The Trinity algorithm: local gyrokinetics + global transport = predictive model of core plasma dynamics

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## Overview

- Motivation
- Multiscale approach
- Trinity simulation results
- Conclusions









## Scale separation in ITER

$$\frac{\partial n}{\partial t} + \nabla \cdot \mathbf{\Gamma} = S_n$$

$$\frac{\partial p}{\partial t} + \nabla \cdot \mathbf{Q} + \dots = S_{p}$$

Physics	Perpendicular spatial scale	Temporal scale
Turbulence from ETG modes	$k_{\perp}^{-1}$ ~ 0.001 - 0.1 cm	$\omega_{*}$ ~ 0.5 - 5.0 MHz
Turbulence from ITG modes	$k_{\perp}^{-1}$ ~ 0.1 - 8.0 cm	$\omega_{*}$ ~ 10 - 100 kHz
Transport barriers	Measurements suggest width ~ 1 - 10 cm	100 ms or more in core?
Discharge evolution	Profile scales ~ 100 cm	Energy confinement time ~ 2 - 4 s

#### Direct simulation cost

- Grid spacings in space (3D), velocity (2D) and time:
  - $\Delta x \sim 0.001 \ cm, \ L_x \sim 100 \ cm$
  - $\Delta v \sim 0.1 v_{th}, \quad L_v \sim v_{th}$
  - $\Delta t \sim 10^{-7} s, \ L_t \sim 1 s$
- Grid points required:



- $(L_x/\Delta x)^3 \times (L_v/\Delta v)^2 \times (L_t/\Delta t) \sim 10^{24}$
- Factor of ~10<sup>10</sup> more than largest fluid turbulence calculations
- Direct simulation not possible; need physics guidance

#### Improved simulation cost

- Field-aligned coordinates take advantage of  $k_{||} \ll k_{\perp}$ : savings of ~1000
- Statistical periodicity in poloidal direction takes advantage of  $k_{\perp}^{-1} \ll L_{\theta}$ : savings of ~100
- Total saving of ~10<sup>5</sup>
- Factor of ~10<sup>5</sup> more than largest fluid turbulence calculations
- Simulation still not possible; need multiscale approach

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#### Gyrokinetic multiscale assumptions

- Fluctuation amplitude small compared with equilibrium:  $f=F_0+\delta f, \ \ \delta f/F_0\sim\epsilon\equiv\rho/L$
- Separation of turbulence and equilibrium space scales:  $abla F_0 \sim F_0/L, \ \ 
  abla_{\parallel} \delta f \sim \delta f/L, \ \ 
  abla_{\perp} \sim \delta f/
  ho$

Separation of turbulence and equilibrium time scales:



$$\partial_t F_0 \sim \tau^{-1} F_0, \quad \partial_t \delta f \sim \omega \delta f \sim \nu \delta f$$
  
 $\tau^{-1} \sim \epsilon^2 \omega \sim \epsilon^3 \Omega$ 

- Sub-sonic drifts:  $v_E \sim \epsilon v_{th}$
- "Smooth" velocity space:  $\partial_v f \sim f/v_{th}$

#### Key results: turbulence and transport

Gyrokinetic equation for turbulence:  $\partial h/\partial t + v_{\parallel} \hat{\mathbf{b}} \cdot \nabla h + \langle \mathbf{v}_{\chi} \rangle_{\mathbf{R}} \cdot \nabla (F_0 + h) + \mathbf{v}_{\mathbf{B}} \cdot \nabla h = \frac{qF_0}{T_0} \frac{\partial \langle \chi \rangle_{\mathbf{R}}}{\partial t} + \langle C[h] \rangle_{\mathbf{R}}$ 

Moment equations for equilibrium evolution:

$$\begin{aligned} \frac{\partial n_s}{\partial t} &= -\frac{\partial \psi}{\partial V} \frac{\partial}{\partial \psi} \left[ \frac{\partial V}{\partial \psi} \left\langle \mathbf{\Gamma}_s \cdot \nabla \psi \right\rangle \right] \\ \frac{3}{2} \frac{\partial \left( n_s T_s \right)}{\partial t} &= -\frac{\partial \psi}{\partial V} \frac{\partial}{\partial \psi} \left[ \frac{\partial V}{\partial \psi} \left\langle \mathbf{Q}_s \cdot \nabla \psi \right\rangle \right] \\ &+ T_s \left( \frac{\partial \ln n_s}{\partial \psi} - \frac{3}{2} \frac{\partial \ln T_s}{\partial \psi} \right) \left\langle \mathbf{\Gamma}_s \cdot \nabla \psi \right\rangle + \frac{\partial \ln T_s}{\partial \psi} \left\langle \mathbf{Q}_s \cdot \nabla \psi \right\rangle \\ &- \left\langle \int d^3 v \frac{h_s T_s}{F_{0,s}} \left\langle C(h_s) \right\rangle_{\mathbf{R}} \right\rangle + n_s \nu_{\epsilon}^{su} \left( T_u - T_s \right) \end{aligned}$$

## Multiscale grid



## Flux tube spatial simulation domain for microturbulence

- Small regions of fine grid (for turbulence) embedded in "coarse" radial grid (for equilibrium)
- Turbulent fluxes and heating in small regions calculated using flux tubes (equivalent to flux surfaces)
- Flux tubes = radial grid points in large-scale transport equations
- Small regions of fine grid (for turbulence) embedded in "coarse" time grid (for equilibrium)
- Steady-state (time-averaged) turbulent fluxes and heating in this volume simulated using flux tubes
- Flux tube sim = time grid point in long-time transport equations

Flux tube temporal simulation domain for microturbulence



#### Flux tubes minimize flux surface grid points



Image of MAST simulation courtesty of G. Stantchev

#### More flux tube savings

# • (Near) perfect parallelization:

- Only communication between flux tubes occurs when solving transport equations, which is infrequent
- Flux tube calculations are independent

Strong scaling of a single flux tube simulation (GS2)



#### Validity of flux tube approximation

- Lines represent global simulations from GYRO
- Dots represent local (flux tube) simulations from GS2
- Excellent agreement for  $\rho_* \ll 1$



\*J. Candy, R.E. Waltz and W. Dorland, The local limit of global gyrokinetic simulations, Phys. Plasmas **11** (2004) L25.

#### Multiscale simulation cost

- Grid spacings in radius and velocity (2D) roughly unchanged
- Major savings is in time domain:

Turbulence:  $\Delta t \sim 10^{-7} s$ ,  $L_t \sim 10^{-4} s$ 

Transport:  $\Delta au \sim 0.1 \ s, \ L_{ au} \sim 1 \ s$ 

• Required number of grid points:



 $(L_r/\Delta r) \times (L_\theta/\Delta \theta) \times (L_\phi/\Delta \phi) \times (L_v/\Delta v)^2 \times (L_t/\Delta t) \times (L_\tau/\Delta \tau) \sim 10^{15}$ 

• Savings of ~10<sup>3</sup> over conventional numerical simulation

Coarse space-time grid

#### **Trinity schematic**



#### **Trinity schematic**



### Sampling profile with flux tubes



## Sampling profile with flux tubes



### Sampling profile with flux tubes



#### **Trinity schematic**



• Transport equations are stiff, nonlinear PDEs:

$$\frac{\partial n}{\partial t} = H(r) \frac{\partial G[\Gamma, n, T, r]}{\partial r} \qquad \qquad \Gamma \to \frac{\Gamma}{n v_{th}}$$

• General (single-step) time discretization:

$$\frac{n^{m+1} - n^m}{\Delta \tau} = \alpha \left[ H \frac{\partial G}{\partial r} \right]^{m+1} + (1 - \alpha) \left[ H \frac{\partial G}{\partial r} \right]^m$$

• 2nd order centered difference in radial coordinate (equally spaced grid):  $\frac{\partial G}{\partial r} = \frac{G_{j+1/2} - G_{j-1/2}}{\Delta r}$ 

\*S.C. Jardin, G. Bateman, G.W. Hammett, and L.P. Ku, On 1D diffusion problems with a gradient-dependent diffusion coefficient, J. Comp. Phys. **227**, 8769 (2008).

 Treat nonlinear terms implicitly with (single-iteration) Newton's Method

$$G_{j}^{m+1} \approx G_{j}^{m} + (\mathbf{y} - \mathbf{y}^{m}) \left. \frac{\partial G_{j}}{\partial \mathbf{y}} \right|_{\mathbf{y}^{m}}$$
$$\mathbf{y} = \left[ \{n_{k}\}, \{p_{i_{j}}\}, \{p_{e_{j}}\} \right]^{T}$$

 Simplifying assumption: normalized fluxes depend primarily on gradient scale lengths

 $\Gamma_j \approx \Gamma[(R/L_n)_j, (R/L_{p_i})_j, (R/L_{p_e})_j]$ 

 $\Rightarrow \frac{\partial \Gamma_j}{\partial n_k} \approx \frac{\partial \Gamma_j}{\partial (R/L_n)_j} \frac{\partial (R/L_n)_j}{\partial n_k} \qquad \left(\frac{R}{L_n}\right)_{\pm} \approx \mp \frac{R}{\Delta r} \frac{n_{\pm} - n_j}{n_{\pm}}$ 

- Calculating flux derivative approximations:
  - at every radial grid point, simultaneously calculate  $\Gamma_j[(R/L_n)_j^m]$  and  $\Gamma_j[(R/L_n)_j^m+\delta]$  using 2 different flux tubes
  - use 2-point finite differences:

$$\frac{\partial \Gamma_j}{\partial (R/L_n)_j} \approx \frac{\Gamma_j[(R/L_n)_j^m] - \Gamma_j[(R/L_n)_j^m + \delta]}{\delta}$$

- Example calculation with 10 radial grid points:
  - evolve density and electron/ion pressures
  - simultaneously calculate fluxes for equilibrium profile and for 3 separate profiles (one for each perturbed gradient scale length)
  - total of 40 flux tube simulations running simultaneously
  - ~2000-4000 processors per flux tube => scaling to over 100,000 processors with >80% efficiency



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### Preliminary tests for Trinity

- Test 1: choose fluxes so that heat transport equations reduce to diffusion equation
- Test 2: choose fluxes so that heat transport equations reduce to advection equation



#### Preliminary results from Trinity



- Single ion species
- Adiabatic electrons
- Electrostatic
- 60 MW external heat source into ions
- Local equilibrium model with circular flux surfaces
- 8 radial grid points (flux tubes)
- Temperature at r=0.8a fixed at 4 keV
- Only ion temperature evolved
- Takes ~20 minutes on ~2000 processors

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### Conclusions

- Multiscale nature of turbulent transport can be exploited to provide significant savings in time domain
- Implemented in coupled flux tube code, Trinity
- Highly parallelized -- should scale to over 10<sup>5</sup> processors
- Trinity (interfaced with GS2) is capable of running with multiple species, electromagnetic effects, realistic geometry, physical collisional effects (such as heating), etc.
- Missing model for momentum transport -- current area of activity
- Turbulent transport/heating code like Trinity necessary part of full numerical tokamak simulations -- MCF will not work if we can not reliably reduce turbulent transport

- Nonlinear turbulence simulation runs until fluxes converged
  - convergence criterion:

$$\epsilon \equiv \sqrt{\frac{1}{m-j} \sum_{i=j}^{m-1} \left(\overline{\Gamma}_m^2 - \overline{\Gamma}_i^2\right)} < \epsilon_0 \qquad \overline{\Gamma}_m = \frac{1}{t_m} \sum_{i=1}^m \Gamma_i \ (\Delta t)_i$$

- turbulence for new transport time step initialized to saturated state from previous transport time step -- faster convergence
- Sources and initial profiles are analytically specified. In process of adding capability to read in experimental profiles for these quantities
- Option to use model fluxes (IFS-PPPL, offset linear, quasilinear, etc.)
- Boundray conditions:
  - fixed n and T at outer edge of simulation domain
  - zero flux boundary condition at magnetic axis
  - innermost point at 0.5\*delta r and outermost point specified

- Nonlinear implicit method<sup>\*</sup> (single-iteration Newton)
- 2nd order centered difference in radial coordinate
- Simplifying physics assumption: normalized fluxes depend primarily on gradient scale lengths:

 $\Gamma_j \approx \Gamma[(R/L_n)_j, (R/L_{p_i})_j, (R/L_{p_e})_j]$ 

$$\Rightarrow \frac{\partial \Gamma_j}{\partial n_k} \approx \frac{\partial \Gamma_j}{\partial (R/L_n)_j} \frac{\partial (R/L_n)_j}{\partial n_k}$$

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