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Oxford - 23/02/2016



In collaboration with Óscar J. C. Dias, Ben Niehoff and Benson Way





# Seven Pillars of Black Hole Wisdom (sorry T. E. Lawrence):

- **1** Black holes have two Killing isometries: rigidity theorems.
- 2 Black holes have no hair: described by conserved charges.
- **3** The laws of black hole mechanics/thermodynamics:

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  - a) 0<sup>th</sup> law: constant temperature rigidity theorems.
  - b) 1<sup>st</sup> law:  $dE = TdS + \Omega_i dJ^i$ .
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- **5** Asymptotically flat black holes are stable.
- 6 Cosmic Censorship protect us from naked singularities.
- If a gravitational system is linearly stable, it ought to be nonlinearly stable.



2 Seemingly different instabilities in AdS

3 Geons as special solutions

4 One black hole to interpolate them all and in the darkness bind them

# 5 Outlook

Motivation				
1 Spoiler alert:				

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In a longer talk, I would argue that all known SUSY black holes in  $AdS_5$  are nonlinearly unstable.

Seemingly different instabilities in AdS

# Superradiance

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Superradiance instability

• The Kerr-AdS<sub>4</sub> black hole (aka Carter solution - '68):

$$ds^{2} = -\frac{\Delta_{r}}{r^{2} + x^{2}} \left[ dt - (1 - x^{2}) d\phi \right]^{2} + \frac{\Delta_{x}}{r^{2} + x^{2}} \left[ dt - (1 + r^{2}) d\phi \right]^{2} + a^{2} (r^{2} + x^{2}) \left( \frac{dr^{2}}{\Delta_{r}} + \frac{dx^{2}}{\Delta_{x}} \right) ,$$

where

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- ∂<sub>t</sub> and ∂<sub>φ</sub> are commuting Killing fields; decompose perturbations in Fourier modes: e<sup>-iωt+imφ</sup>.
- Unstable if quasi-normal modes with  $Im(\omega) > 0$  exist.








Black holes with a single Killing vector field

Seemingly different instabilities in AdS



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# The nonlinear stability of AdS

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In particular, if a geodesically complete spacetime is perturbed, does it remain "complete"?

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#### Claim:

Some generic small (but finite) perturbations of AdS become large and eventually form black holes.

• The energy cascades from low to high frequency modes in a manner reminiscent of the onset of turbulence.

Black holes with a single Killing vector field — Seemingly different instabilities in AdS

## Heuristics

• AdS acts like a confining finite box. Any generic finite excitation which is added to this box might be expected to explore all configurations consistent with the conserved charges of AdS - including small black holes.

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  - These solutions are special since they are exactly periodic in time and invariant under a single continuous symmetry.
  - Geons are analogous to nonlinear gravitational plane waves.
- This Heuristic argument has been observed numerically for certain types of initial data, but fails for other types.

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- Certain types of initial data do not do this: do not seem to form black holes at late times! Balasubramanian et. al.
- Understand why special fine tuned solutions Geons exist.

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Unclear if they can have the same energy, *i.e.* coexist, with large AdS black holes!

Black holes with a single Killing vector field

One black hole to interpolate them all and in the darkness bind them

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### Surfing the Geon:

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Evades rigidity theorem because the only Killing field is the horizon generator!

• We have constructed these solutions: ten coupled 3D nonlinear partial differential equations of Elliptic type.

### Black resonators 1/3:

• One helical Killing field:  $\partial_T = \partial_t + \Omega_H \partial_\phi$ .
One black hole to interpolate them all and in the darkness bind them

# Black resonators 1/3:

- One helical Killing field:  $\partial_T = \partial_t + \Omega_H \partial_\phi$ .
- Their line element can be adapted to  $\partial_T$ :

$$ds^{2} = \frac{L^{2}}{(1-y^{2})^{2}} \left[ -y^{2} A \Delta_{y} \left( dT + y \chi_{1} dy \right)^{2} + \frac{4y_{+}^{2} B dy^{2}}{\Delta_{y}} + \frac{4y_{+}^{2} S_{1}}{2-x^{2}} \left( dx + yx \sqrt{2-x^{2}} \chi_{3} dy + y^{2} x \sqrt{2-x^{2}} \chi_{2} dT \right)^{2} + (1-x^{2})^{2} y_{+}^{2} S_{2} \left( d\Psi + y^{2} \Omega dT + \frac{x \sqrt{2-x^{2}} \chi_{4} dx}{1-x^{2}} + y \chi_{5} dy \right)^{2} \right]$$

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• 2D moduli space:

$$T \equiv \frac{1+3y_+^2}{4\pi y_+} \quad \text{and} \quad \varepsilon \equiv \int_0^\pi \mathrm{d}\phi \chi_4(0,1,\phi) \sin(m\,\phi) \, .$$

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• Bifurcating Killing sphere - Killing horizon generated by  $\partial_T$ .

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• Black resonators extend from the onset of superradiance instability to the Geons ('onset of turbulent instability').

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- Black resonators extend from the onset of superradiance instability to the Geons ('onset of turbulent instability').
- Black resonators exist in regions where the Kerr-AdS solution is beyond extremality.

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 When black resonators coexist with Kerr-AdS solutions, they have higher entropy - 2<sup>nd</sup> order phase transition.

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- When black resonators coexist with Kerr-AdS solutions, they have higher entropy  $2^{\rm nd}$  order phase transition.
- Their horizon is deformed along the  $\phi$  direction along which they rotate embedding in 3D spacetime  $\delta Z \equiv Z \overline{Z}$ .

# Possible Endpoint of the Superradiance Instability - 1/3 ?

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- However, recall that m = 2 becomes stable in a region where Kerr-AdS is unstable to perturbations with m > 2!
- In addition, the cloud of gravitons hair never backreacts very strongly on the geometry central black hole really looks like Kerr-AdS.
- Finally, higher *m* black resonators seem to have increasing entropy with increasing *m*.

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Conjecture: there is no endpoint -Dias. Horowitz and JES '11

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### Outlook:

- What is the field theory interpretation of this phenomenon?
- Can we make a connection with glassy physics?
- . . .

Outlook

# Thank You!