer only. Core dominates average

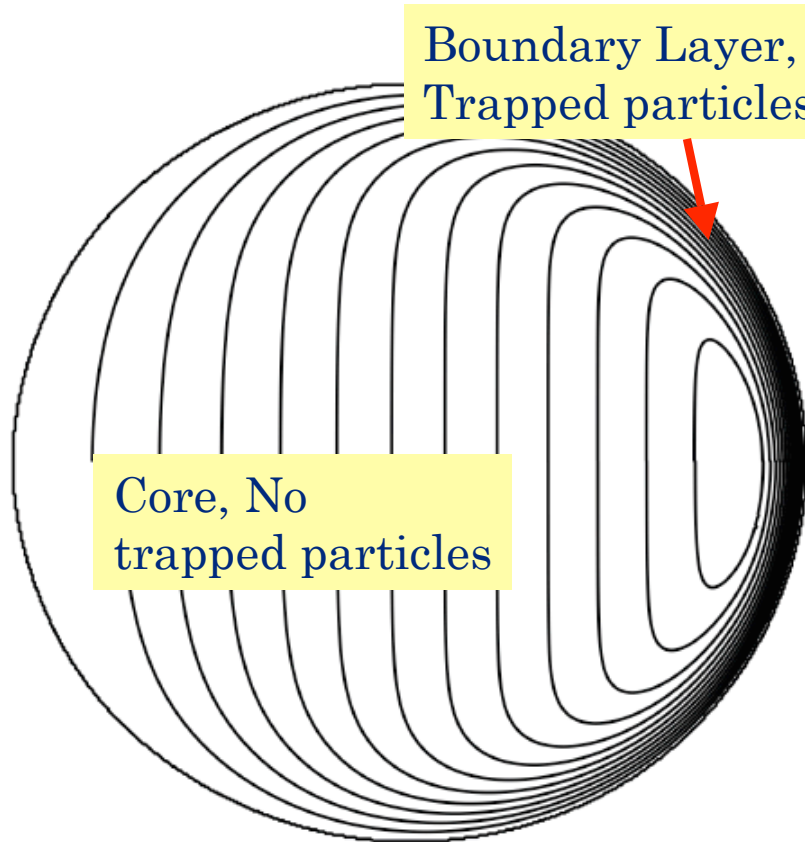
$$\langle \nabla p \cdot (\mathbf{b} \cdot \nabla \mathbf{b}) \rangle \sim -\frac{p}{aR}$$

Mercier stable and tearing mode stable.
Cowley Phys. Fluids B 1991.

FIG. 3: Comparison of an equilibrium solution computed in CUBE (top) and the same solution calculated using the analytic theory (bottom).

Good properties.

2. Small trapped particle fraction



$|\mathbf{B}|$ constant on flux surface in core \Rightarrow no bounce points in core.
Trapped particle fraction.....

$$f_T \sim \left(1 - \frac{B_{min}}{B_{max}}\right)^{1/2} \frac{\Delta V}{V}$$

As beta increases both factors decrease. $|\mathbf{B}|$ constant on flux surface in BL too (omnidigeneity). The volume fraction in BL is

$$\frac{\Delta V}{V} \sim \sqrt{\frac{\epsilon}{q^2 \beta}} \ll 1$$

Banana width is squeezed by strong B_p . Neoclassical transport reduced by more than $\sim \sqrt{\frac{\epsilon}{q^2 \beta}}$ *SCC PPPL report.*

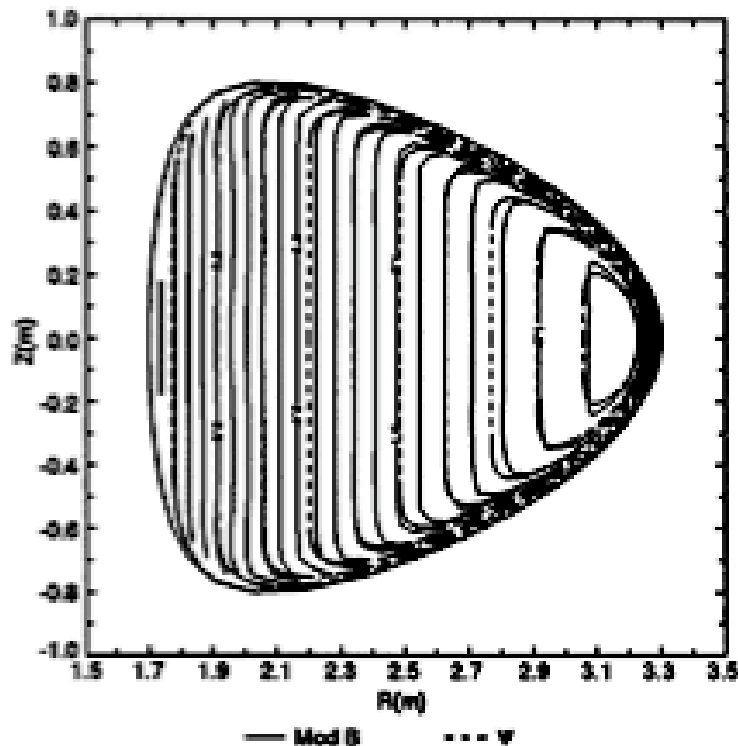


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Good properties. 3. Magnetic well

$$\mathbf{B}^2 = \frac{F^2}{R^2} + \frac{|\nabla\psi|^2}{R^2}$$

$$p + \frac{\mathbf{B}^2}{2} = \text{constant}$$



$|\mathbf{B}|$ is small in the center of
The plasma.

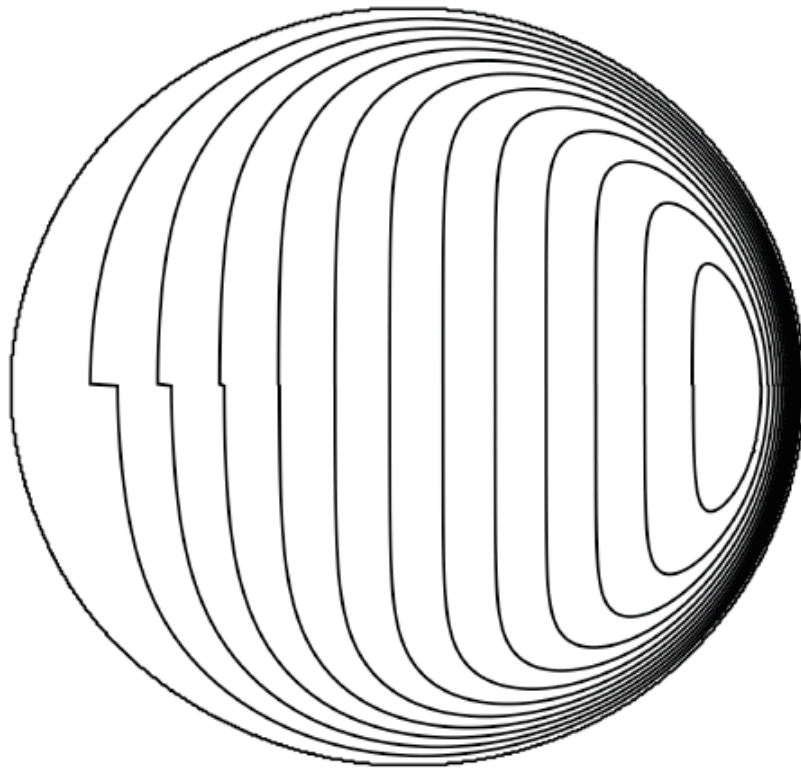
$$\mu = \frac{v_{\perp}^2}{B} = \text{constant}$$

Got to give particles energy
To get them out. Helps
Stability, *Taylor 1963*



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Good properties.
4. Short Connection length



B_p is large in BL so distance
along field from bad to
Good curvature is

$$L_c \sim qR \sqrt{\frac{\epsilon}{q^2 \beta}}$$

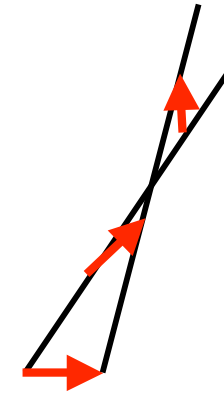
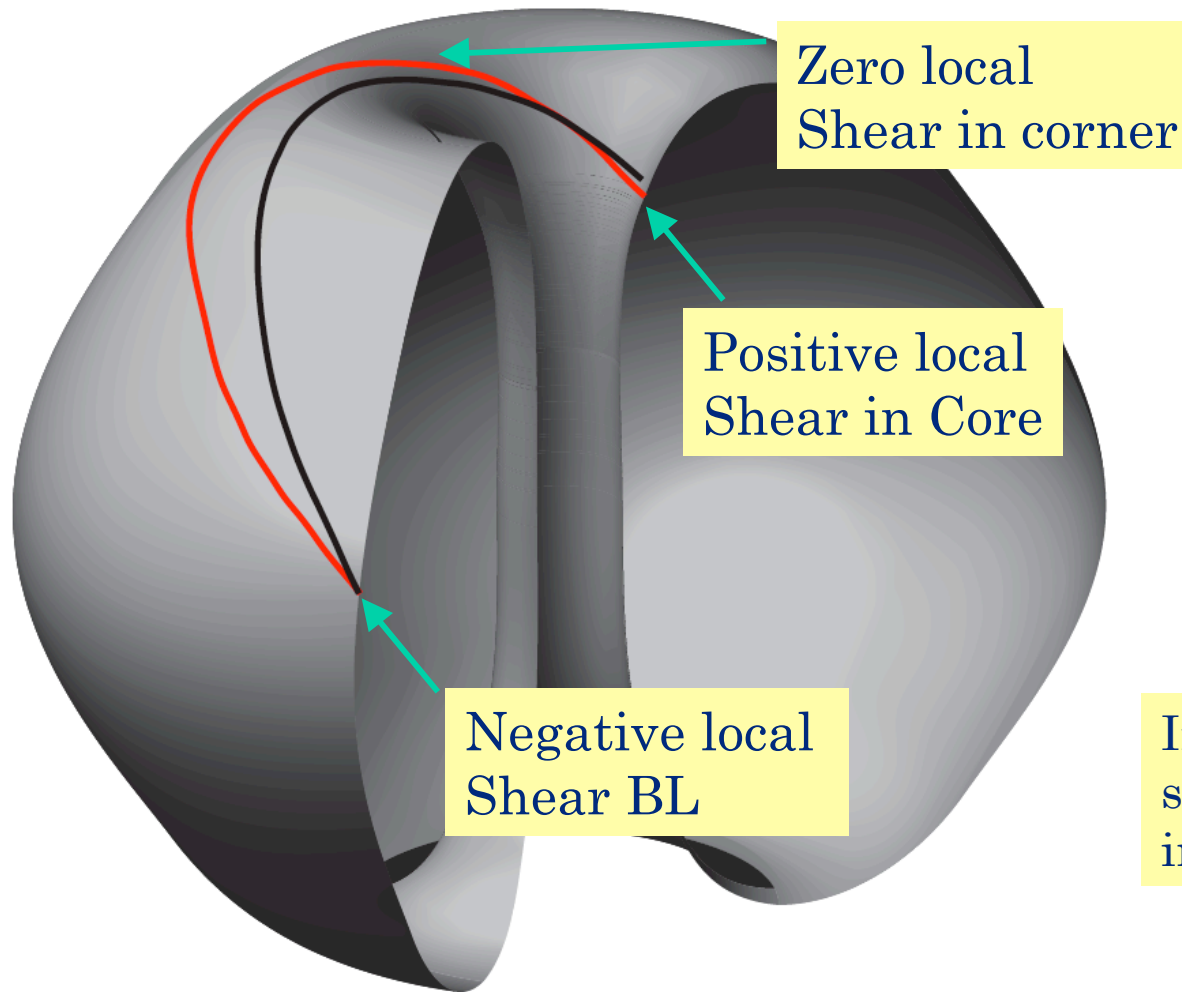
Stabilizing.



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Good properties.

5. Strong negative local shear in BL



Shear length in boundary layer =

$$L_s \sim \mathcal{O}(a/q\beta) \ll L_c$$

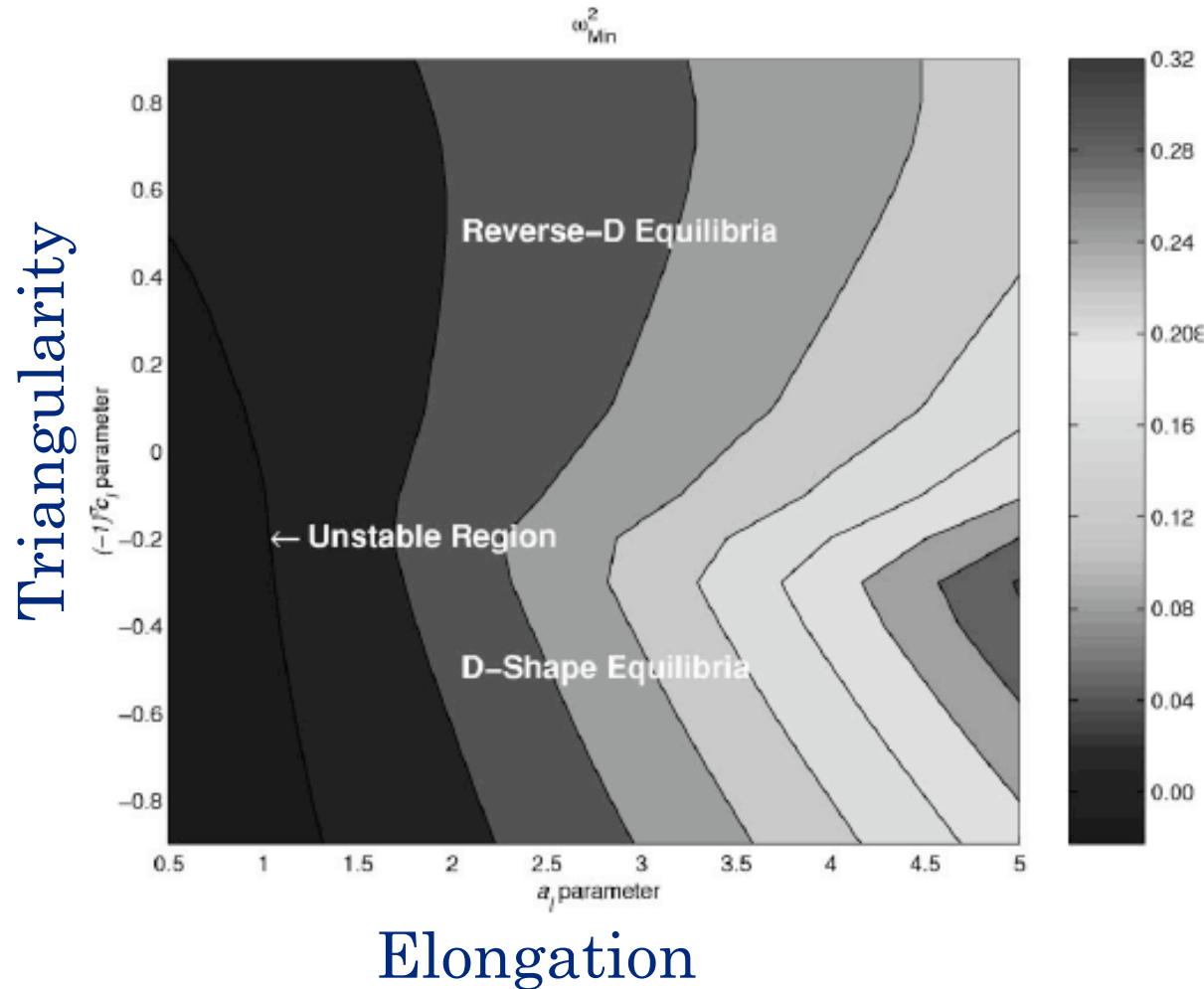
Instabilities are heavily sheared by Magnetic shear in BL Stabilizes ballooning.



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Internal Kink Stability

Analytic result just depends on shape -- not on β or q .



Hurricane et. al.
Phys. Plas. 2000

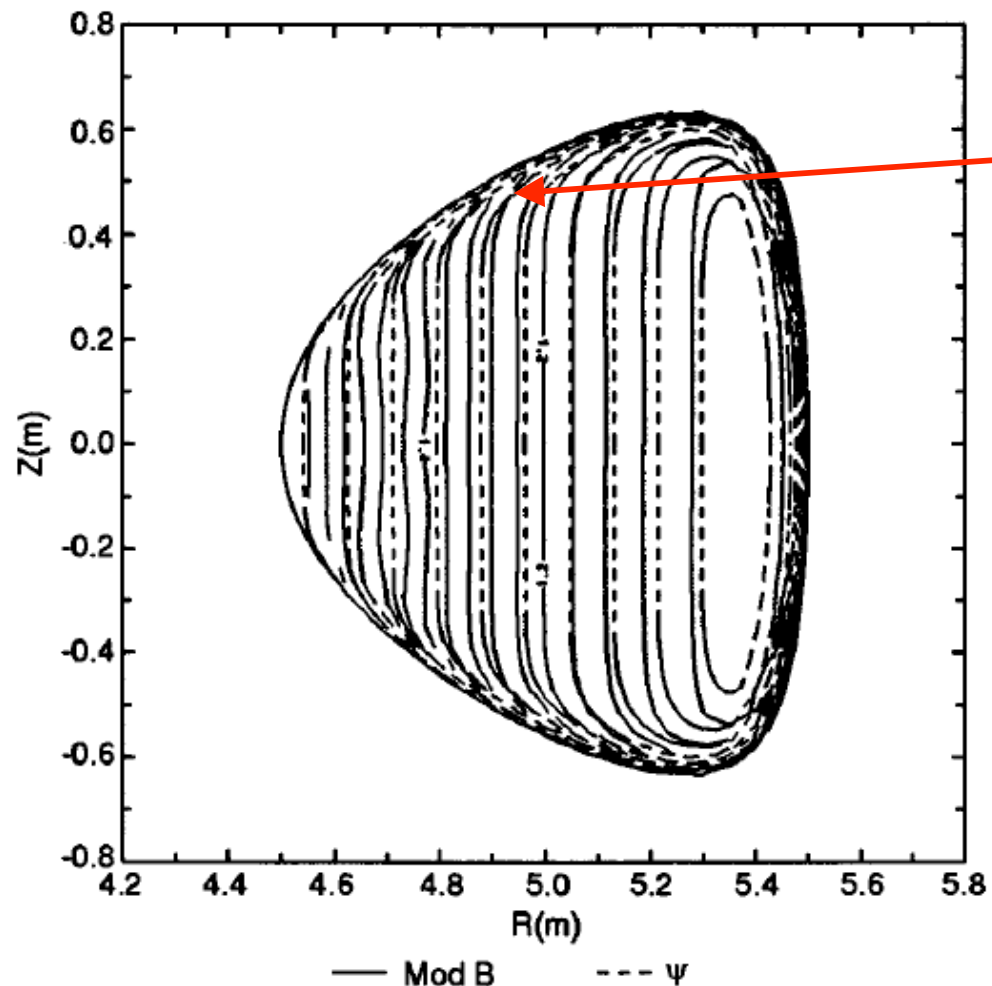
What about
Rotation?

External kinks
are the problem.



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Negative Triangularity - Reverse D.



Zero local Shear in
Corner firmly in good
Curvature region.

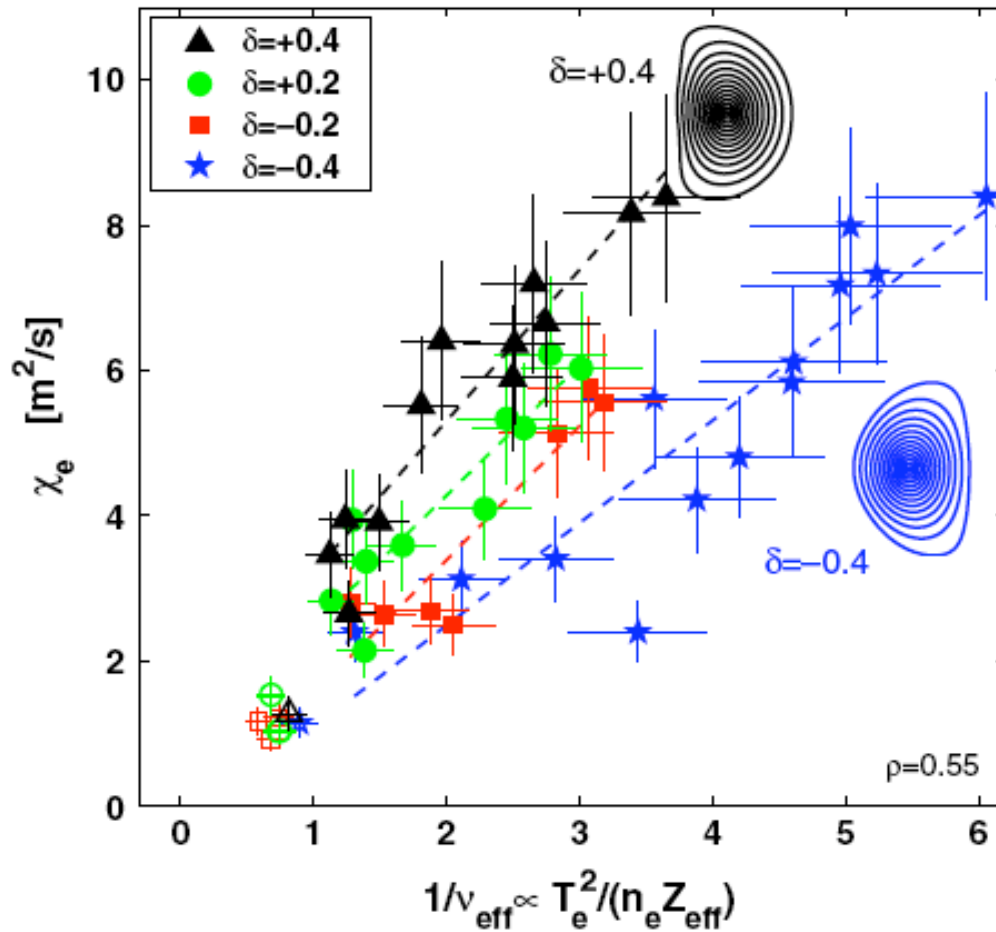
Particles drift reversed

Bill Dorland will show
Interesting results on
Similar configuration.



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Negative Triangularity - Reverse D. TCV



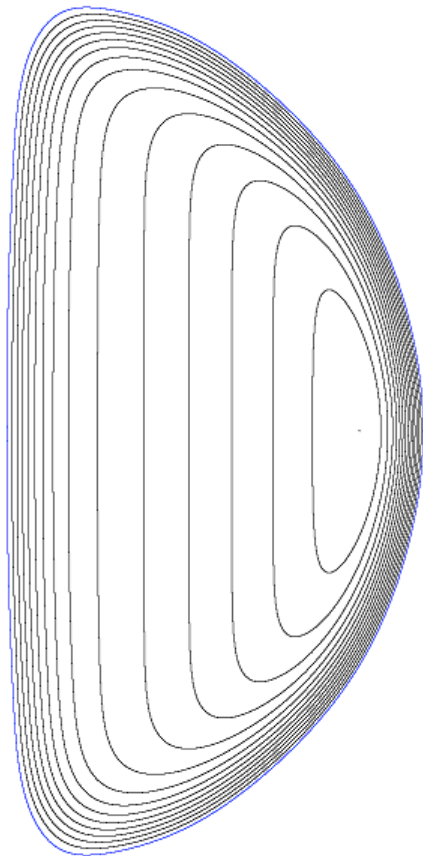
Less transport
In reverse D

Camenen et. al. Nucl Fus. 2007

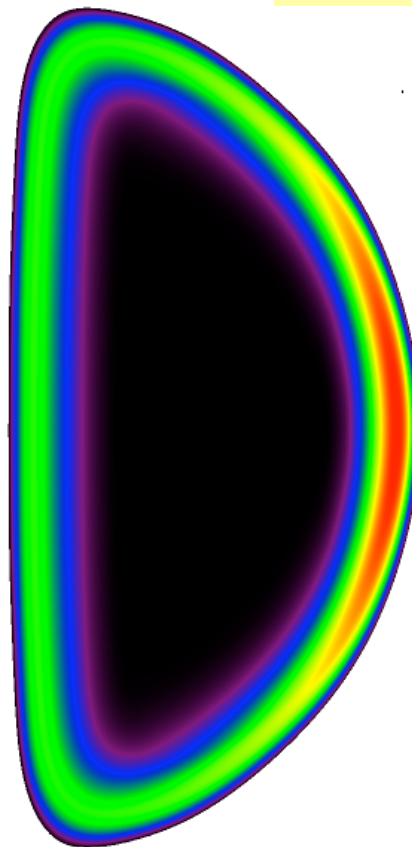
Unity beta current hole equilibrium

This equilibrium is stable to all ideal MHD criteria including internal and external modes for $n = 1, 2$ and $3 \dots$ *Note that the β_N is “small” despite the large value of beta.*

Pierre Gourdain's work



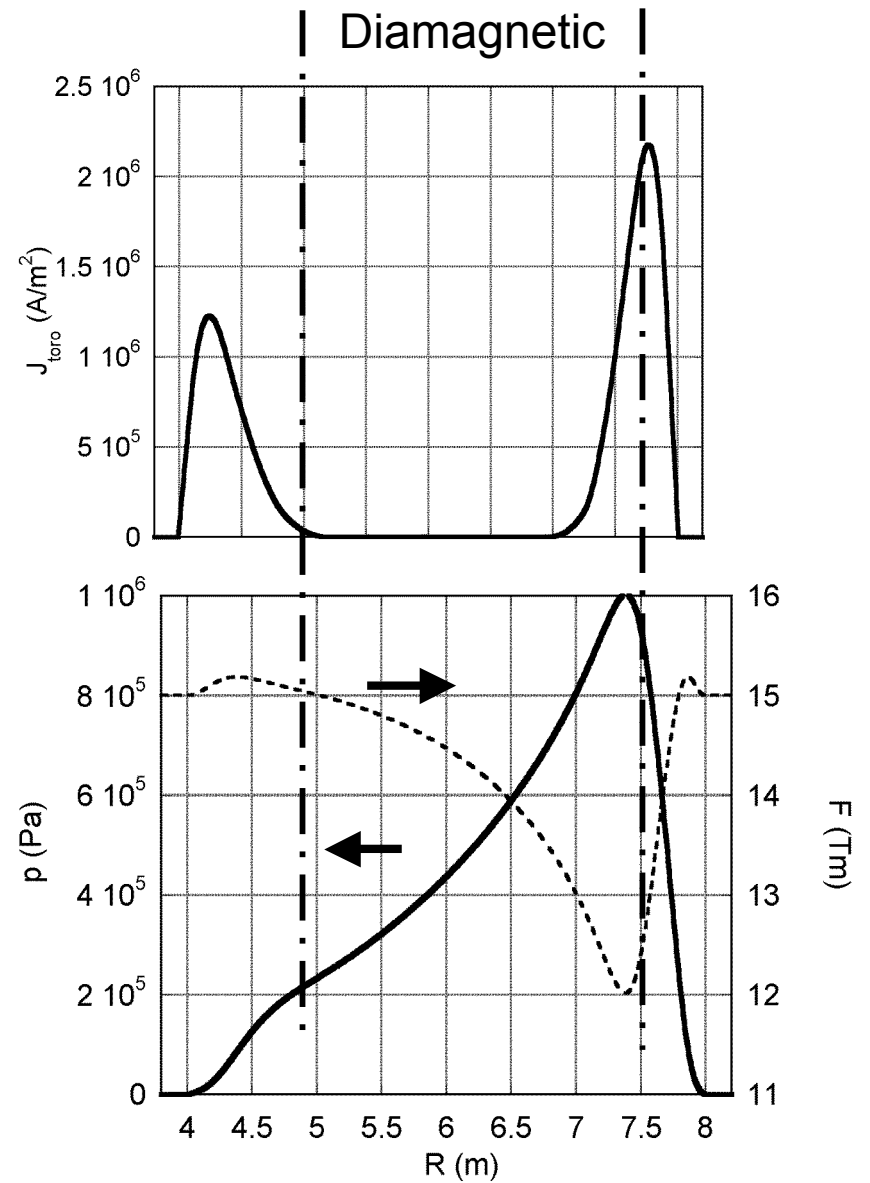
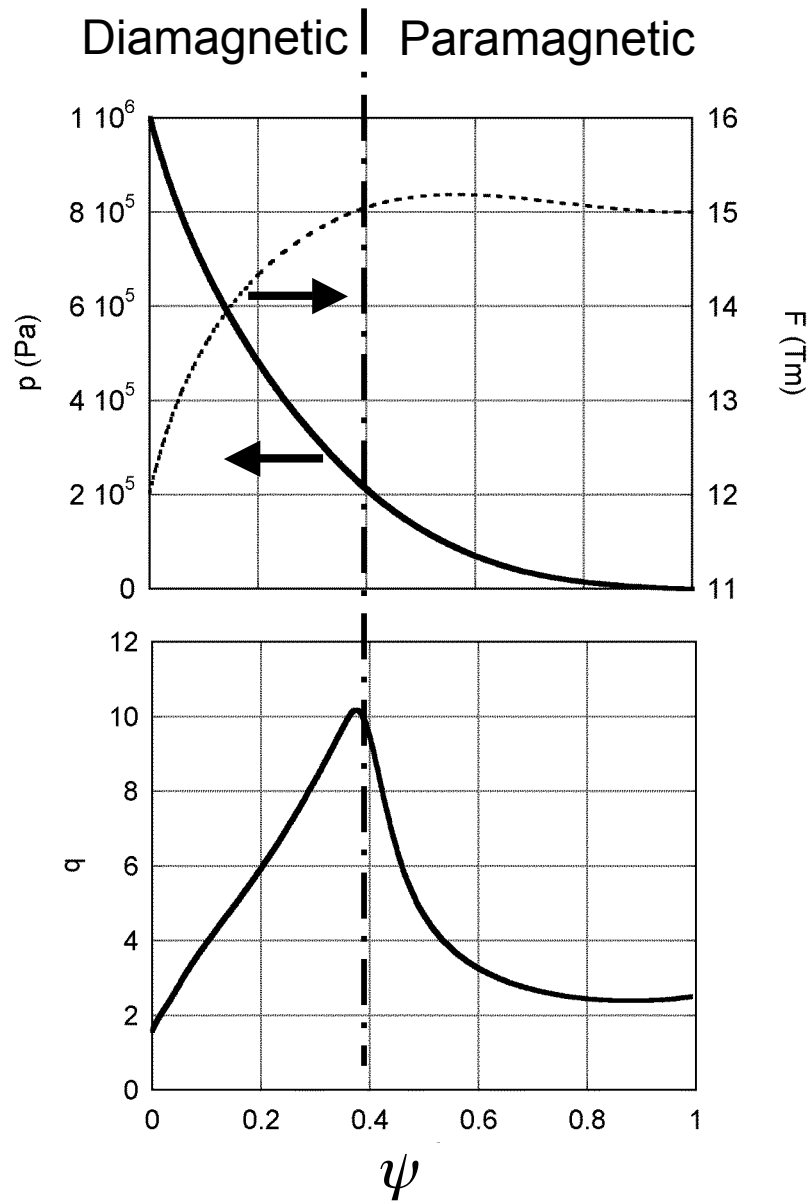
Flux Surfaces



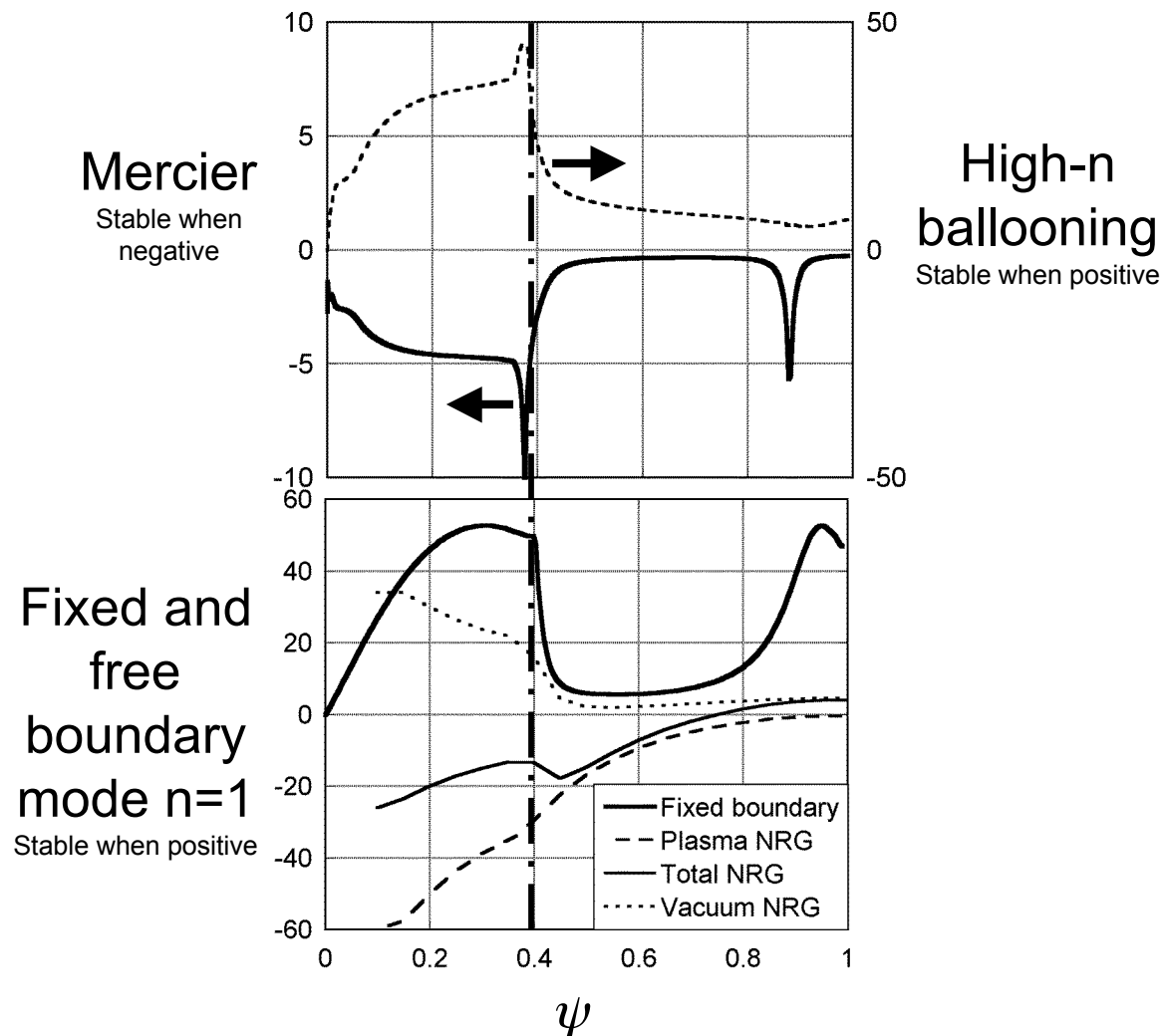
Toroidal Current Distribution

R	6 m
a	2 m
B_T	2.5 T
β	100%
$\langle \beta \rangle$	12%
β_N	4.6
q_{min}	1.5

Unity beta current hole profiles



Internal and external kink stability



DCON finds stability for Mercier, high-n ballooning as well as fixed boundary kink modes ($n=1$).

The free boundary mode $n=1$ is also stable (stability criteria obtained for $\psi = 1$).

Stability for $n=2$ and $n=3$ was also demonstrated.



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Better Bottle -- better shape?

- It certainly isn't clear that we can find a better bottle. But we should use our best tools to look hard.
- Start with transport considerations -- gyro-kinetics?