ASTROPARTICLE



Hilary 2021



PHYSICS

Axford hysics

Oxford Master Course in Mathematical and Theoretical Physics

The universe observed Relativistic world models **Reconstructing the thermal history Big bang nucleosynthesis** Dark matter: astrophysical observations Dark matter: relic particles Dark matter: direct detection Dark matter: indirect detection Cosmic rays in the Galaxy \diamond ♦ Antimatter in cosmic rays Ultrahigh energy cosmic rays High energy cosmic neutrinos The early universe: constraints on new physics The early universe: baryo/leptogenesis The early universe: inflation & the primordial density perturbation Cosmic microwave background & large-scale structure

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http://www-thphys.physics.ox.ac.uk/user/SubirSarkar/astropartphys.html

Mass scale	Particle	Symmetry/ Quantum #	Stability	Production	Abundance
Λ_{QCD}	Nucleons	Baryon number	$\tau > 10^{33} \text{ yr}$ (dim-6 OK)	'freeze-out' from thermal equilibrium	$\Omega_{\rm B} \sim 10^{-10}$ cf. observed $\Omega_{\rm B} \sim 0.05$

We have a good theory for why baryons are massive and (cosmologically) stable



However, in the standard cosmology ~none should be left-over from the Big Bang!

WHAT IS THE EXPECTED RELIC ABUNDANCE OF BARYONS?



 $n_B + n_{\bar{B}}$

The Standard $SU(3)_c \times SU(2)_L \times U(1)_\gamma$ Model provides an exact description of all microphysics (up to some high energy cut-off M)

Higgs mass divergence $+M^{4} + M^{2} \Phi^{2} m_{H}^{2} \simeq \frac{h_{t}^{2}}{16\pi^{2}} \int_{0}^{M^{2}} dk^{2} = \frac{h_{t}^{2}}{16\pi^{2}} M^{2} \qquad \text{super-renormalisable}$ $\mathcal{L}_{eff} = F^{2} + \bar{\Psi} \not{D} \Psi + \bar{\Psi} \Psi \Phi + (D\Phi)^{2} + V(\Phi) \qquad \text{renormalisable}$ $-\mu^{2} \phi^{\dagger} \phi + \frac{\lambda}{4} (\phi^{\dagger} \phi)^{2}, m_{H}^{2} = \lambda v^{2}/2 \qquad -\mu^{2} \phi^{\dagger} \phi + \frac{\lambda}{4} (\phi^{\dagger} \phi)^{2}, m_{H}^{2} = \lambda v^{2}/2 \qquad \text{non-renormalisable}$

The effect of new physics beyond the SM (neutrino mass, nucleon decay, FCNC) \Rightarrow **non-renormalisable operators** suppressed by M^n ... which 'decouple' as $M \rightarrow M_P$

But as *M* is raised, the effects of the **super-renormalisable operators** are exacerbated

One solution for 2^{nd} term \rightarrow 'softly broken' supersymmetry at $M \sim 1$ TeV

This suggests possible mechanisms for **baryogenesis**, candidates for **dark matter**, ... (as also do other proposed extensions of the SM, e.g. new dimensions @ TeV scale)

For example, the lightest supersymmetric particle (typically the neutralino χ), *if* protected against decay by *R*-parity, is a candidate for thermal dark matter

But if the Higgs is composite (as in **technicolour** models of $SU(2)_L \times U(1)_Y$ breaking) then there is *no* need for supersymmetry ... and light TC states can be dark matter THERMAL RELICS

$$\dot{n} + 3Hn = -\langle \sigma v \rangle (n^2 - n_{\rm T}^2)$$

Chemical equilibrium is maintained as long as the annihilation rate exceeds the Hubble expansion rate

- 'Freeze-out' can occur either when the annihilating particles are:
- \blacktriangleright Relativistic: $n \sim n_{oldsymbol{\gamma}}$
- > Non-relativistic: $n \sim n_{\gamma} \mathrm{e}^{-m/2}$



Example 1 : $\sum \Omega_{\nu} h^2 \simeq m_{\nu_i} / 93 \text{eV}$

But how might this mass scale arise?

➤ natural for Fermi scale mass/coupling

Example 2 :
$$\Omega_{\chi} h^2 \simeq \frac{3 \times 10^{-27} \text{cm}^3 \text{s}^{-1}}{\langle \sigma_{\text{ann}} v \rangle_{T=T_{\text{f}}}}$$

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$\Lambda_{ m Fermi} \sim \ G_{ m F}^{-1/2}$	Neutralino?	R-parity?	violated?	'freeze-out' from thermal equilibrium	$\Omega_{\rm LSP} \sim 0.3$
$H = \bigcup_{i}^{t} H = \bigcup_{H}^{t} H$ $H = \bigcup_{i}^{t} H$ $H = \bigcup_{H}^{t} H$			U d V ₃ C Quarks	Standard particles	SUSY particles

For (softly broken) **supersymmetry** we have the 'WIMP miracle':

$$\Omega_{\chi}h^2 \simeq \frac{3 \times 10^{-27} \mathrm{cm}^{-3} \mathrm{s}^{-1}}{\langle \sigma_{\mathrm{ann}} v \rangle_{T=T_{\mathrm{f}}}} \simeq 0.1, \text{ since } \langle \sigma_{\mathrm{ann}} v \rangle \sim \frac{g_{\chi}^4}{16\pi^2 m_{\chi}^2} \approx 3 \times 10^{-26} \mathrm{cm}^3 \mathrm{s}^{-1}$$

But why should a *thermal* relic have an abundance comparable to *non-thermal* baryons?

The more we fail to find SUSY particles, the *higher* their relic abundance is expected to be!



Figure 4

Regions of minimal supergravity (m_0 , $M_{1/2}$) parameter space for fixed $A_0 = 0$, tan $\beta = 10$, and $\mu > 0$. The green (dark yellow) region is cosmologically favored with $0.20 < \Omega_{\chi} < 0.28$ ($0.2 < \Omega_{\chi} < 0.6$). The names of cosmologically favored regions (focus point, bulk, and coannihilation) are indicated, along with regions with too much and too little dark matter. The lower right red shaded region is excluded by collider bounds on chargino masses; the upper left red region is excluded by the presence of a stable charged particle. Contours are for neutralino dark matter mass m_{χ} in gigaelectronvolts. Adapted from Feng, Matchev & Wilczek (2001).



'Natural' parameter space in the CMSSM

Heavy sparticles \rightarrow fine tuning of terms ... with measure: $\Delta(a_i) = \left| \frac{a_i}{M_z} \frac{\partial M_z}{\partial a_i} \right|$ 1000 Castell, Ghilencea & Ross, Nucl.Phys. B835:110,2010 Relic density unrestricted 500 SUSY particle masses 200 $3.20 < 10^4~{\rm Br}(b \rightarrow s \gamma) < 3.84$ 100 Δ $Br(b \to \mu\mu) < 1.8 \times 10^{-8}$ 50 $\delta a_{\mu} < 292 \times 10^{-11}$ 20 10 $-0.0007 < \delta \rho < 0.0012$ $\Delta_{Min} = 9, \quad m_h = 114 \pm 2GeV$ 80 12090 100110 1000 Relic density restricted 500 h^0 resonant annihilation 200 \tilde{h} t-channel exchange 100 2 Δ 2 50 $\tilde{\tau}$ co-annihilation 3 20 \tilde{t} co-annihilation 10 • 5 A^0 / H^0 resonant annihilation 110 120 80 90 100 $< 3\sigma$ WMAP: $\Delta_{Min} = 18$, $m_h = 115.9 \pm 2GeV$ Higgs Mass /GeV

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$\Lambda_{ m Fermi} \sim G_{ m F}^{-1/2}$	Neutralino?	R-parity?	violated?	'freeze-out' from thermal equilibrium	$\Omega_{\rm LSP} \sim 0.3$

This yields the 'WIMPless miracle' (Feng & Kumar, PRL **101**:231301,2008) since *generic* hidden sector matter $(g_h^2/m_h \sim g_\chi^2/m_\chi \sim F/16\pi^2 M)$... gives the required abundance as before!



Mass scale	Particle	Symmetry/ Quantum #	Stability	Production	Abundance
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$\Lambda_{ m Fermi} \sim$ $G_{ m F}^{-1/2}$	Neutralino? Technibaryon?	R-parity? (walking) Technicolour	violated? $\tau \sim 10^{18} \text{ yr}$ e ⁺ excess?!	'freeze-out' from thermal equilibrium Asymmetric (like the observed baryons)	$\Omega_{\rm LSP} \sim 0.3$ $\Omega_{\rm TB} \sim 0.3$

A new electroweak-scale mass particle which shares in this asymmetry (e.g. technibaryon) would have the right abundance to be dark matter ... and *explain* the ratio of dark to baryonic matter (Nussinov, PL B165:55,1985; Dodelson, Phys.Rev. D40:3252,1989)



Sterile (right-handed) neutrinos can also be the dark matter ...



Hence although they may never come into equilibrium, the relic abundance will be of order the dark matter for a mass of order KeV (Dodelson & Widrow, PRL **72**:17,1994)

... however there is no natural motivation for such a mass scale

Sterile neutrino and 3.5 keV line



DON'T FORGET THE (OXFORD) 17 KEV ANOMALY!



FIG. 2 Part of the distribution of energies of electrons from the β decay $\overline{2}$ of tritium reported by Simpson¹. The smooth curve is the expected spectrum with only a massless neutrino produced, while the dashed curve shows the expected effect of adding a massive neutrino of 17.1 keV, with a coupling *R*, of 2.3% rather than the 3% given in the text. The next point to the left of the lowest-energy point is off-scale at 1.6×10^6 counts. *E*_{th} is the value in keV at the kink.

~7σ evidence from a laboratory experiment that turned out to be a conspiracy of systematic effects!



AXION DARK MATTER

The SM admits a term which would lead to CP violation in strong interactions, hence an (unobserved) electric dipole moment for neutrons \rightarrow requires $\theta_{\text{QCD}} < 10^{-10}$

To achieve this without fine-tuning, θ_{QCD} must be made a dynamical parameter, through the introduction of a new $U(1)_{\text{Peccei-Quinn}}$ symmetry which must be broken ... the resulting (pseudo) Nambu-Goldstone boson is the QCD **axion** which later acquires a small mass through its mixing with the pion (the pNGB of QCD): $m_a = m_{\pi} (f_{\pi}/f_{\text{PQ}})$ (Kim, Phys.Rep.**150**:1,1987, Rev.Mod.Phys.**82**:557,2010; Raffelt, Phys.Rep.**198**:1,1990)



When the temperature drops to $\Lambda_{\rm QCD}$ the axion potential turns on and the coherent oscillations of relic axions contain energy density that behaves like cold dark matter with $\Omega_{\rm a}h^2 \sim 10^{11}~{\rm GeV}/f_{\rm PQ}$... however the natural P-Q scale is probably $f_{\rm PQ} \sim 10^{18}~{\rm GeV}$

Hence QCD axion dark matter would need to be *significantly diluted*, i.e. its relic abundance is not predictable (or seek anthropic explanation for why θ_{OCD} is small?)

CURRENT LIMITS ON AXIONS AND AXION-LIKE PARTICLES



Mass scale	Lightest stable particle	Symmetry/ Quantum #	Stability ensured?	Production	Abundance
A _{QCD}	Nucleons	Baryon number	$\tau > 10^{33} \text{ yr}$	'Freeze-out' from thermal equilibrium Asymmetric baryogenesis (how?)	$\Omega_{\rm B} \sim 10^{-10}$ cf. observed $\Omega_{\rm B} \sim 0.05$
Λ_{QCD} , ~ $5\Lambda_{QCD}$	Dark baryon	U(1) _{DB}	?NT	Asymmetric (like the observed baryons)	$\Omega_{DB}\!\sim 0.3$
$\Lambda_{ m Fermi} \ \sim G_{ m F}^{-1/2}$	Neutralino?	<i>R</i> -parity?	violated?	freeze-out' from thermal equilibrium	$\Omega_{LSP}\!\sim 0.3$
	Technibaryon?	(walking) Technicolour	τ~10 ¹⁸ yr	Asymmetric (like the observed baryons)	$\Omega_{TB}\!\sim 0.3$
$\Lambda_{ m hidden \ sector} \ \sim (\Lambda_{ m F} M_{ m P})^{1/2}$	Crypton? hidden valley?	Discrete (model-	$\tau \gtrsim 10^{18} \mathrm{yr}$	Varying gravitational field during inflation	$\Omega_X \sim 0.3?$
$\Lambda_{see-saw}$ $\sim \Lambda_{Fermi}^2 / \Lambda_{B-L}$	Neutrinos	dependent) Lepton number	Stable _.	Thermal (like CMB)	$\Omega_{v} > 0.003$
M _{string} M _{Planck}	Kaluza-Klein states?	? Peccei-Quinn	?	?	?
- miner	Axions		stable	Field oscillations	$\Omega_{\rm a} \gg 1!$