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The Wave Equation

Transverse displacements of an elastic string of linear density (kg/m) ρ



 $\therefore T$

 $T_1 \cos \theta_1 = T_2 \cos \theta_2$ for small θ , $\cos \theta \sim 1 \implies T_1 = T_2 = T$ Resolve horizontal forces : $T\sin\theta_2 - T\sin\theta_1 = \left(\rho\delta x\right)\frac{\partial^2 y}{\partial t^2}$ Resolve vertical forces $\left(\sin\theta \approx \tan\theta = \frac{\partial y}{\partial x}\right)$ $\left(\frac{\partial y}{\partial x}\right)$ $\left(\frac{\partial y}{\partial x}\right)_{1} = \rho \delta x \frac{\partial^2 y}{\partial t^2}$

$$T\left[\left(\frac{\partial y}{\partial x}\right)_2 - \left(\frac{\partial y}{\partial x}\right)_1\right] = \rho \delta x \frac{\partial^2 y}{\partial t^2}$$

$$\left(\frac{\partial y}{\partial x}\right)_2 = \left(\frac{\partial y}{\partial x}\right)_1 + \frac{\partial}{\partial x}\left(\frac{\partial y}{\partial x}\right)\delta x + \dots \implies T\left(\frac{\partial^2 y}{\partial x^2}\right)\delta x = \rho \frac{\partial^2 y}{\partial t^2}\delta x$$

$$\implies \frac{\partial^2 y}{\partial x^2} = \frac{\rho}{T} \frac{\partial^2 y}{\partial t^2} \equiv \frac{1}{c^2} \frac{\partial^2 y}{\partial t^2}$$

This is a WAVE EQUATION with velocity $c = \sqrt{T / \rho}$ (hence larger tension or lighter string leads to faster waves)

Remarks

• Style of analysis similar to study of N oscillating masses on a string; instead of N coordinates $y_p(t)$, p = 1, ..., Nnow we have y(x, t), x = continuous variable

• Motion given by partial differential equation (PDE) in x and t: the wave equation.

2nd-order linear PDE of hyperbolic type with constant coefficients

 \bullet Recall from N oscillators case: onset of wave behavior from looking at sinusoidal dependence on p

2nd-order linear PDEs

 $A(x, y)u_{xx} + 2B(x, y)u_{xy} + C(x, y)u_{yy} + D(x, y)u_x + E(x, y)u_y + F(x, y)u = R(x, y)u_{yy} + C(x, y)u_{yy} + D(x, y)u_{yy} + E(x, y)u_{yy} + F(x, y$

- $B^2 AC < 0$ elliptic. Ex.: Laplace eqn. $u_{xx} + u_{yy} = 0$
- $B^2 AC = 0$ parabolic. Ex.: diffusion eqn. $u_t \alpha u_{xx} = 0$
- $B^2 AC > 0$ hyperbolic. Ex.: wave eqn. $u_{tt} c^2 u_{xx} = 0$

♦ General solutions of PDEs depend on arbitrary functions
 (analogous to solutions of ODEs depending on arbitrary constants)
 → boundary conditions to determine such functions

• Cauchy boundary conditions: assign function u and normal derivative $\partial u/\partial n$ on given curve γ in xy plane

(relevant to hyperbolic PDEs)

D'Alembert's solution	$\frac{\partial^2 y}{\partial x^2} = \frac{1}{c^2} \frac{\partial^2 y}{\partial t^2}$
Change variables	u = x - ct $v = x + ct$
$\partial^2 v = 1 \ \partial^2 v = (\partial \partial \partial)^2 = 1$	$\frac{\partial}{\partial x} = \frac{\partial u}{\partial x}\frac{\partial}{\partial u} + \frac{\partial v}{\partial x}\frac{\partial}{\partial v} = \frac{\partial}{\partial u} + \frac{\partial}{\partial v}$ $\frac{\partial}{\partial t} = \frac{\partial u}{\partial t}\frac{\partial}{\partial u} + \frac{\partial v}{\partial t}\frac{\partial}{\partial v} = -c\frac{\partial}{\partial u} + c\frac{\partial}{\partial v}$ $\left(-\frac{\partial}{\partial t} - \frac{\partial}{\partial t}\right)^{2} - \left(-\frac{\partial}{\partial t} - \frac{\partial}{\partial t}\right)^{2}$
$\frac{\partial y}{\partial x^2} = \frac{1}{c^2} \frac{\partial y}{\partial t^2} \implies \left(\frac{\partial}{\partial u} + \frac{\partial}{\partial v}\right) y = \frac{1}{c^2}$ $\implies \frac{\partial^2 y}{\partial u \partial v} = \frac{\partial^2 y}{\partial u \partial v}$	$\frac{1}{2} \left(-c \frac{d}{\partial u} + c \frac{d}{\partial v} \right) y = \left(\frac{d}{\partial u} - \frac{d}{\partial v} \right) y$ $0 \qquad \qquad \text{f, g arbitrary functions}$
General solution $y(u,v) = f(u) + g(v)$ <i>i.e.</i> $y(x,t) = f(x-ct) + g(x+ct)$	

 \Diamond D'Alembert realizes that $y_{xx} = c^{-2}y_{tt}$ implies that y is a function of x - ct + a function of x + ct

 \Diamond Interpretation of D'Alembert analysis in terms of traveling waves:

• Take e.g. the part of the solution y = f(x - ct). At time $t = t_1$: $y(x, t_1) = f(x - ct_1)$ At time $t_2 > t_1$: $y(x, t_2) = f(x - ct_2) = f[x - ct_1 - c(t_2 - t_1)]$ $= f[(\underbrace{x - c(t_2 - t_1)}_{x'}) - ct_1]$

i.e., y at time t_2 and position x is the same as y was at time t_1 and position $x' = x - c(t_2 - t_1)$ shifted leftwards \Rightarrow wave has traveled to the right with speed c

Travelling waves y(x,

$$y(x,t) = f(x-ct)$$



Wave moves to right with speed c





Wave moves to left with speed c

Summary

• transverse displacements y(x,t) of elastic string

 \longrightarrow wave equation $y_{xx} = c^{-2}y_{tt}$

• hyperbolic 2nd-order linear PDE

• solution of wave equation y(x,t) = f(x-ct) + g(x+ct)

 \checkmark \nearrow traveling waves

 \bullet next: determine f and g from initial conditions