

Reduced transport regions in rotating tokamak plasmas

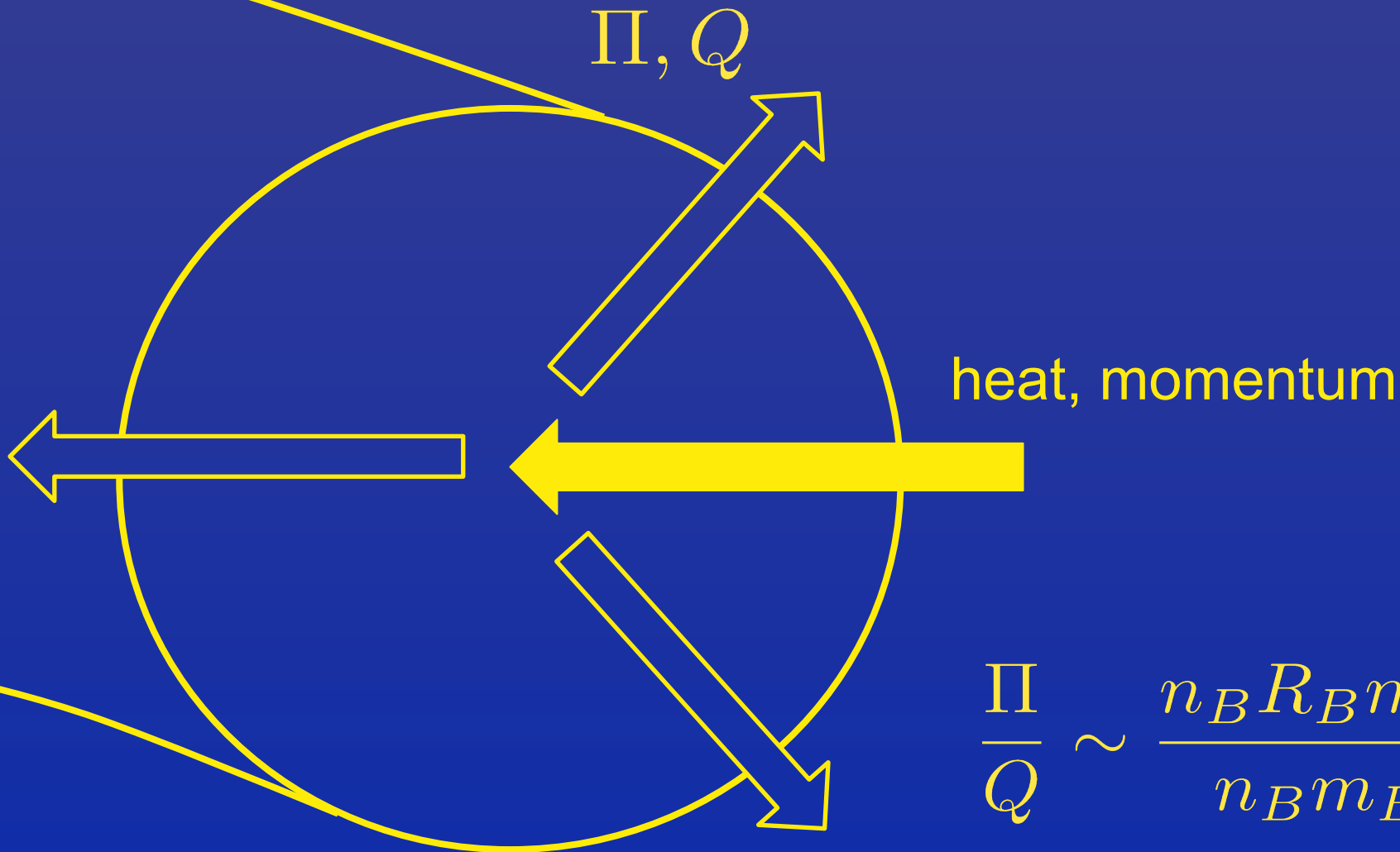
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Power/Torque balance for beam injection



$$\frac{\Pi}{Q} \sim \frac{n_B R_B m_B v_B^2}{n_B m_B v_B^3} = \frac{R_B}{v_B}$$

Model fluxes

- Simple model for fluxes with parameters chosen to fit zero magnetic shear results from GS2:

$$Q = Q_t + Q_n$$

$$\Pi = \Pi_t + \Pi_n$$

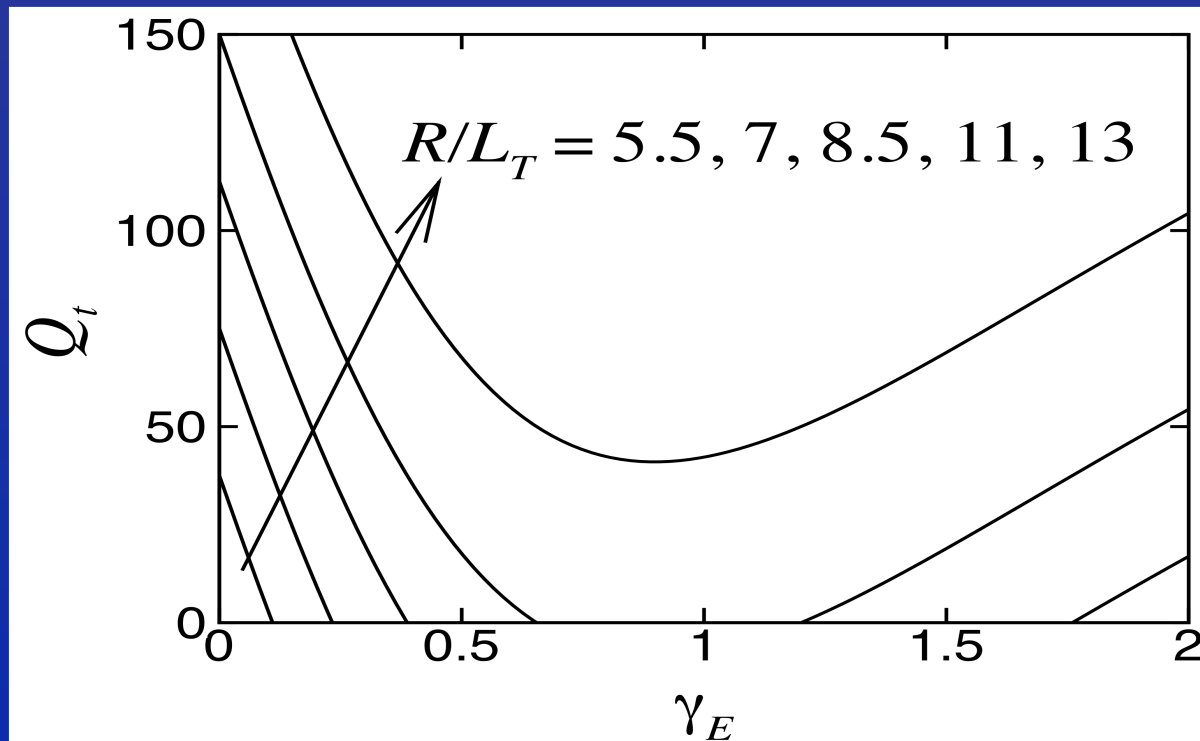
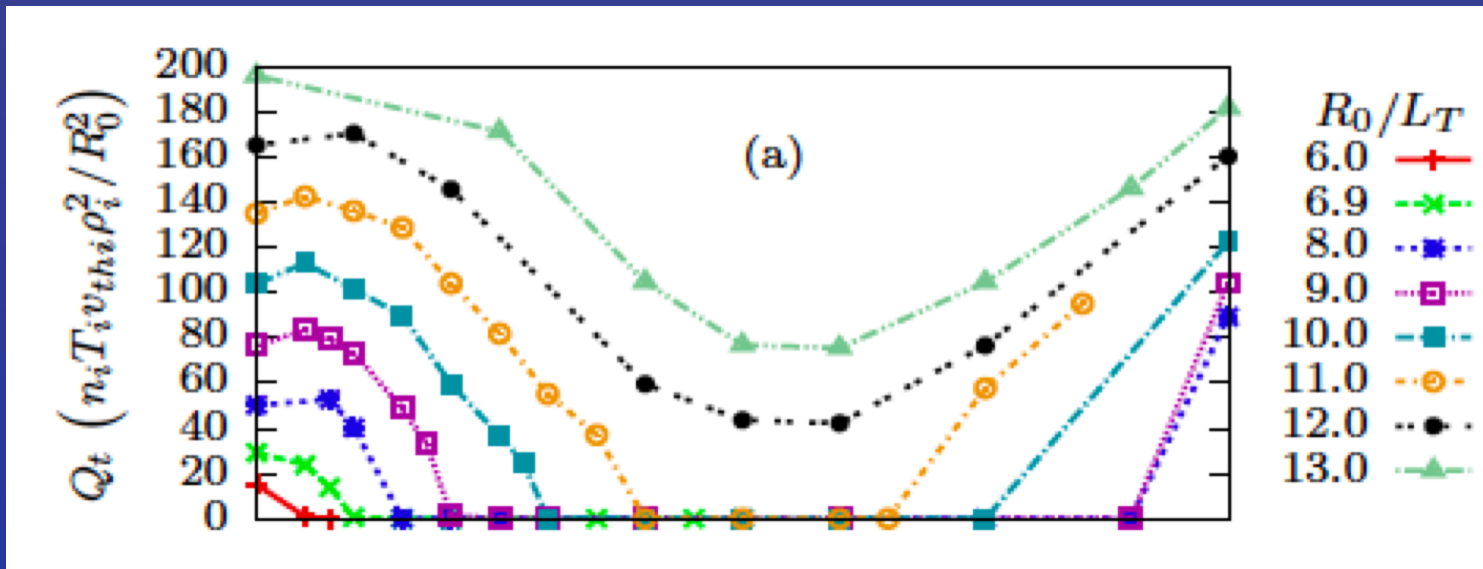
$$\bar{Q}_t \equiv \frac{Q_t}{nTv_{th}} \left(\frac{R}{\rho} \right)^2 \equiv \chi_t \left[\frac{R}{L_T} - \left(\frac{R}{L_T} \right)_c \right]$$

$$\bar{Q}_n \equiv \frac{Q_n}{nTv_{th}} \left(\frac{R}{\rho} \right)^2 \equiv \frac{\chi_n}{T^2} \frac{R}{L_T}$$

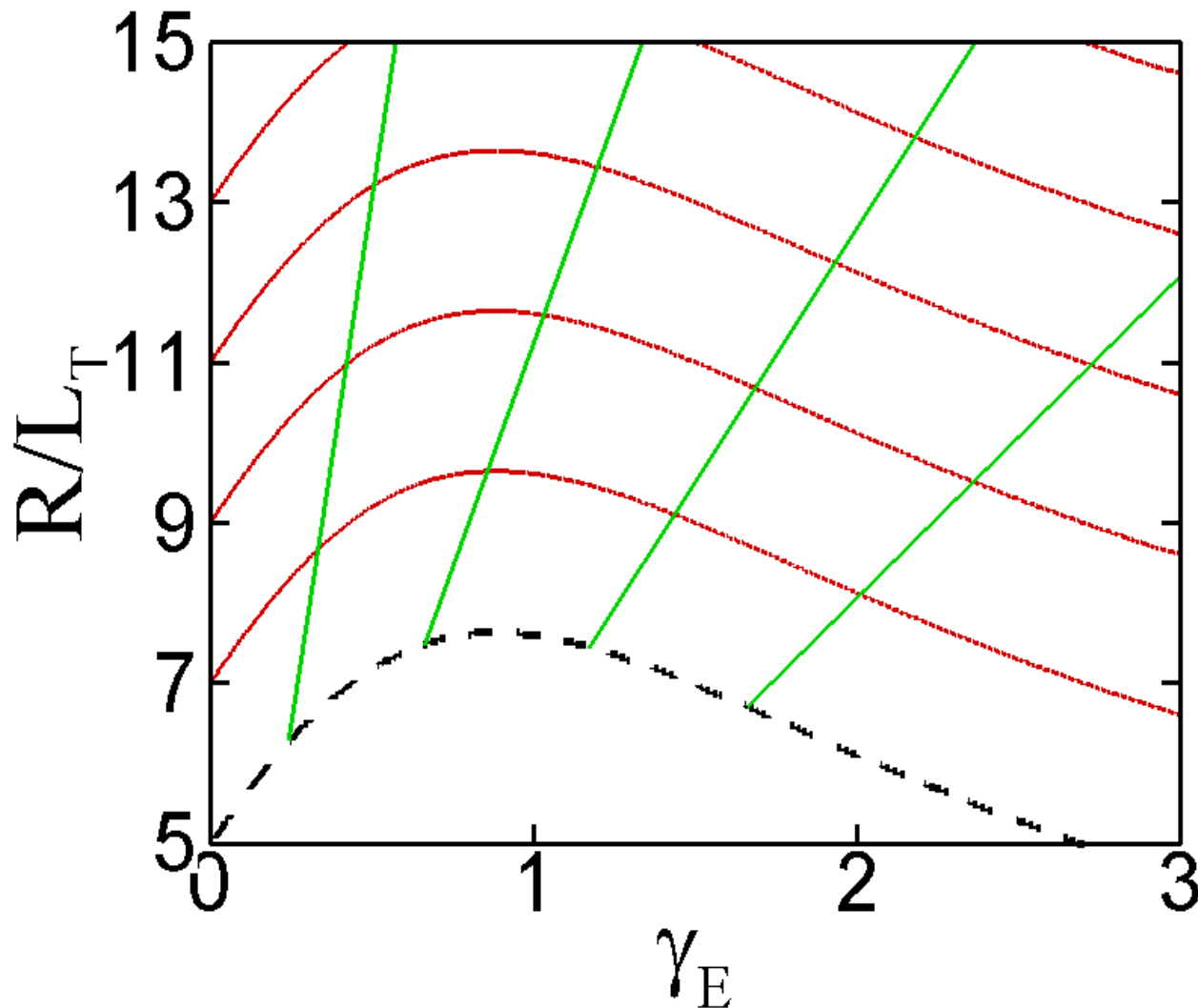
$$\left(\frac{R}{L_T} \right)_c \equiv \frac{\alpha_1 \gamma_E + (R/L_T)_{c0}}{1 + \alpha_2 \gamma_E^2}$$

$$\bar{\Pi}_{t,n} \equiv \frac{\Pi_{t,n}}{mnRv_{th}^2} \left(\frac{R}{\rho} \right)^2 = \bar{Q}_{t,n} \text{Pr}_{t,n} \frac{\gamma_E}{R/L_T}$$

Model fluxes



Balance w/o neoclassical

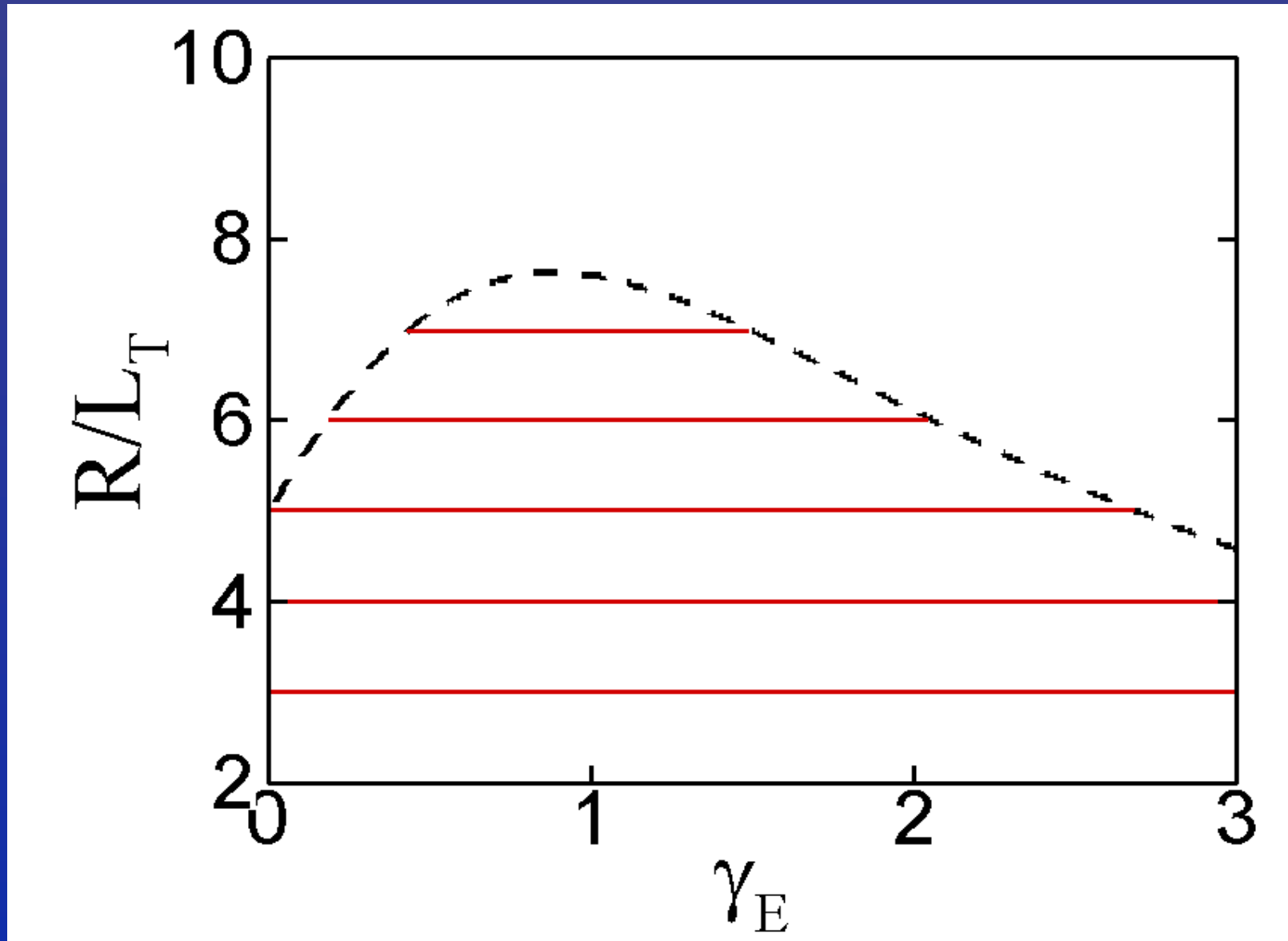


- \bar{Q} = red lines
- $\bar{\Pi}/\bar{Q}$ = green lines

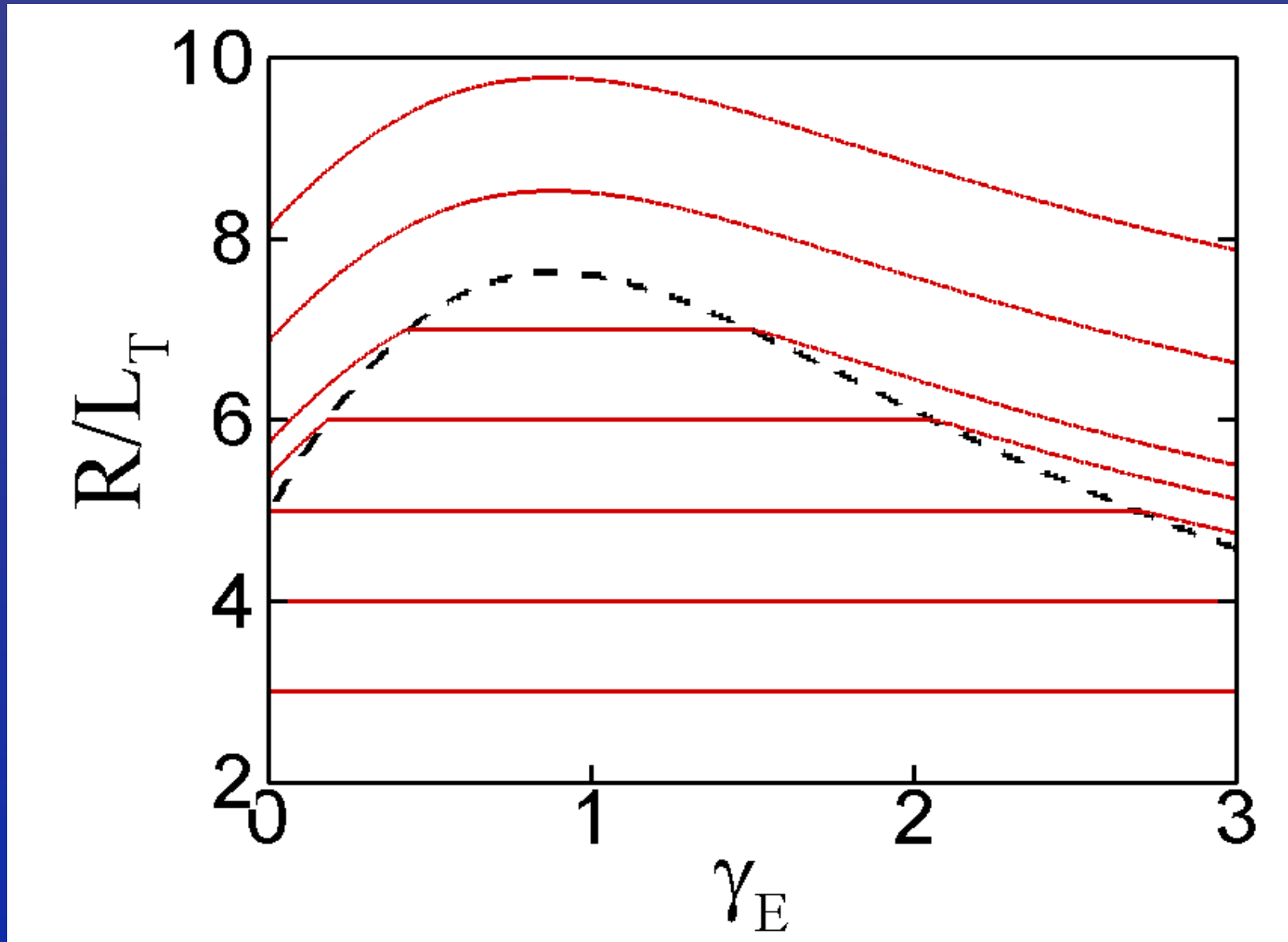
$$\frac{R}{L_t} = \frac{Pr_t}{\bar{\Pi}/\bar{Q}} \gamma_E$$

- Critical gradient = dashed line
- For given $\bar{\Pi}/\bar{Q}$ and \bar{Q} , only one solution

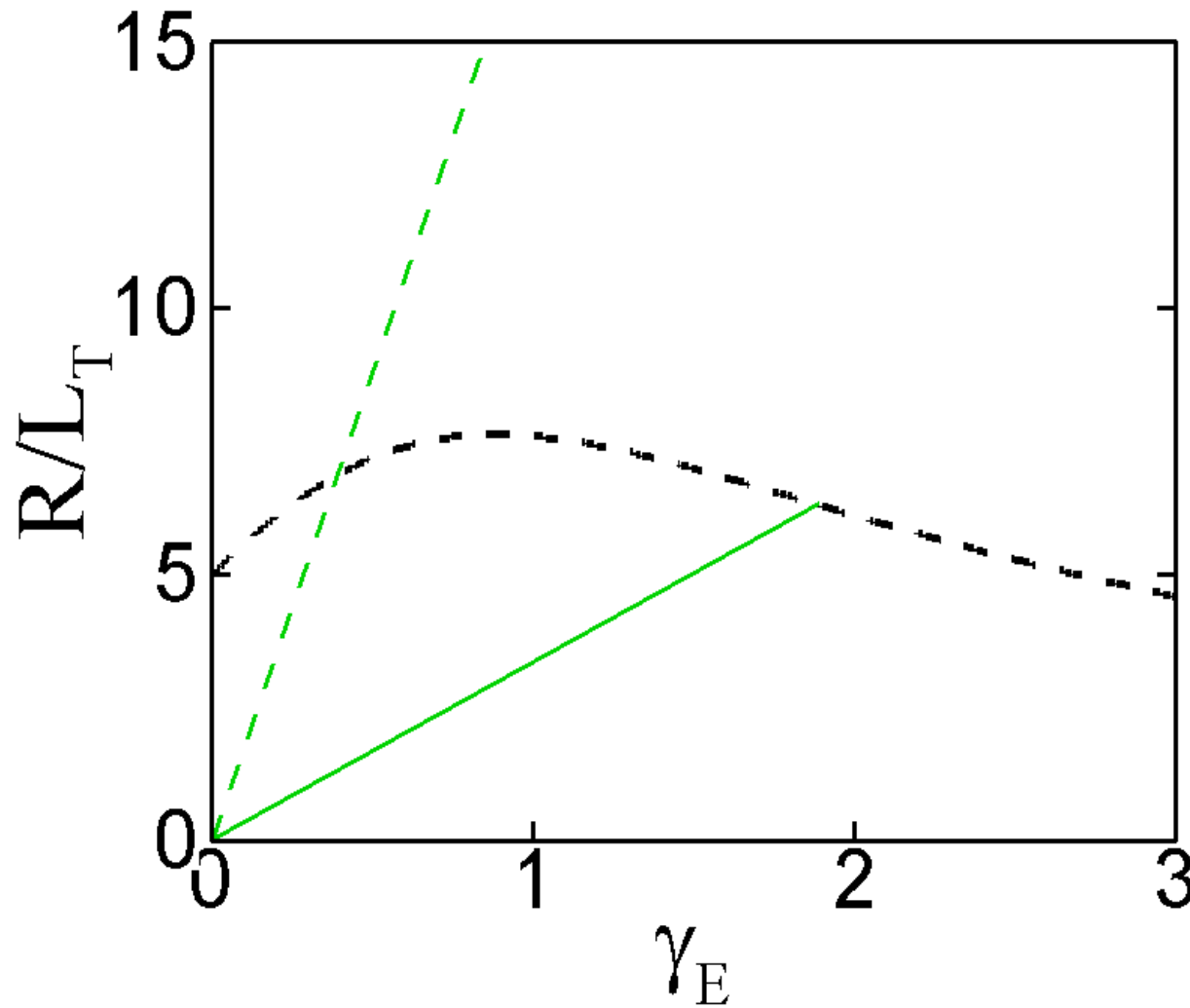
Neoclassical energy flux



Balance with neoclassical



Curves of constant $\bar{\Pi}/\bar{Q}$



- Neoclassical

$$\frac{R}{L_t} = \frac{\text{Pr}_n}{\bar{\Pi}/\bar{Q}} \gamma_E$$

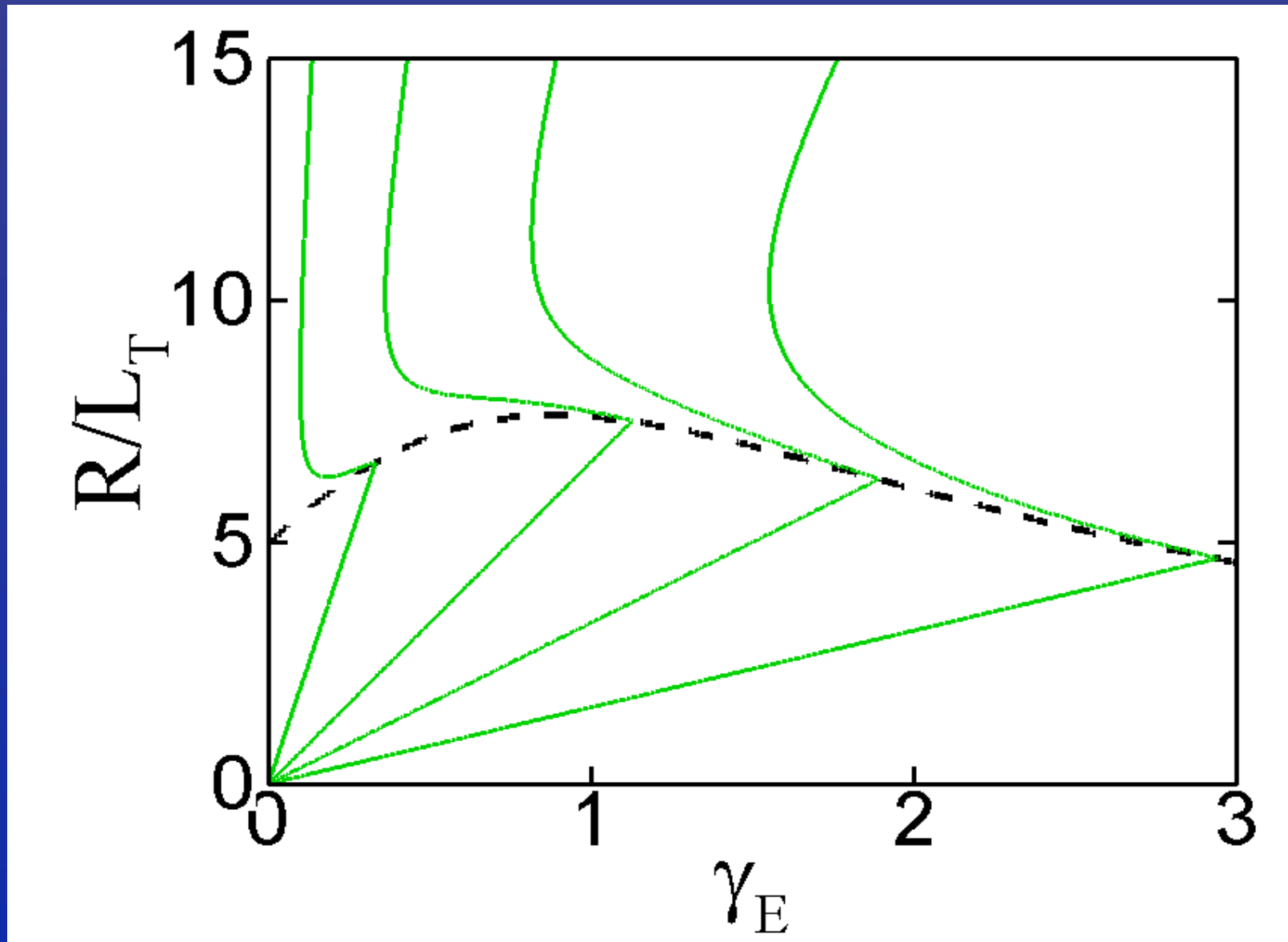
- Turbulent

$$\frac{R}{L_t} = \frac{\text{Pr}_t}{\bar{\Pi}/\bar{Q}} \gamma_E$$

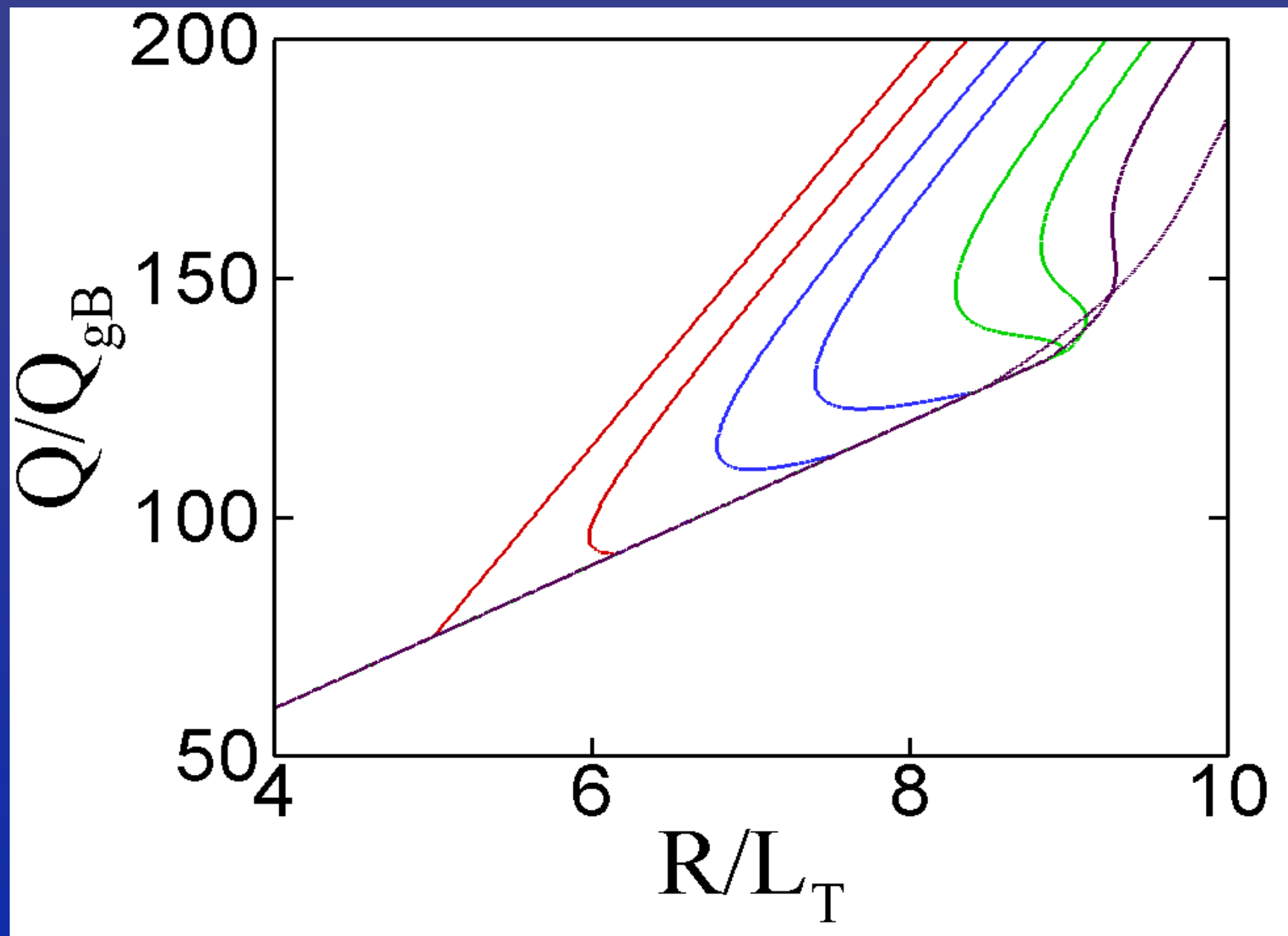
- Prandtl numbers

$$\text{Pr}_n \ll \text{Pr}_t$$

Curves of constant $\bar{\Pi}/\bar{Q}$



Total energy flux



Temperature dependence

- Have been using \bar{Q} and $\bar{\Pi}/\bar{Q}$ for local analysis, but these quantities depend on local temperature:

$$\bar{Q} \sim \frac{Q}{T^{5/2}} \qquad \frac{\bar{\Pi}}{\bar{Q}} \sim \frac{\Pi}{Q} T^{1/2}$$

- Consequently, \bar{Q} is a label for radius, and the contours of constant $\bar{\Pi}/\bar{Q}$ vary from radius to radius
- Better to consider $\bar{\Pi}/\bar{Q}^{4/5}$ or $\hat{\Pi}/\hat{Q}$, which are independent of temperature

$$\hat{Q} = \frac{Q}{T_0^{5/2}} \qquad \hat{\Pi} = \frac{\Pi}{T_0^2}$$

Solving for radial profiles

- Expressions for fluxes:

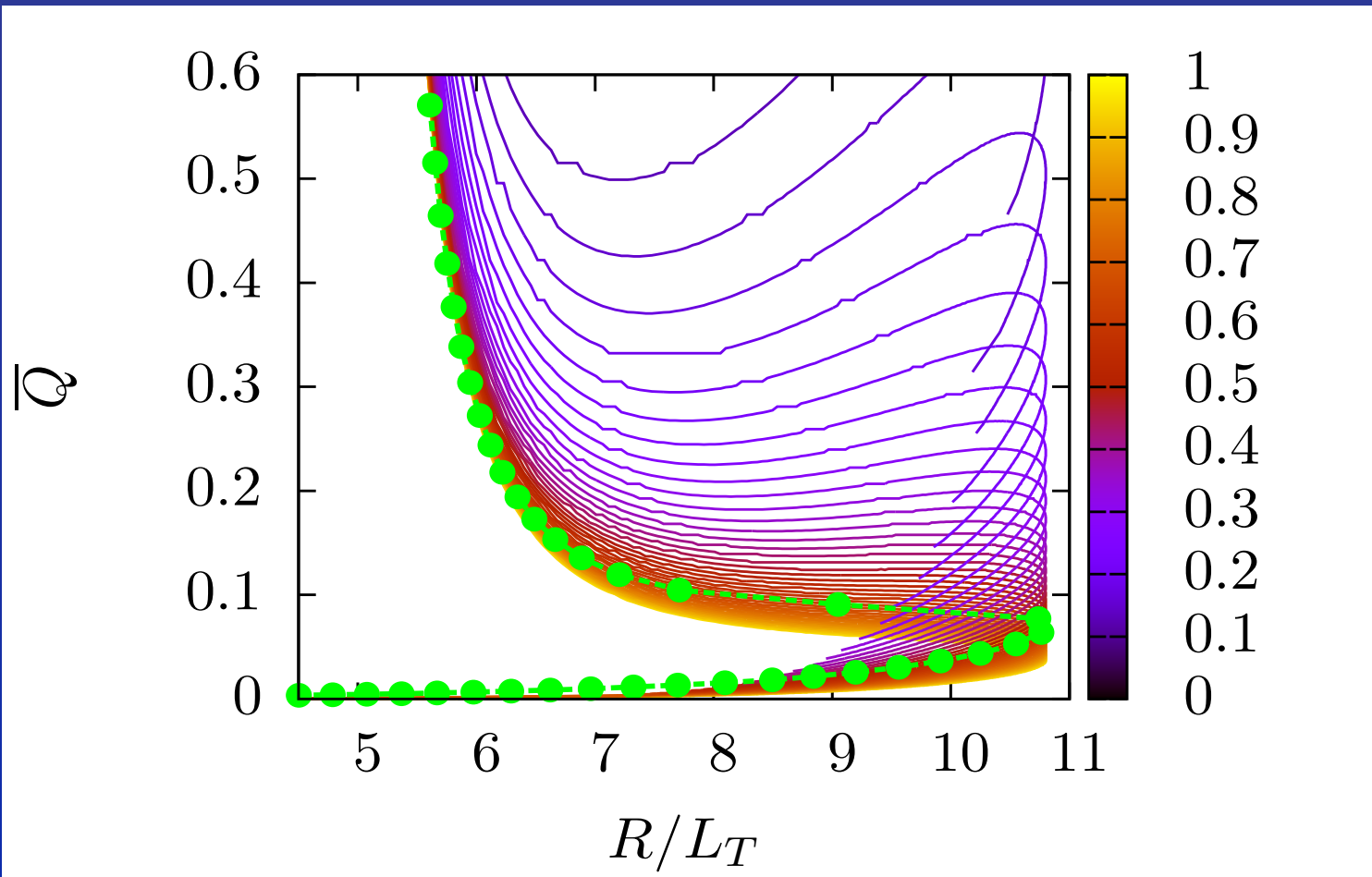
$$\hat{Q}(\kappa, \gamma_E, T) = \hat{T}^{5/2} \left(\hat{\chi}_t (\kappa - \kappa_c) + \frac{\hat{\chi}_n}{\hat{T}^2} \kappa \right)$$

$$\hat{\Pi}(\kappa, \gamma_E, T) = \gamma_E \left(\hat{\chi}_t \left(1 - \frac{\kappa_c}{\kappa} \right) \text{Pr}_t \hat{T}^2 + \hat{\chi}_n \text{Pr}_n \right)$$

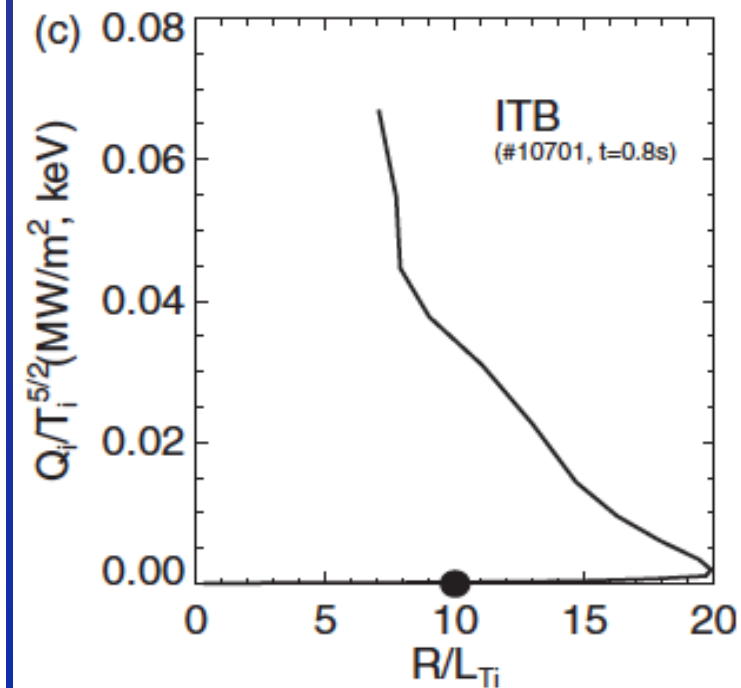
- Radial profiles of \hat{Q} and $\hat{\Pi}$ are inputs. Given \hat{T} at one radius, we can solve for γ_E and κ at that radius. With \hat{T} and κ , we can obtain \hat{T} at nearby radii. Repeat process to construct radial profiles.

Numerical results

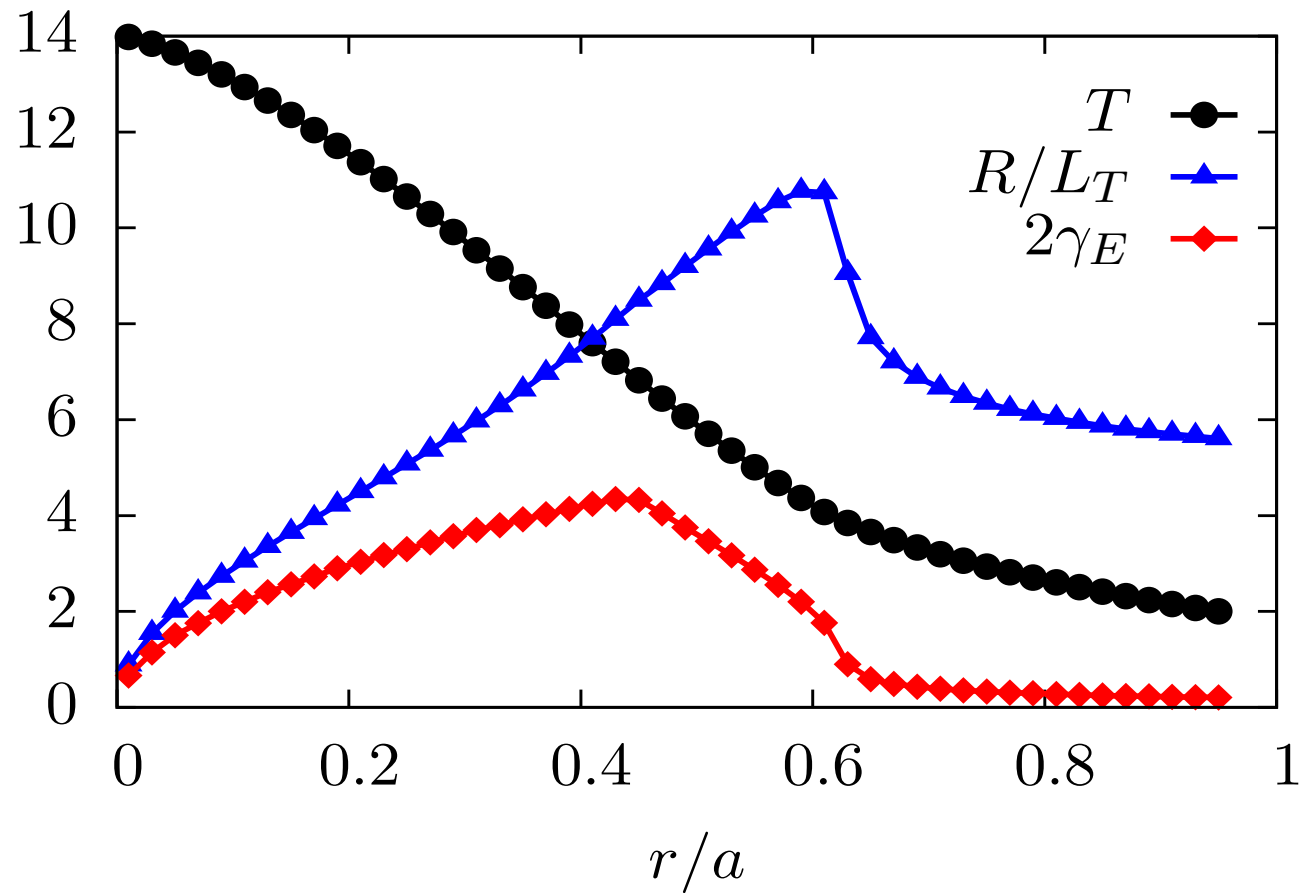
- Here, $Q \sim \sqrt{r/a}$, $P_i/Q=0.1$, Edge $T=2$ keV

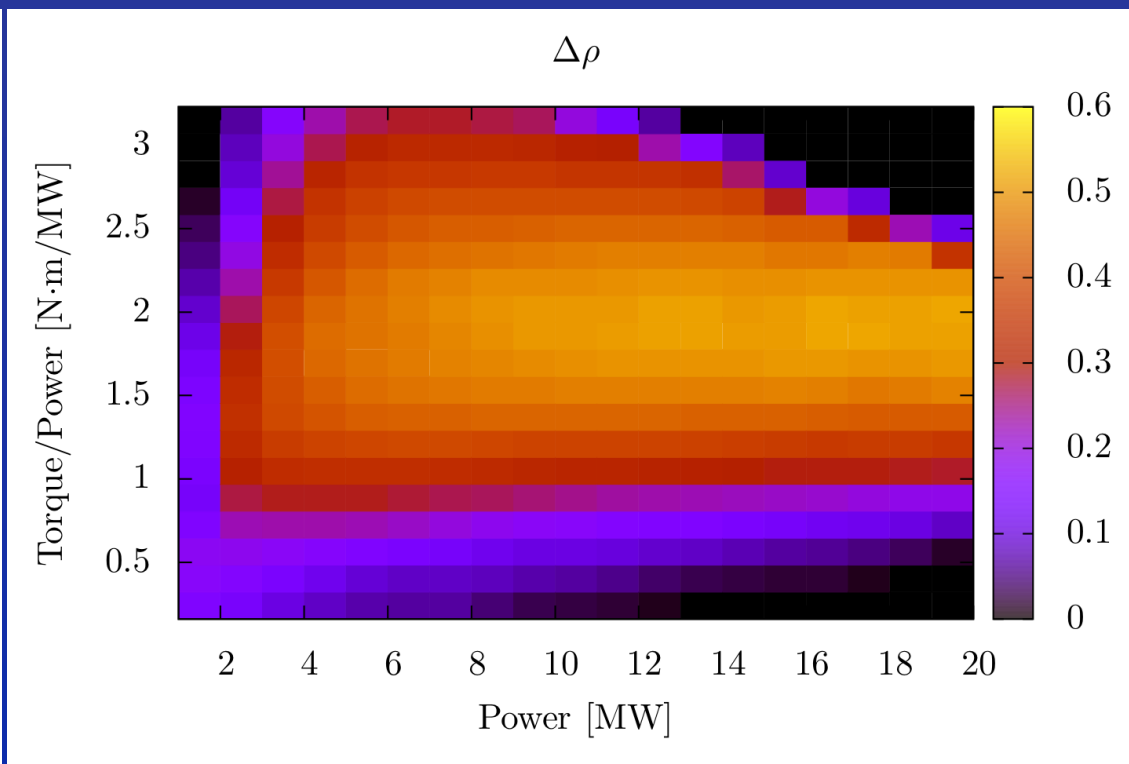
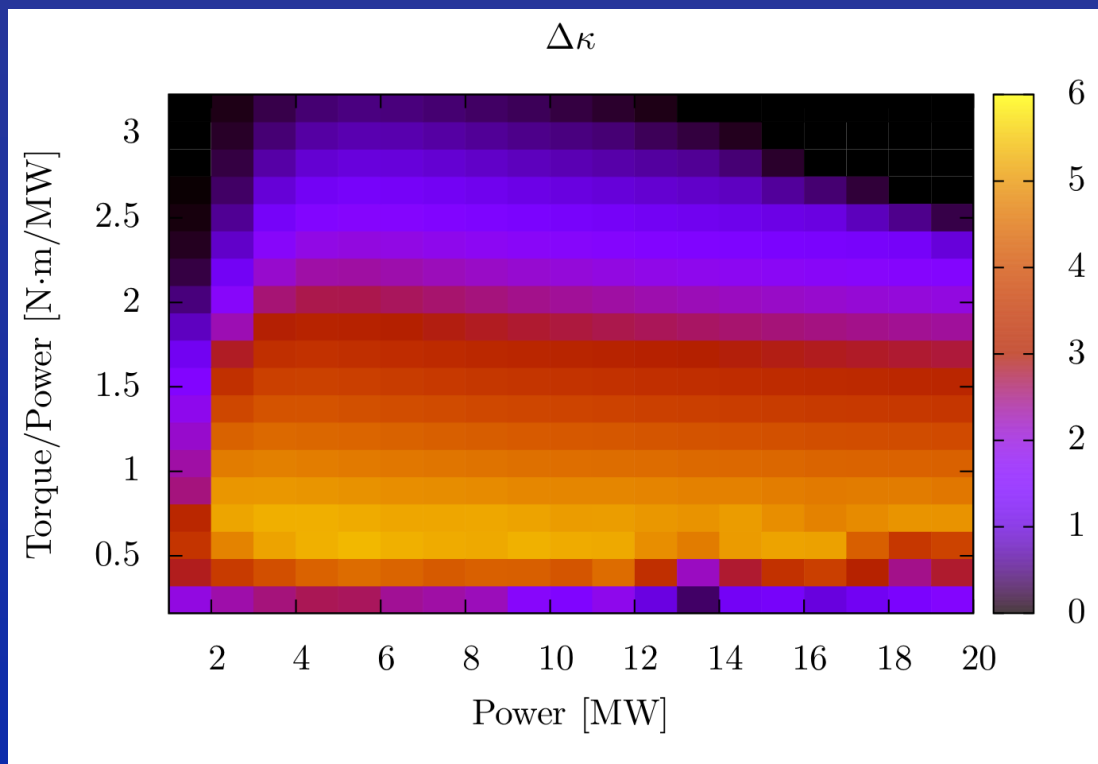
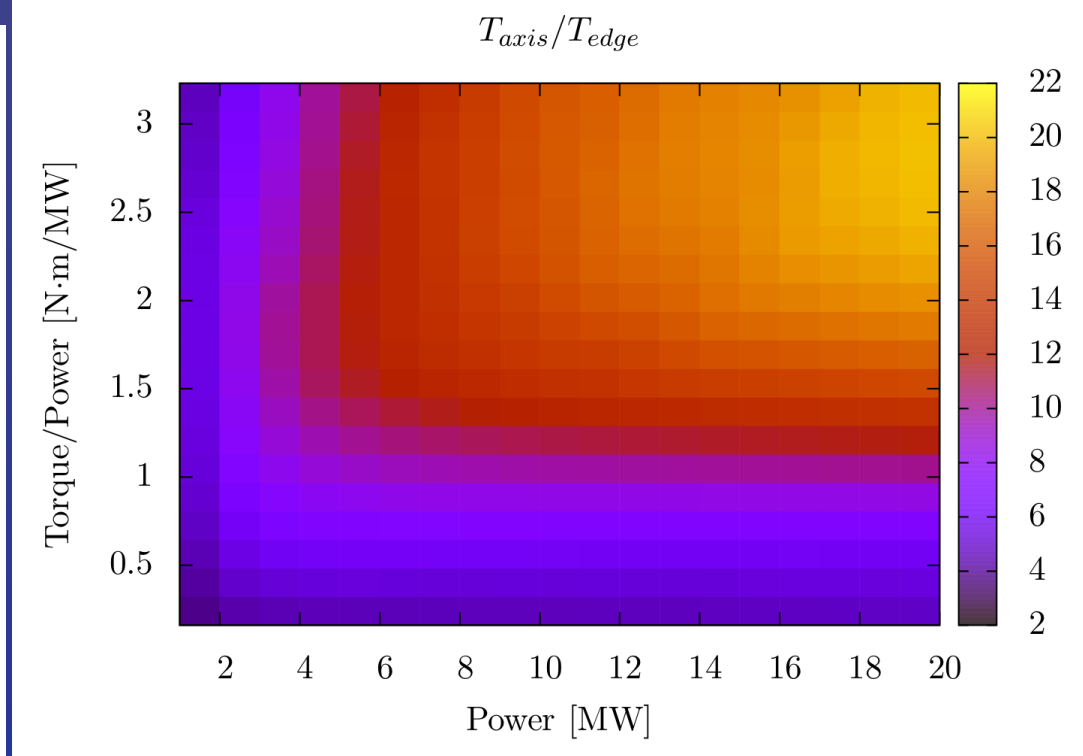
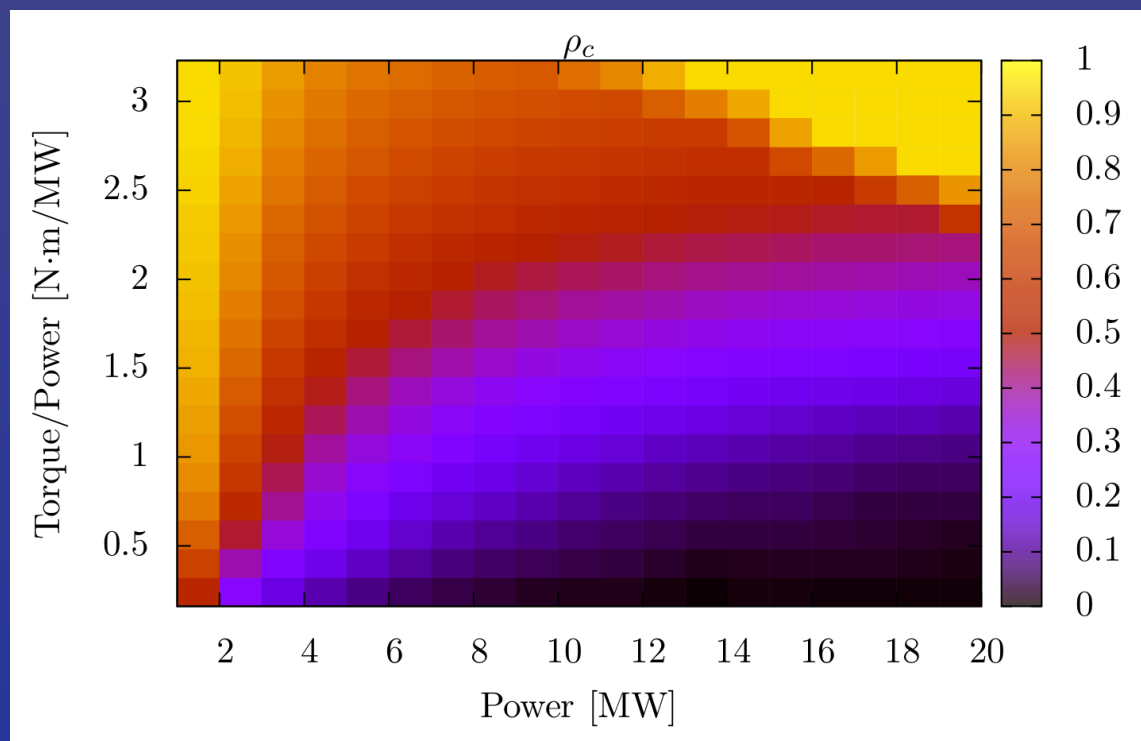


Wolf, PPCF 45 (2003)



Extension to 1D (radial)





What have we learned?

- Significant enhancement of temperature gradient obtained solely through flow shear and magnetic shear
- Enhancement for simple geometry not sufficient to account for strong ITBs
 - Shafranov shift
 - Plasma shaping
- Multiple solutions not necessary to obtain localized enhancement of temperature gradient
- Understanding transition to enhanced gradient regime is work in progress

The future: multiscale simulation

- In TRINITY [Barnes *et al.*, PoP **17**, 056109 (2010)], turbulent fluctuations calculated in small regions of fine space-time grid embedded in “coarse” grid (for mean quantities)

Flux tube simulation domain

