

## Further Quantum Mechanics HT 2014

### Problems 2 (Easter Vacation)

#### Radiative transitions

**2.1\*** Let  $|E, l, m\rangle$  denote a stationary state of an atom with orbital angular-momentum quantum numbers  $l, m$ , and let  $x_{\pm} = x \pm iy$  be complex position operators while  $L_{\pm} = L_x \pm iL_y$  are the usual orbital angular-momentum ladder operators. Show that  $x_{\pm}|E, l, m\rangle$  is an eigenket of  $L_z$  with eigenvalue  $m \pm 1$ . Show also that

$$[L_+, x_+] = [L_-, x_-] = 0 \quad \text{and} \quad [L_+, x_-] = -[L_-, x_+] = 2z.$$

Hence show that

$$\langle E', l', m | z | E, l, m \rangle = \alpha_+(l, m) \langle E', l', m | x | E, l, m + 1 \rangle - \alpha_-(l', m) \langle E', l', m - 1 | x | E, l, m \rangle.$$

where  $\alpha_{\pm}(l, m) \equiv \sqrt{l(l+1) - m(m \pm 1)}$ . [Hint: compute  $\langle E', l', m | x | E, l, m + 1 \rangle$ ]

**2.2** Given that  $a_0 = \hbar/(\alpha m_e c)$  show that the product  $a_0 k$  of the Bohr radius and the wavenumber of a photon of energy  $E$  satisfies

$$a_0 k = \frac{E}{\alpha m_e c^2}. \quad (2.1)$$

Hence show that the wavenumber  $k_{\alpha}$  of an H $\alpha$  photon satisfies  $a_0 k_{\alpha} = \frac{5}{72}\alpha$  and determine  $\lambda_{\alpha}/a_0$ . What is the connection between this result and our estimate that  $\sim 10^7$  oscillations are required to complete a radiative decay. Does it imply anything about the way the widths of spectral lines from allowed atomic transitions vary with frequency?

**2.3** Given that a system's Hamiltonian is of the form

$$H = \frac{p^2}{2m_e} + V(\mathbf{x}) \quad (2.2)$$

show that  $[x, [H, x]] = \hbar^2/m_e$ . By taking the expectation value of this expression in the state  $|k\rangle$ , show that

$$\sum_{n \neq k} |\langle n | x | k \rangle|^2 (E_n - E_k) = \frac{\hbar^2}{2m_e}, \quad (2.3)$$

where the sum runs over all the other stationary states.

The **oscillator strength** of a radiative transition  $|k\rangle \rightarrow |n\rangle$  is defined to be

$$f_{kn} \equiv \frac{2m_e}{\hbar^2} (E_n - E_k) |\langle n | x | k \rangle|^2. \quad (2.4)$$

Show that  $\sum_n f_{kn} = 1$ . What is the significance of oscillator strengths for the allowed radiative transition rates of atoms?

**2.4** With  $|nlm\rangle$  a stationary state of hydrogen, which of these matrix elements is non-zero?

$$\begin{array}{lll} \langle 100 | z | 200 \rangle & \langle 100 | z | 210 \rangle & \langle 100 | z | 211 \rangle \\ \langle 100 | z | 300 \rangle & \langle 100 | z | 310 \rangle & \langle 100 | z | 320 \rangle \\ \langle 100 | x | 200 \rangle & \langle 100 | x | 210 \rangle & \langle 100 | x | 211 \rangle \end{array}$$

**2.5** With  $|nlm\rangle$  a stationary state of hydrogen and given that

$$\langle \mathbf{x} | 100 \rangle = \frac{2e^{-r/a_0}}{a_0^{3/2}} Y_0^0; \quad \langle \mathbf{x} | 210 \rangle = \frac{r}{a_0} \frac{e^{-r/2a_0}}{\sqrt{3}(2a_0)^{3/2}} Y_1^0(\theta, \phi); \quad Y_1^0(\theta, \phi) = \sqrt{\frac{6}{8\pi}} \cos \theta,$$

show that

$$\langle 100 | z | 210 \rangle = 2^{5/2} (2/3)^5 a_0$$

Hence show that Einstein's  $A$  coefficient for the Lyman  $\alpha$  transition is

$$A = \frac{1}{3} 2^4 (2/3)^8 \pi \alpha^3 \nu.$$

[Hint: recall that  $\mathcal{R} = \frac{1}{2}\alpha^2 m_e c^2$  and  $a_0 = \hbar/\alpha m_e c$ .]

What is the characteristic timescale in ns for the radiative decay of an isolated hydrogen atom that starts from the  $n = 2$  level?

**2.6** With  $|nlm\rangle$  a stationary state of hydrogen, and given that

$$Y_1^0(\theta, \phi) = \sqrt{\frac{6}{8\pi}} \cos \theta \quad Y_1^1(\theta, \phi) = -\sqrt{\frac{3}{8\pi}} \sin \theta e^{i\phi} \quad Y_1^{-1}(\theta, \phi) = \sqrt{\frac{3}{8\pi}} \sin \theta e^{-i\phi},$$

show that

$$\begin{aligned} \langle 100|(x - iy)|211\rangle &= -\sqrt{2}\langle 100|z|210\rangle \\ \langle 100|(x - iy)|21 - 1\rangle &= 0. \end{aligned}$$

Write down the values of  $\langle 100|(x + iy)|21 - 1\rangle$  and  $\langle 100|(x + iy)|211\rangle$  and hence show that with

$$|\psi\rangle \equiv \frac{1}{\sqrt{2}}(|211\rangle - |21 - 1\rangle),$$

$\langle 100|x|\psi\rangle = -\langle 100|z|210\rangle$ . Explain the physical significance of this result.

**2.7** Derive the selection rules

$$\begin{aligned} \langle n'l'm'|x_+|nlm\rangle &= 0 \quad \text{unless} \quad m' = m + 1 \\ \langle n'l'm'|x_-|nlm\rangle &= 0 \quad \text{unless} \quad m' = m - 1. \end{aligned}$$

where  $x_{\pm} = x \pm iy$ . From this selection rule one infers that when the atom sits in a magnetic field along the  $z$  axis and the spectrometer looks along the  $z$  axis, the detected photons will be circularly polarised. Show that linearly polarised photons *can* be detected from an atom that's in a magnetic field.

From the above rules it might be argued that photons emitted along the  $z$  axis will be circularly polarised even in the absence of a magnetic field. Why is this argument bogus?