Hilary Term 2012

## Nonabelian Quantum Field Theory

## Problem set A (due February 16)

P1) Using the Fourier mode expansion for the Dirac field,

$$\psi(x) = \sum_{s} \int \frac{d^3k}{(2\pi)^3 2k^0} \left( a_s(k)u_s(k)e^{-ikx} + b_s^{\dagger}(k)v_s(k)e^{ikx} \right)$$

obtain the expression of the hamiltonian

$$H = \int d^3x \,\overline{\psi} \left( -i\gamma^j \partial_j + m \right) \psi$$

in terms of  $a_s(k)$ ,  $b_s(k)$ , and discuss the role of a and b being anticommuting Grassmann variables.

P2) Let  $\chi$  be a Weyl spinor field. Show that if  $\chi$  transforms in the (1/2, 0) representation of the Lorentz group, then  $\varepsilon \chi^*$ , where  $\varepsilon = \begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix}$ , transforms in the (0, 1/2) representation.

P3) Write Dirac spinor  $\psi$  in terms of two left-handed Weyl spinors  $\xi$ ,  $\chi$  as

$$\psi = \begin{pmatrix} \xi \\ -\varepsilon \chi^* \end{pmatrix} \ .$$

a) Obtain the expression for Dirac mass lagrangian in terms of the two left-handed Weyl fields. b) For  $\psi$  belonging to representation R of an internal symmetry group, and  $\xi$  and  $\chi$  belonging to complex conjugate representations R and  $\overline{R}$ , show the invariance of the Dirac mass lagrangian under the internal symmetry.

P4) Consider the theory of a single left-handed Weyl spinor field  $\chi$ . (a) Write the Majorana mass and kinetic energy lagrangian, and corresponding equations of motion. (b) Consider internal symmetry transformations represented by unitary matrices U on  $\chi$ 

$$\chi \to U\chi$$
.

Show that the lagrangian mass term in part (a) is invariant under these transformations only if the field transforms in a real representation of the internal symmetry.

P5) Consider the quadratic part of the action for the gauge field  $A^a_{\mu}$  in  $n \cdot A = 0$  gauge:

$$S = \frac{1}{2} \int d^4x \ A^a_\mu(x) \left( \partial^2 g^{\mu\nu} - \partial^\mu \partial^\nu - \frac{1}{\xi} n^\mu n^\nu \right) A^a_\nu(x)$$

Show that the propagator in this gauge is given by

$$\widetilde{D}_{F\ \mu\nu}^{(n)\ ab}(k) = \frac{-i\delta^{ab}}{k^2 + i\varepsilon} \left[ g_{\mu\nu} - \frac{k_{\mu}n_{\nu} + k_{\nu}n_{\mu}}{k \cdot n} + (\xi k^2 + n^2)\frac{k_{\mu}k_{\nu}}{(k \cdot n)^2} \right]$$

P6) Consider the BRST transformations

$$\delta A^a_\mu = \lambda D^{ab}_\mu c^b \ , \quad \delta \psi = ig\lambda c^a t^a \psi \ , \quad \delta c^a = -\frac{1}{2}\lambda g f^{abc} c^b c^c$$

,

with  $\lambda$  Grassmann parameter. Let Q be the BRST charge operator. Evaluate the action of  $Q^2$  on the gauge, ghost and matter fields and verify that the charge operator is nilpotent,  $Q^2 = 0$ .

P7) Verify that the generalized Ward identity

$$\int dx \left[ \frac{\delta\Gamma}{\delta A^a_{\mu}} \frac{\delta\Gamma}{\delta u^{\mu a}} + \frac{\delta\Gamma}{\delta c^a} \frac{\delta\Gamma}{\delta v^a} - \frac{1}{g\xi} \frac{\delta\Gamma}{\delta \bar{c}^a} (\partial^{\mu} A^a_{\mu}) \right] = 0$$

can be rewritten in terms of the transformed generating functional  $\Gamma'$ ,

$$\Gamma = \Gamma' - \int dx \; \frac{1}{2\xi} \; (\partial^{\mu} A^a_{\mu})^2 \; .$$

in the form

$$\int dx \left[ \frac{\delta \Gamma'}{\delta u^{\mu a}} \frac{\delta \Gamma'}{\delta A^a_\mu} + \frac{\delta \Gamma'}{\delta v^a} \frac{\delta \Gamma'}{\delta c^a} \right] = 0 \ .$$

[Use the equations of motion to relate  $\delta\Gamma/\delta\bar{c}$  to the spacetime derivative of  $\delta\Gamma/\delta u$ .]

P8) Consider QED Compton scattering of scalar particles:



- (i) Determine the scattering matrix element.
- (ii) Show that it is gauge invariant.

(iii) Consider the laboratory frame in which the scalar particle is initially at rest:



Show that the differential cross section in the solid angle  $\Omega$  of the final photon momentum, at fixed photon polarizations  $\varepsilon$ ,  $\varepsilon'$ , is given by

$$\frac{d\sigma}{d\Omega} = \frac{\alpha^2}{m^2} \left(\frac{\omega'}{\omega}\right)^2 (\varepsilon \cdot \varepsilon')^2 ,$$

where m is the particle's mass, and  $\omega$  and  $\omega'$  are the initial and final photon energies.

(iv) Determine the unpolarized cross section by evaluating the average over polarizations of  $(\varepsilon \cdot \varepsilon')^2$ .

P9) Consider quark-gluon Compton scattering at lowest order in the QCD coupling g:



(i) Show that the sum of graphs (a) and (b) dotted into  $k^{\nu}$  gives

$$M^{(a)+(b)}_{\mu\nu} k^{\nu} = ig^2 [T^a, T^b] \overline{u}(p') \gamma_{\mu} u(p) .$$

(ii) Evaluate the contribution of graph (c) dotted into  $k^{\nu}$ ,

$$M^{(c)}_{\mu\nu} k^{\nu}$$

and show that the sum of all graphs gives  $M_{\mu\nu}k^{\nu} = 0$  provided  $\mu$  is restricted to physical polarizations. Contrast this with the abelian case of electrodynamics.