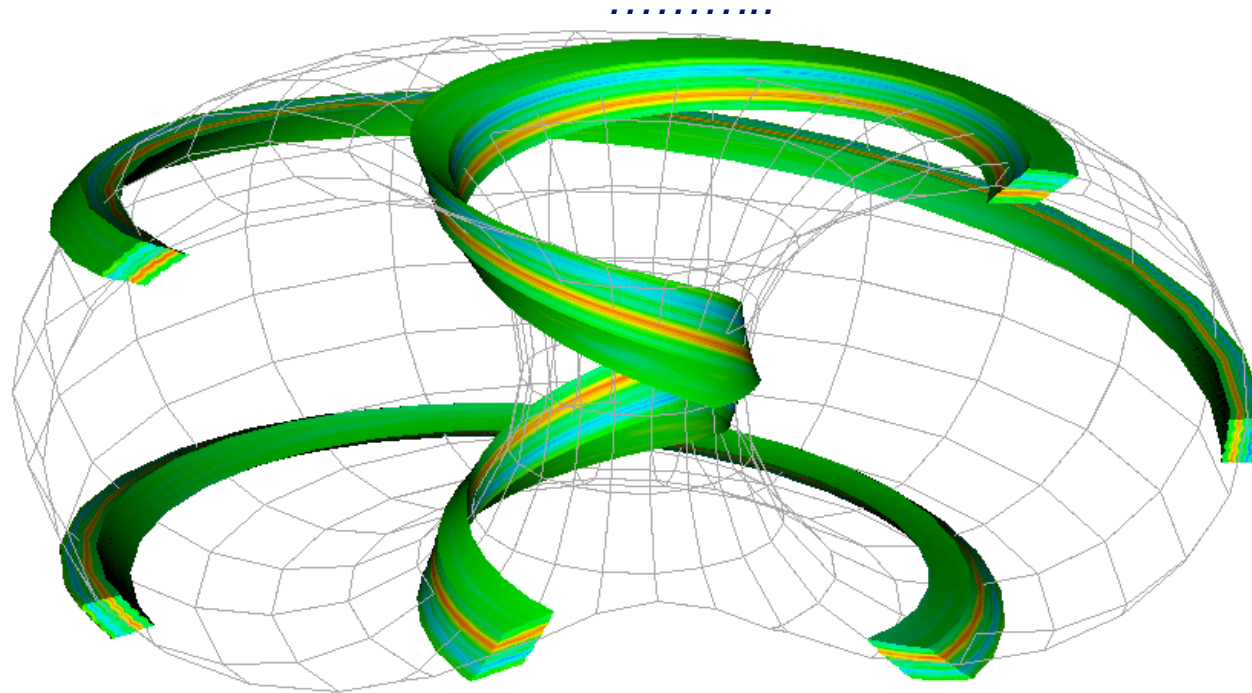


Multiscale Turbulence in Fusion and Gyrokinetics.

Steve Cowley -- Culham, Imperial

Gabe Plunk, Eric Wang, Ian Abel, Alex Schekochihin, Bill Dorland, Greg Hammett, Colin Roach, Michael Barnes, Felix Parra, Francis Casson



What am I going to say?.

- *The standard model of tokamak confinement and turbulence.*
- *Spatial scales -- time scales -- velocity space scales.*
- *Gyro-kinetic expansion.*
- *electron-ion separation.*
- *Possible problems*
 - nonlocality, time/space*
 - loss of scale separation*

ITER

First Sustained Burning Plasma. Starts in 2019.

BASIC PARAMETERS.

Plasma Major Radius 6.2m

Plasma Minor Radius 2.0m

Plasma Current 15.0MA

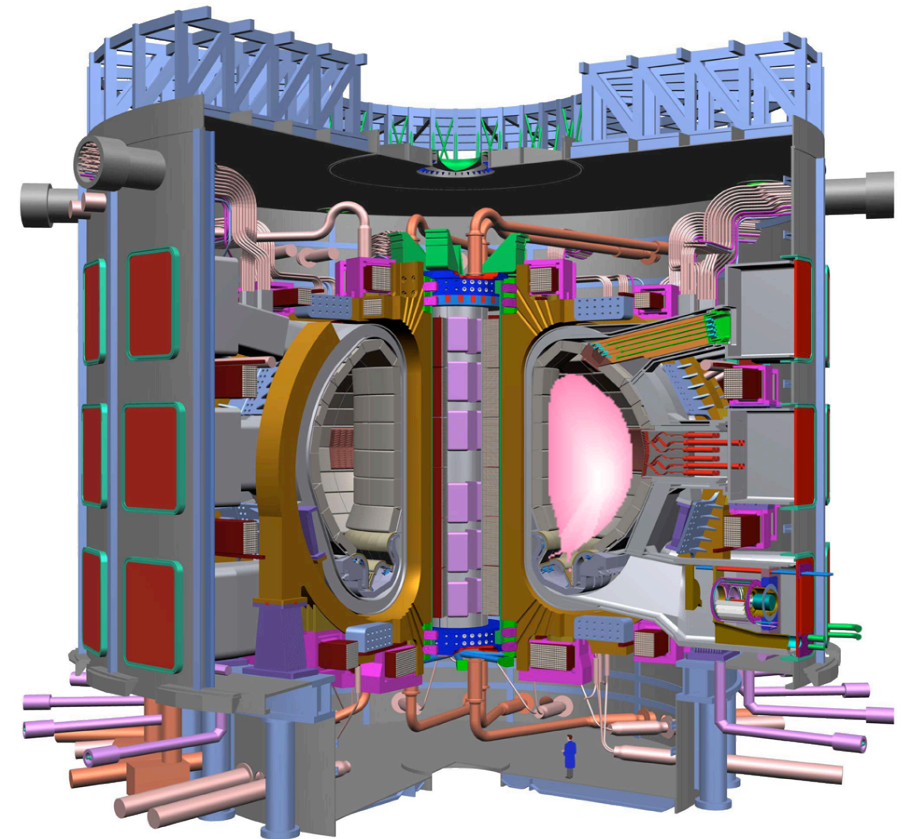
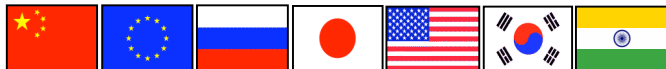
Toroidal Field on Axis 5.3T

Fusion Power 500MW

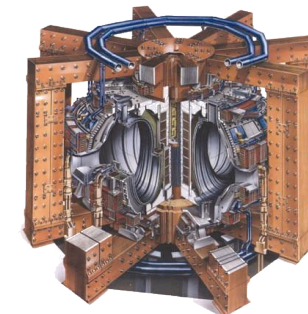
Burn Flat Top >400s

Power Amplification $Q > 10$

Cost is > 12 Billion Euro.

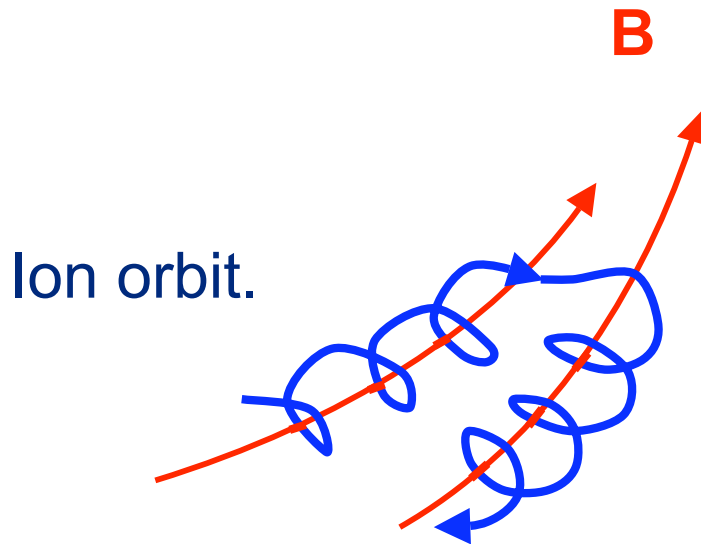


JET



Classical Transport.

Spitzer. 1951.



Random walk:

Step = ρ , larmor/cyclotron radius.

Decorrelation rate = ν = collision rate

Radius of plasma = a .

$$\tau_E \sim \frac{a^2}{\nu \rho^2}$$

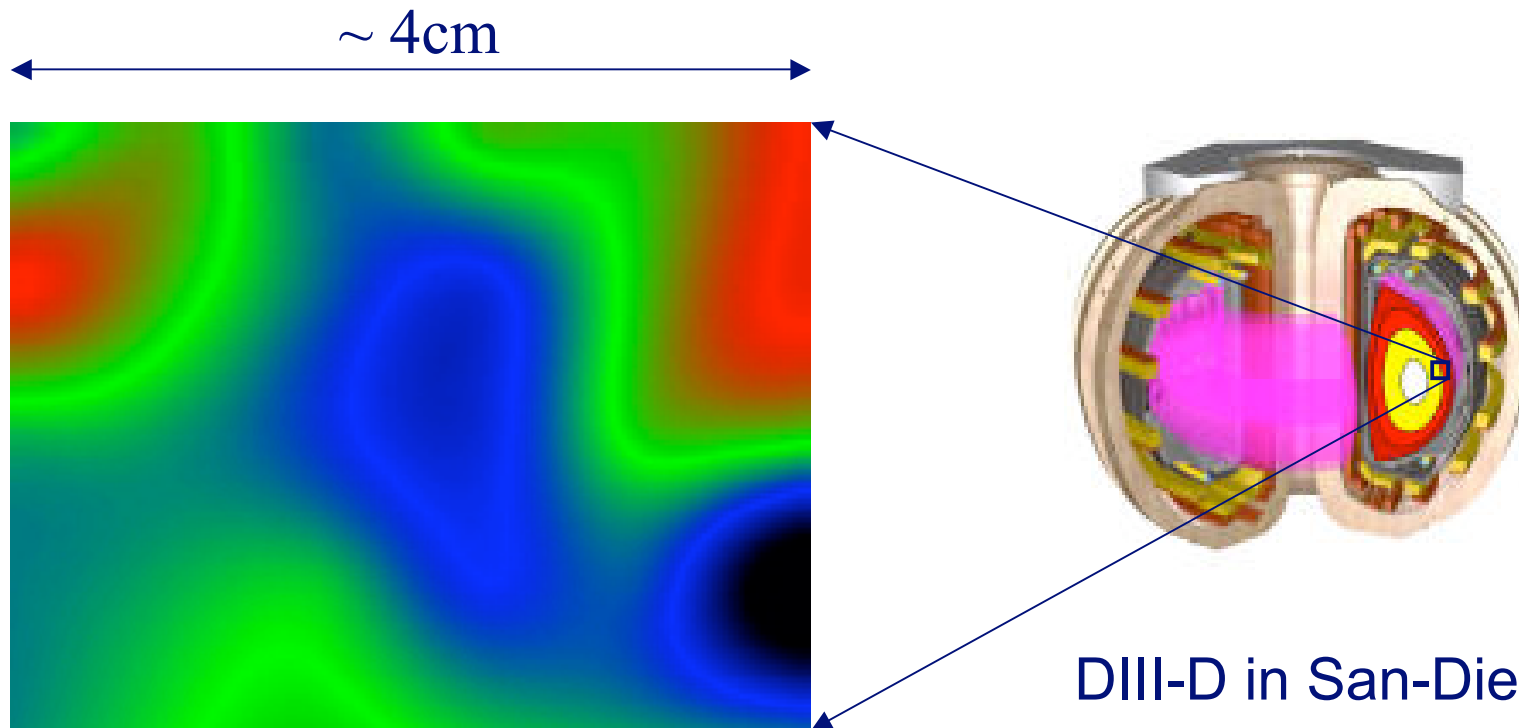
Collisions are rare and classical confinement can be very good.

Spitzer only needed $a = 20\text{cm}$, ($\tau_E > 4\text{s}$) for IGNITION.

Can't be right. Observed transport is much larger.

Turbulence Imaging, Beam Emission Spectroscopy

George McKee and Ray Fonck



Density fluctuations

DIII-D in San-Diego
Plasma is 1m across

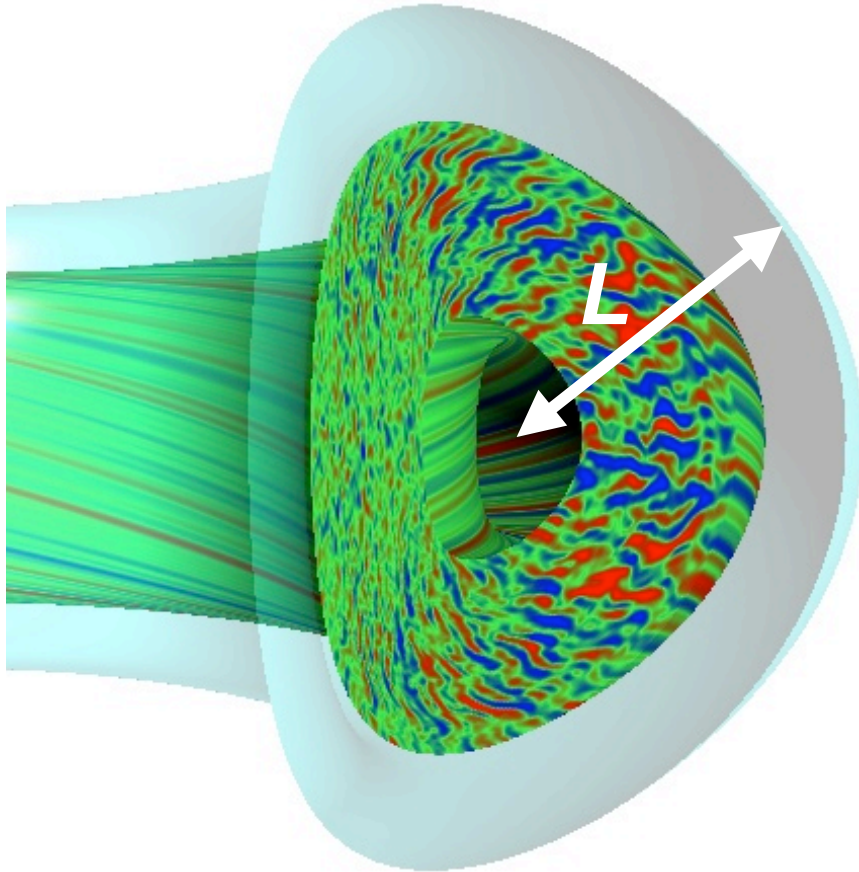
Eddies are small
compared to the device

Gyro-kinetic simulation.

DIII-D Shot 121717

GYRO Simulation
Cray XIE, 256 MSPs

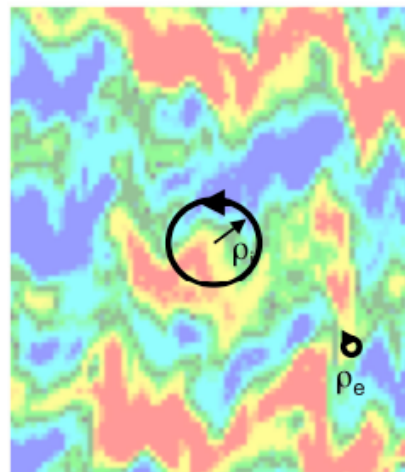
Spatial scales



L = Equilibrium scale and parallel scale of turbulence

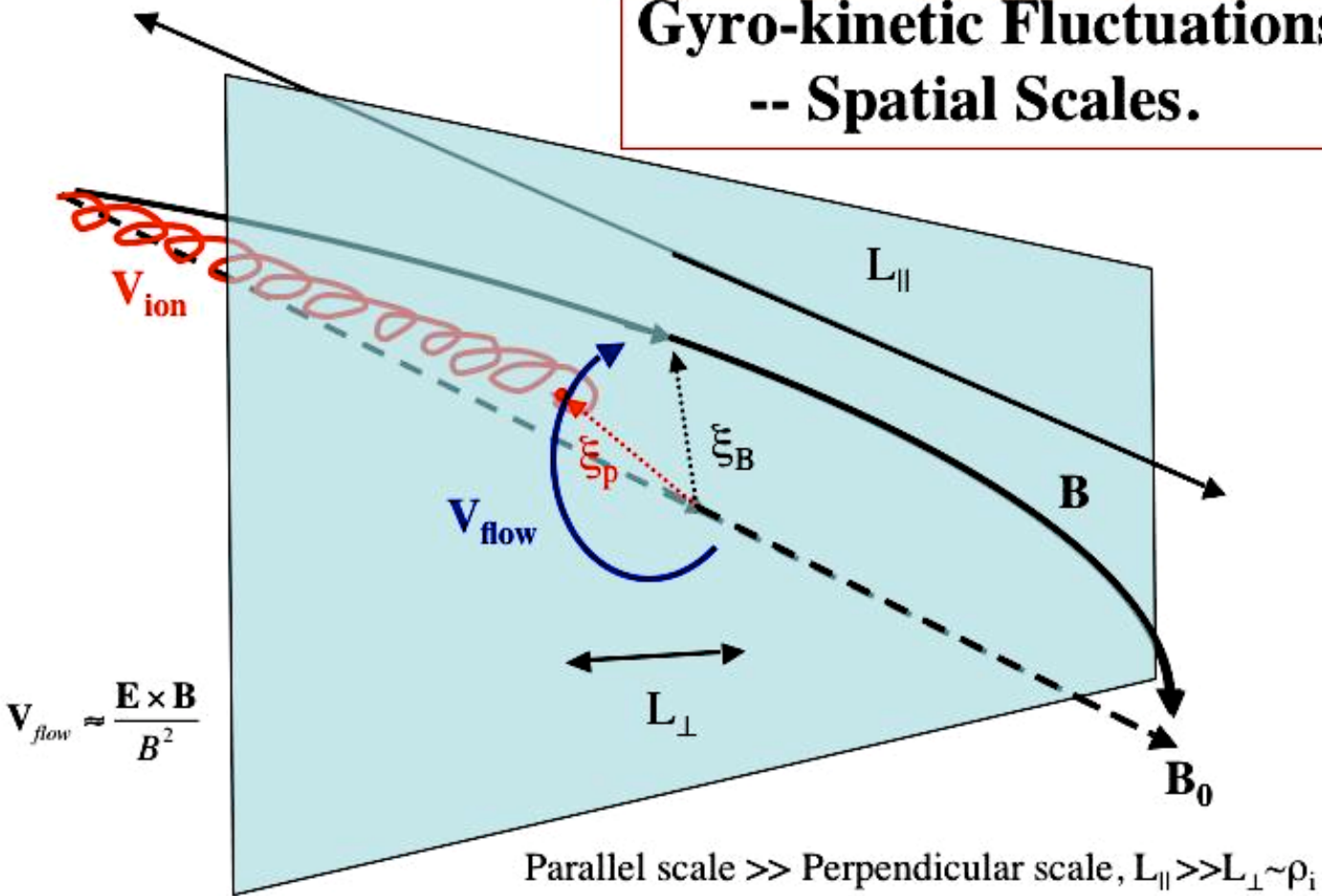
ρ = Ion larmor radius and perpendicular scale of turbulence

$\rho^* = \rho/L \sim 10^{-3}$ in ITER



Turbulent Scales

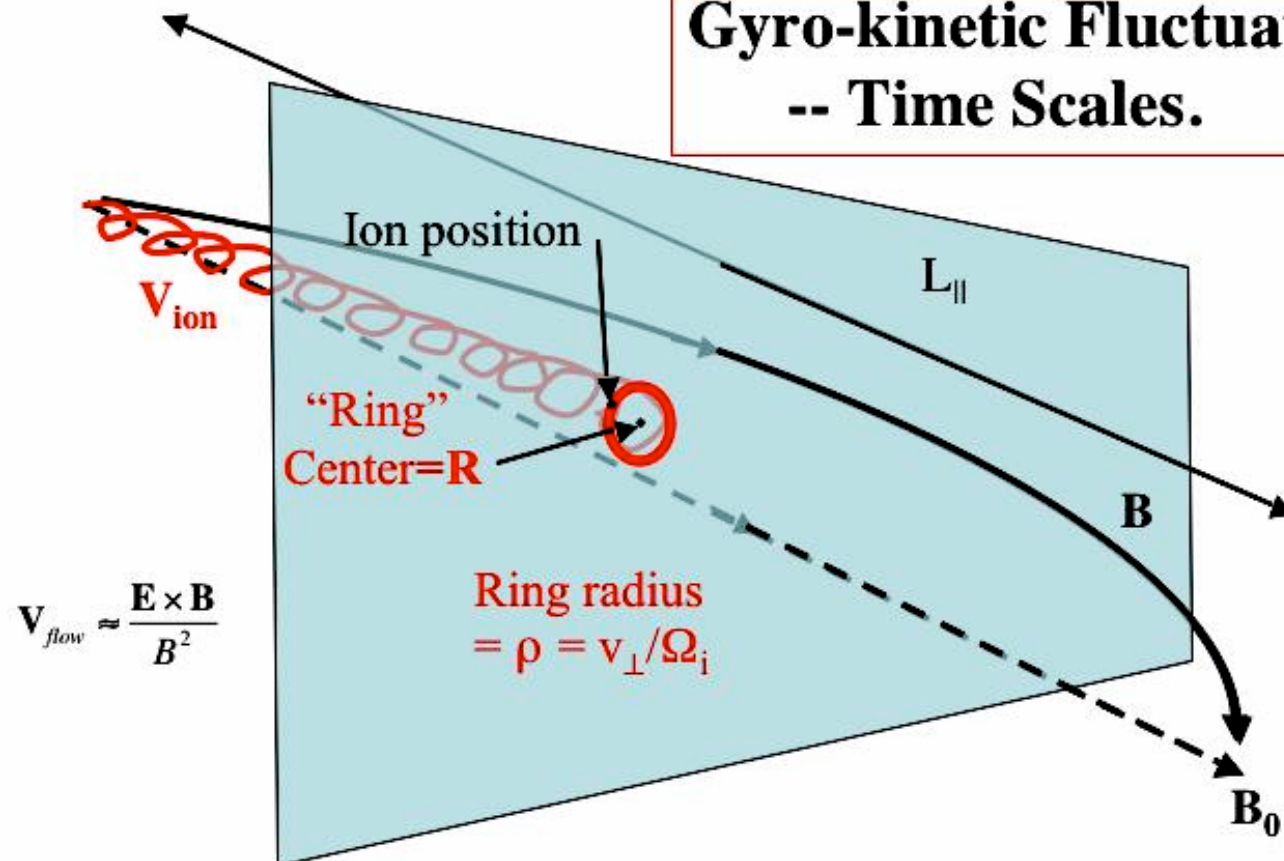
Gyro-kinetic Fluctuations -- Spatial Scales.



Particle displacement, $\xi_p \sim$ Of Order Field line displacement, $\xi_B \sim O(L_{\perp})$

Turbulent Scales

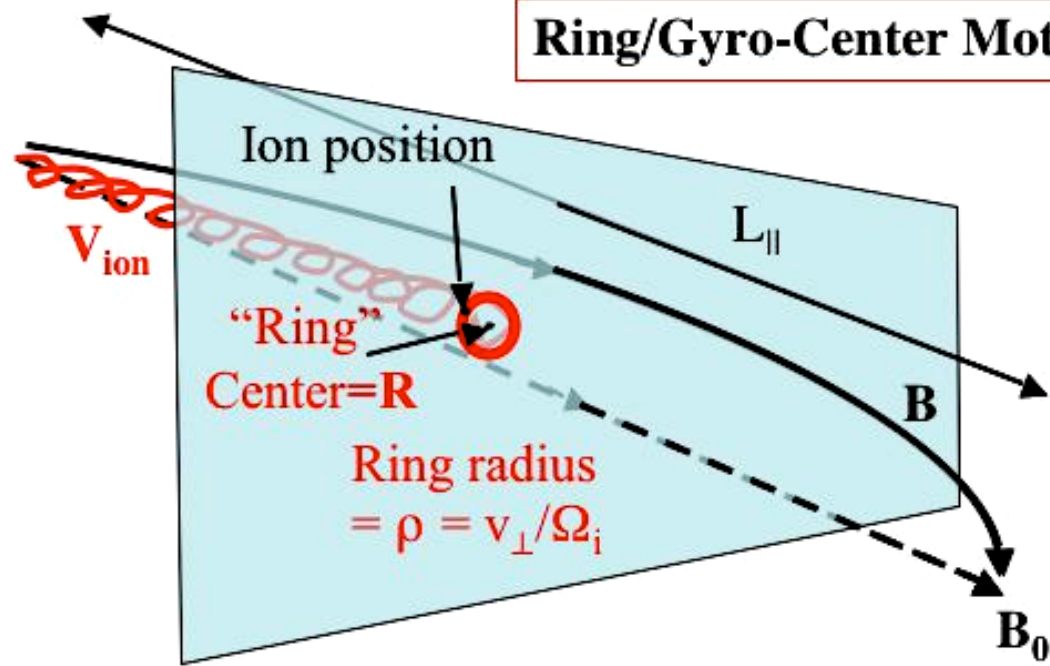
Gyro-kinetic Fluctuations -- Time Scales.



Typical time for fields to change is L_{\parallel} / v_{ion} -- i.e. the time for an ion to go one parallel wavelength. Much longer than the gyration period. \Rightarrow Ion senses the fields averaged over a **RING** of radius $\rho = v_{\perp} / \Omega_i$.

Guiding Centre Motion

Ring/Gyro-Center Motion



$$\left\langle \frac{d\mathbf{R}}{dt} \right\rangle = v_{\parallel} \mathbf{b}_0 - \frac{\partial \langle \chi \rangle_{\mathbf{R}}}{\partial \mathbf{R}} \times \left(\frac{\mathbf{b}_0}{B_0} \right) + v_{\parallel}^2 \left(\frac{\mathbf{b}_0}{\Omega_0} \right) \times \mathbf{b}_0 \cdot \nabla \mathbf{b}_0 + \frac{v_{\perp}^2}{2B_0} \left(\frac{\mathbf{b}_0}{\Omega_0} \right) \times \nabla B_0.$$

Parallel Motion along
Equilibrium Field

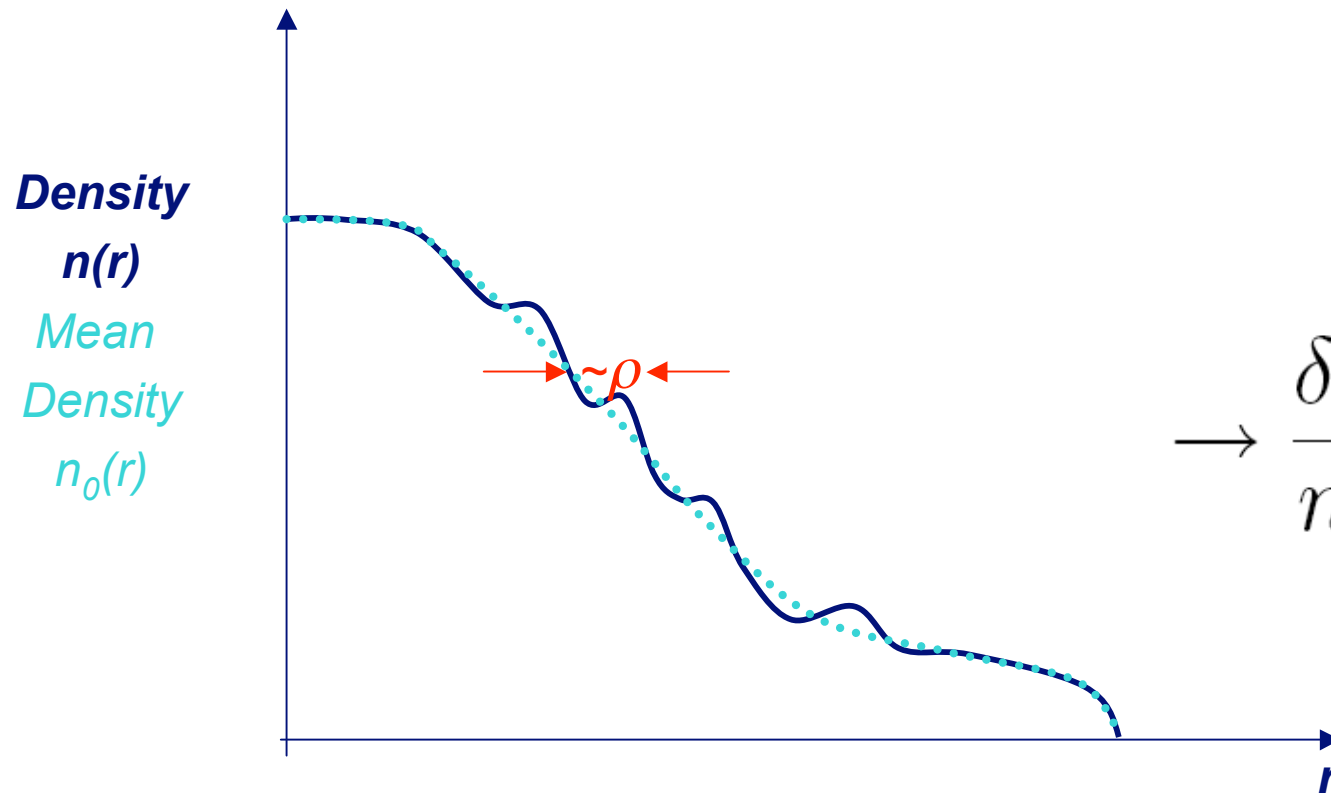
Drifts in
Perturbed Field

Curvature Drift in
Equilibrium Field

∇B Drift in
Equilibrium Field

$$\chi = \phi - \mathbf{v} \cdot \mathbf{A}$$

Ambient Gradient Argument

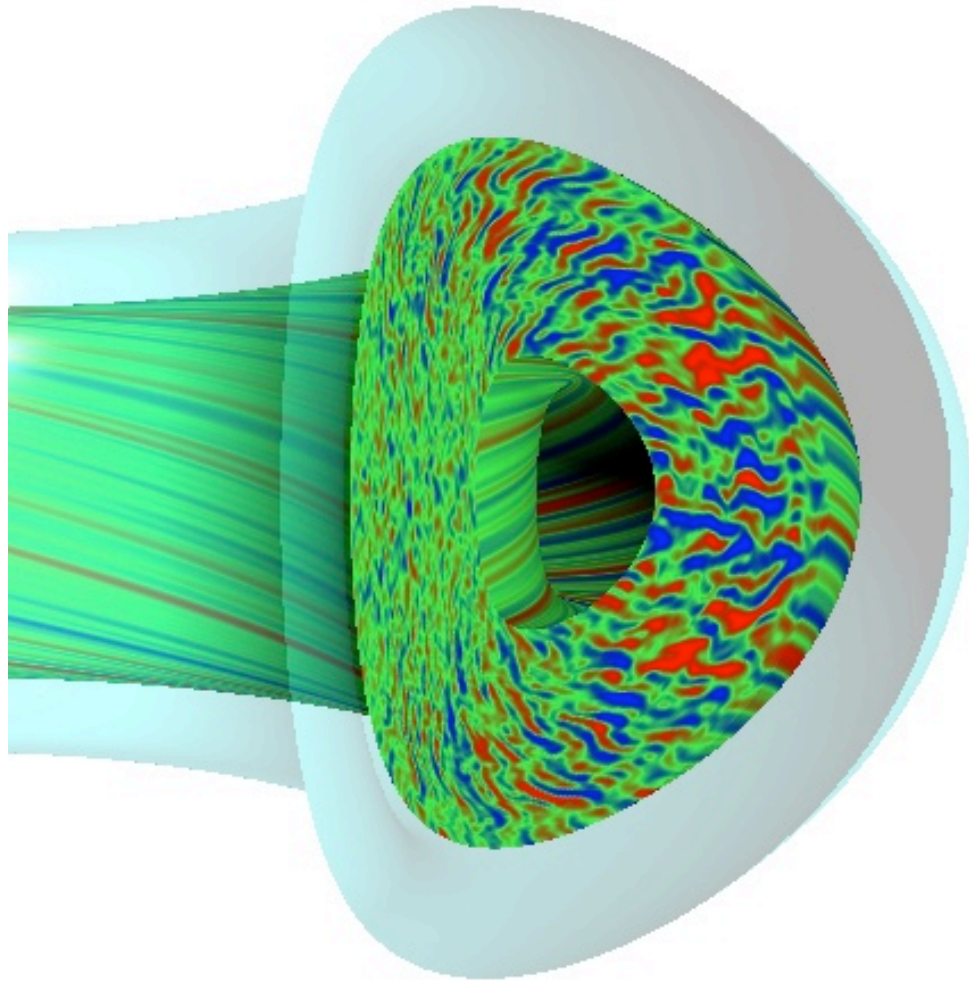


$$\begin{aligned} \nabla \delta n &\sim \nabla n_0 \\ \frac{\delta n}{\rho} &\sim \frac{n_0}{L} \\ \rightarrow \frac{\delta n}{n_0} &\sim \frac{\rho}{L} = \rho^* \end{aligned}$$

WARNING without collisions $\frac{df}{dt} = 0$ \Rightarrow

Collisionless mixing produces Finite fluctuations -- collisions are essential.

Energy Confinement -- Random walk of heat/particles.



L = typical machine size

Δ = radial eddy size \propto Ion larmor

Radius ρ_i = random step.

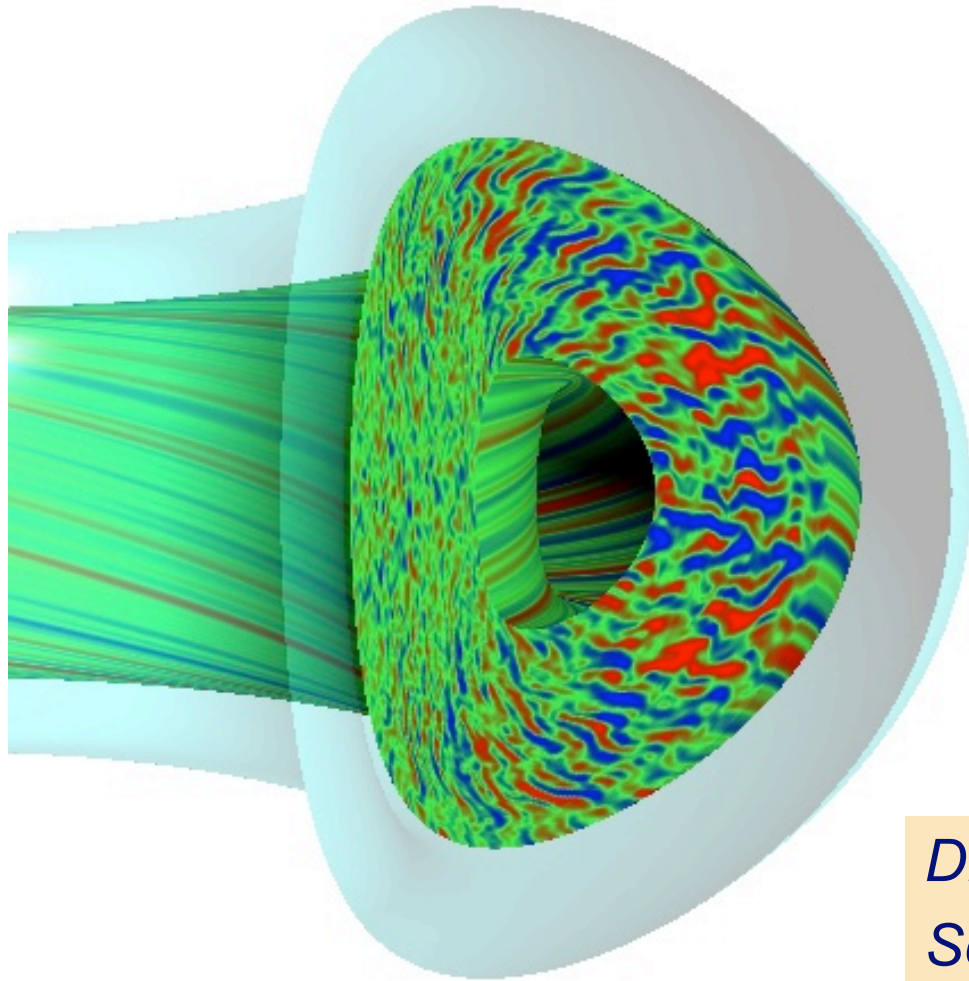
N = number of steps to
random walk out of plasma

$$L \sim \sqrt{N} \rho_i$$

$$\rightarrow N = \left(\frac{L}{\rho_i}\right)^2 = \left(\frac{1}{\rho^*}\right)^2$$

For ITER $N \sim 10^6$.

Energy Confinement -- Random walk of heat/particles.



Eddy turnover time =

$$\tau_{\text{eddy}} = \left(\frac{L}{v_{thi}} \right)$$

$$\tau_E \sim N \tau_{\text{eddy}} \sim \left(\frac{L^3}{\rho_i^2 v_{thi}} \right)$$

$$\propto L^3 B^2 T^{-1}$$

Dramatic scaling with size!

Scaling approximately agrees with data BUT geometry dependant.

Timescales -- ITER numbers -- ordering

Cyclotron -- $\sim 4 \times 10^{-9} \text{s}$ (ions) -- no activity

$$\frac{1}{\Omega_i}$$

Turbulence -- $\sim 10^{-5} \text{s}$ -- fluctuating density, temperature, f , \mathbf{E} .

Collisions -- $\sim 10^{-2} - 10^{-3}$ -- reestablish Maxwellian

$$\left(\frac{1}{\rho^*}\right) \frac{1}{\Omega_i}$$

Transport -- $\sim 4 \text{s}$ -- change mean density, temperature

$$\left(\frac{1}{\rho^*}\right)^3 \frac{1}{\Omega_i}$$

Expand

$$f(\mathbf{r}, \mathbf{v}, t) = \underbrace{F_0(\mathbf{r}, \mathbf{v}, t)}_1 + \underbrace{\frac{q\phi}{T} F_0 + h(\mathbf{R}, \mathcal{E}, \mu, \mathbf{t})}_{\rho^*} + \underbrace{\delta f_2(\mathbf{r}, \mathbf{v}, t)}_{(\rho^*)^2} + \dots$$

*Varies on slow space
and timescales*

*Varies on faster turbulence
space and timescales*

$$\mathbf{B}(\mathbf{r}, t) = \mathbf{B}_0(\mathbf{r}, t) + \delta\mathbf{B}(\mathbf{r}, t), \quad \mathbf{E}(\mathbf{r}, t) = \delta\mathbf{E}(\mathbf{r}, t)$$

At $\mathcal{O}(1)$ collisions establish a local Maxwellian.

$$F_0(\mathbf{r}, \mathbf{v}, t) = n(t, \psi) \left(\frac{m}{2\pi T(t, \psi)} \right)^{3/2} \exp \left[- \left(\frac{(1/2)mv^2}{T(t, \psi)} \right) \right]$$

Key issue is to find the slow evolution of $n(t, \psi)$ and $T(t, \psi)$

Solving for Fine Scale Turbulence

*Evolution of turbulence comes from solving the gyro-kinetic equation:
e.g. the electrostatic version*

$$\frac{\partial h}{\partial t} + v_{\parallel} \frac{\partial h}{\partial Z} + \mathbf{v}_{\perp} \cdot \frac{\partial h}{\partial \mathbf{R}} - \frac{\partial \langle \phi \rangle_{\mathbf{R}}}{\partial \mathbf{R}} \times \left(\frac{\mathbf{b}_0}{B_0} \right) \cdot \frac{\partial h}{\partial \mathbf{R}} - \langle \tilde{C}(h) \rangle_{\mathbf{R}} = q \frac{F_0}{T_0} \frac{\partial \langle \phi \rangle_{\mathbf{R}}}{\partial t} - \frac{\partial \langle \phi \rangle_{\mathbf{R}}}{\partial \mathbf{R}} \times \left(\frac{\mathbf{b}_0}{B_0} \right) \cdot \frac{\partial F_0}{\partial \mathbf{R}}$$

Comes from ring averaged kinetic equation, $\mathcal{O}(\rho^)$*

Solved with quasi neutral approximation:

$$\cancel{\nabla^2} \phi = -\frac{1}{\epsilon_0} (qn_i - en_e)$$

In principle this equation should be solved everywhere

With F_0 held fixed in time letting h converge to statistical equilibrium.

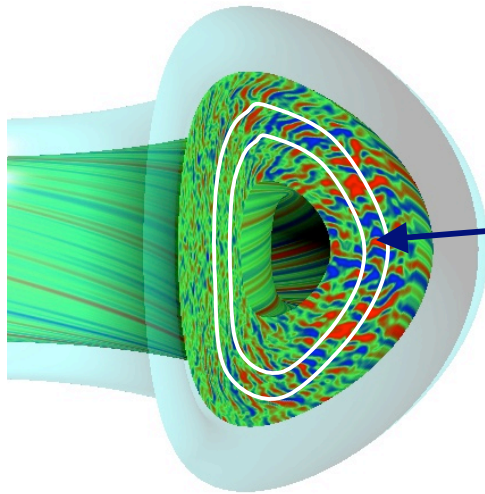
th weak collisions h becomes very fine scale in velocity space due to mixing

Solving for Slow Evolution

$$\mathcal{O}((\rho^*)^2) \quad \frac{\partial F_0}{\partial t} + \frac{\partial \delta f_2}{\partial t} + \dots$$

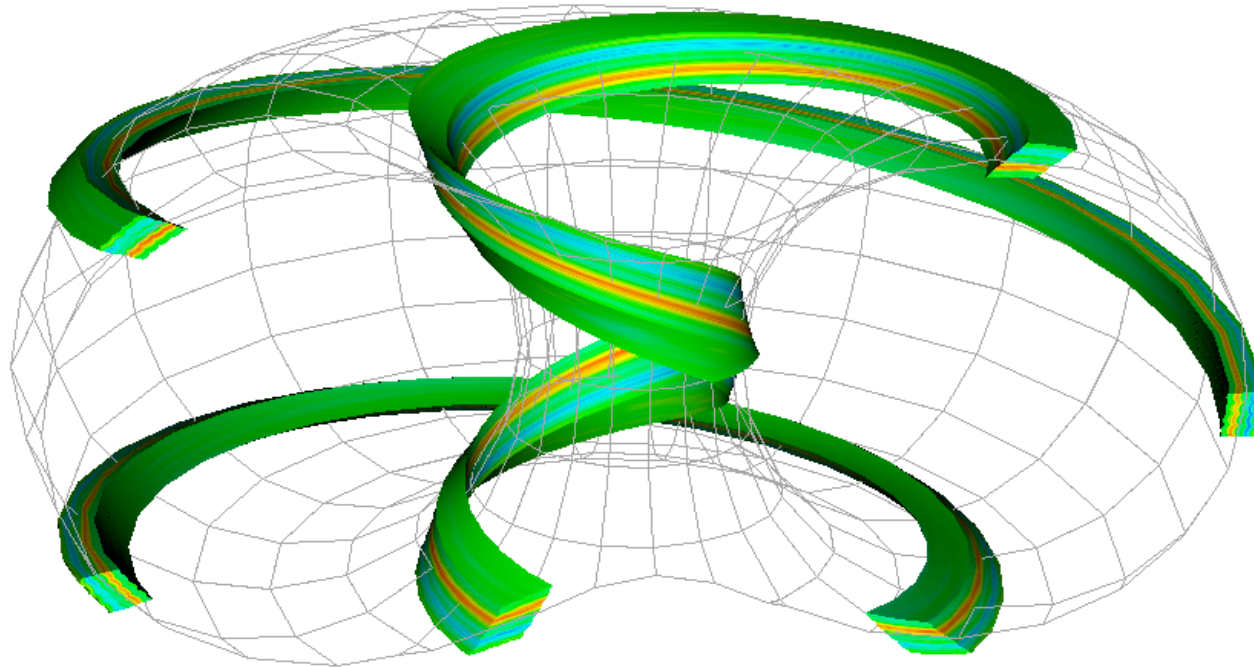
To solve for time evolution of F_0 (n & T) we have to annihilate the fast varying parts in space and time. It is easier to first take moments and use exact moment equations.

$$\frac{\partial n_i}{\partial t} + \nabla \cdot \Gamma_i = 0. \quad \Gamma_i = \int \mathbf{v} f_i d^3 \mathbf{v}. \quad \begin{array}{l} \text{Flux comes} \\ \text{From turbulence.} \end{array}$$



*Average flux over annulus (of thickness much greater than turbulence scale much smaller than radius) and time
To get mean evolution of n and T .*

Local simulation -- flux tubes



Beer 1993.

Is the turbulence in the flux tube determined by the conditions in the tube?

Does turbulence propagate -- correlated or not?

We should look at turbulent Greens function -- response to stirring.

What should we do?

- *Transport barriers -- can we use the “standard model”.*
- *Detailed check of “standard model” needed.*
- *Full coupling -- (e.g. Trinity) needs finishing.*
- *The big prize is to find better configurations.*