

From Full Stopping to Transparency in a Holographic Model of Heavy Ion Collisions

Jorge Casalderrey-Solana

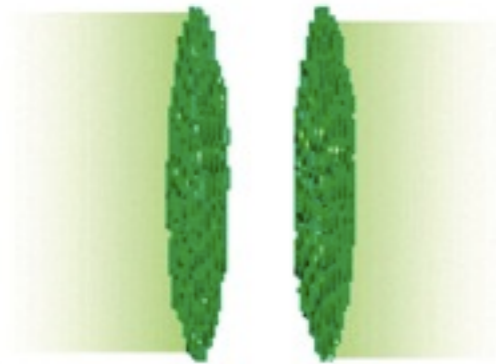
in collaboration with M. Heller, D. Mateos and W. van der Schee



UNIVERSITAT DE BARCELONA



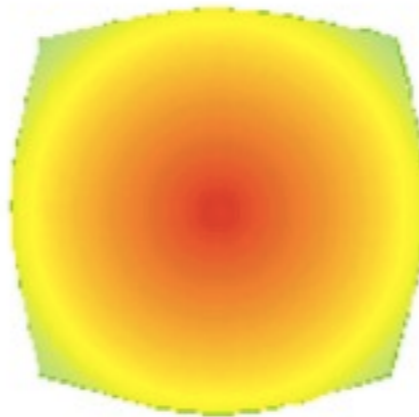
From Initial to Final State in Holography



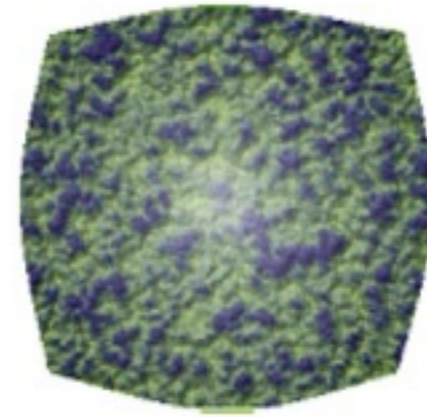
Initial State



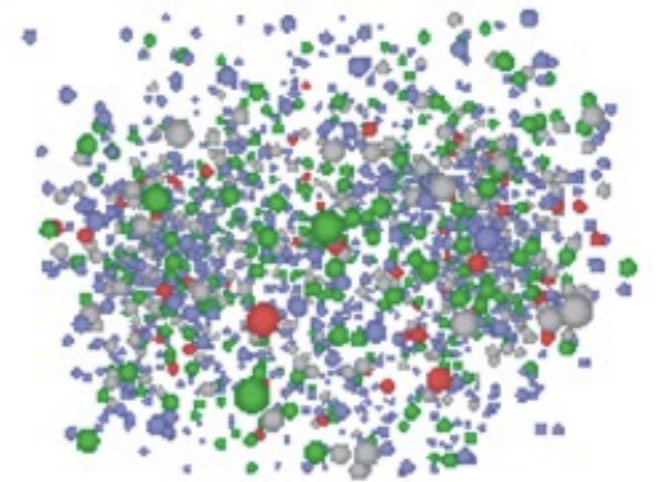
Collision



Out-of-equilibrium

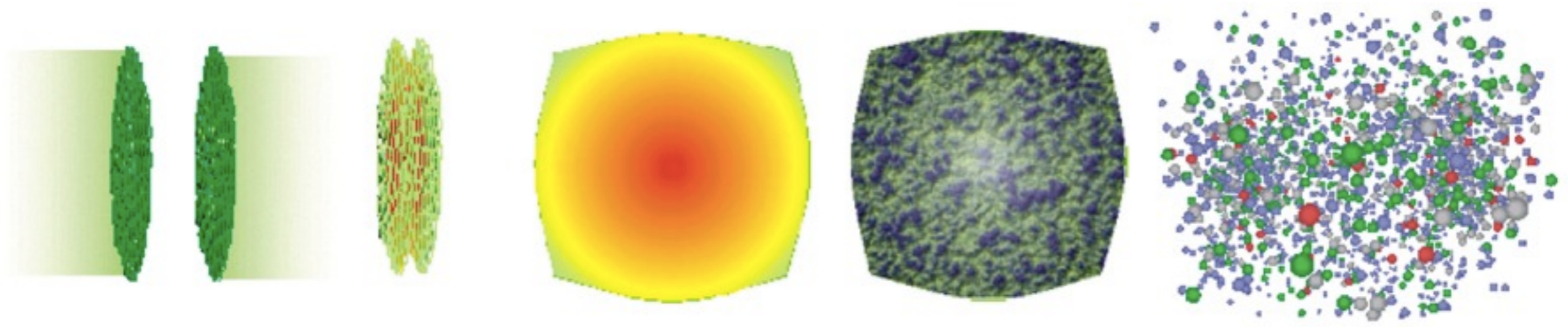


Collective expansion



Final stage

From Initial to Final State in Holography



Initial State

Collision

Out-of-equilibrium

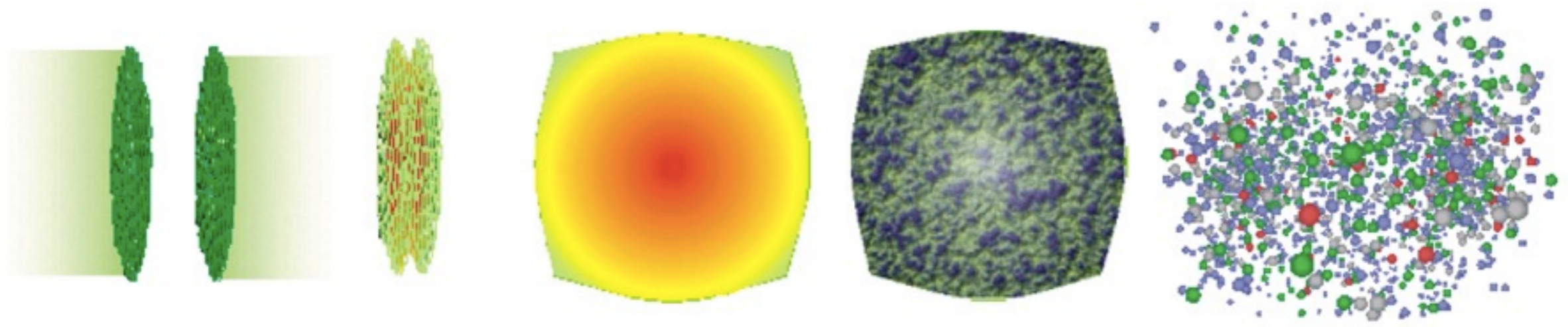
Collective expansion

Final stage



Properties of nuclei
Small x partons
Saturation effects

From Initial to Final State in Holography



Initial State

Collision

Out-of-equilibrium

Collective expansion

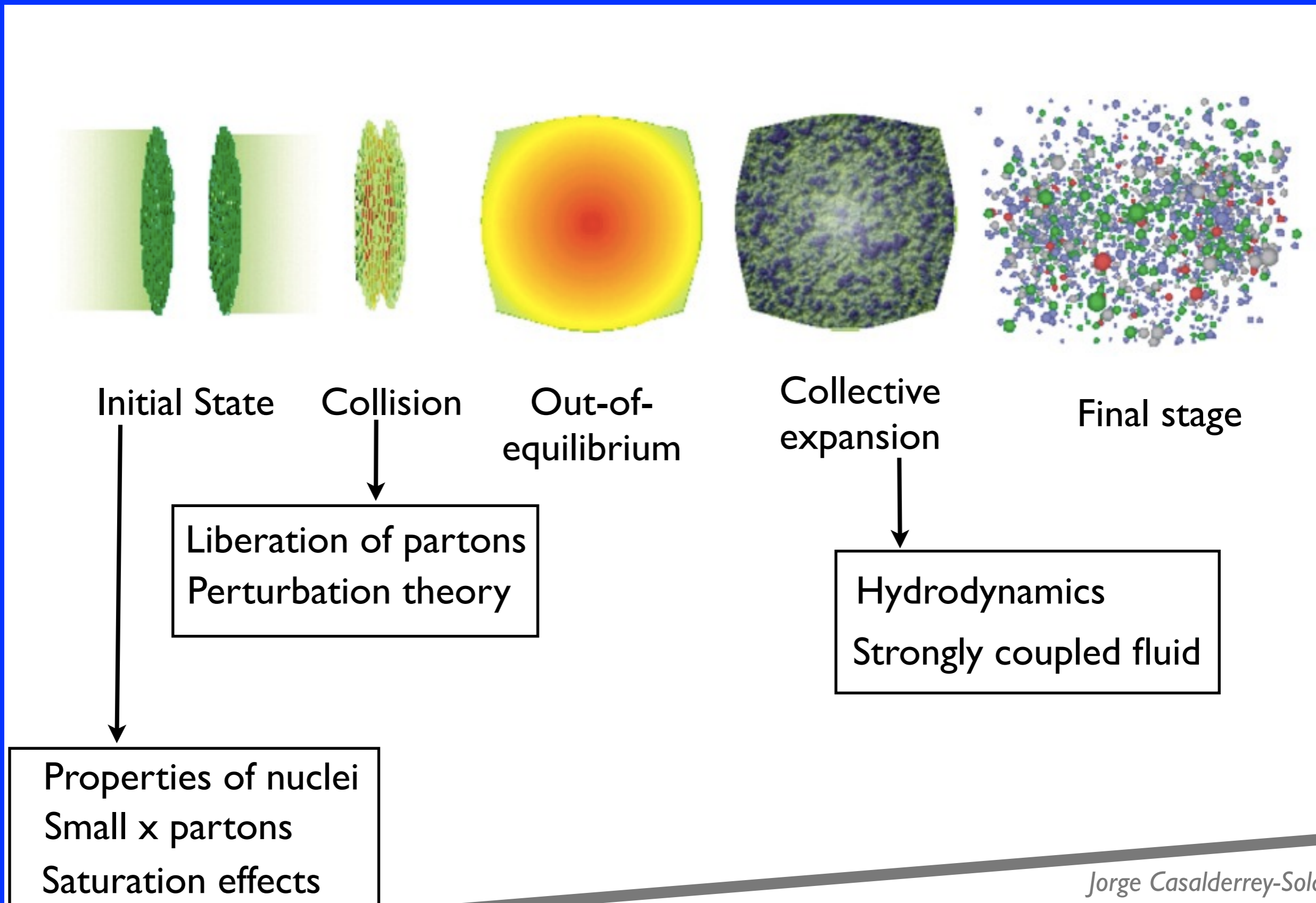
Final stage



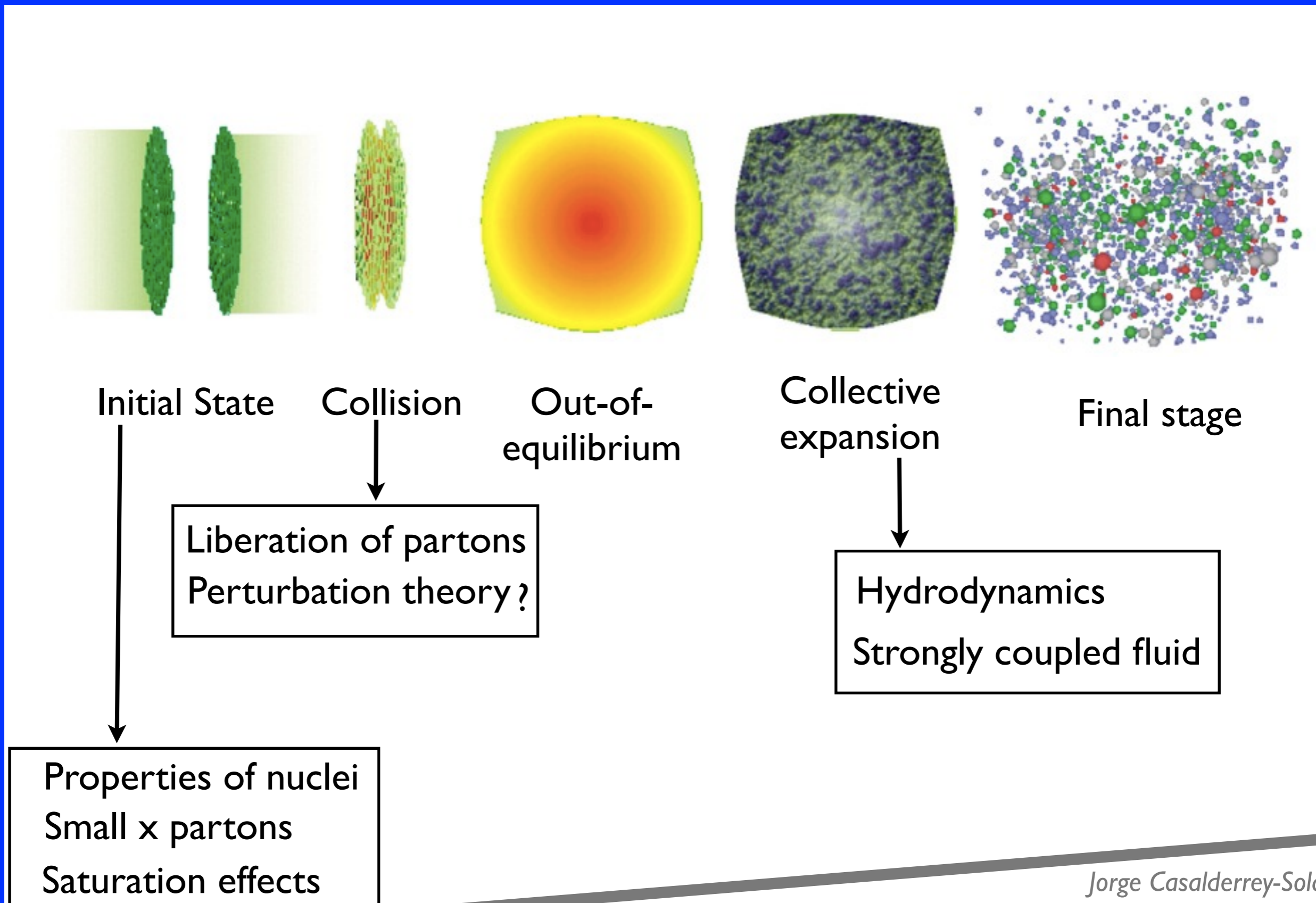
Properties of nuclei
Small x partons
Saturation effects

Hydrodynamics
Strongly coupled fluid

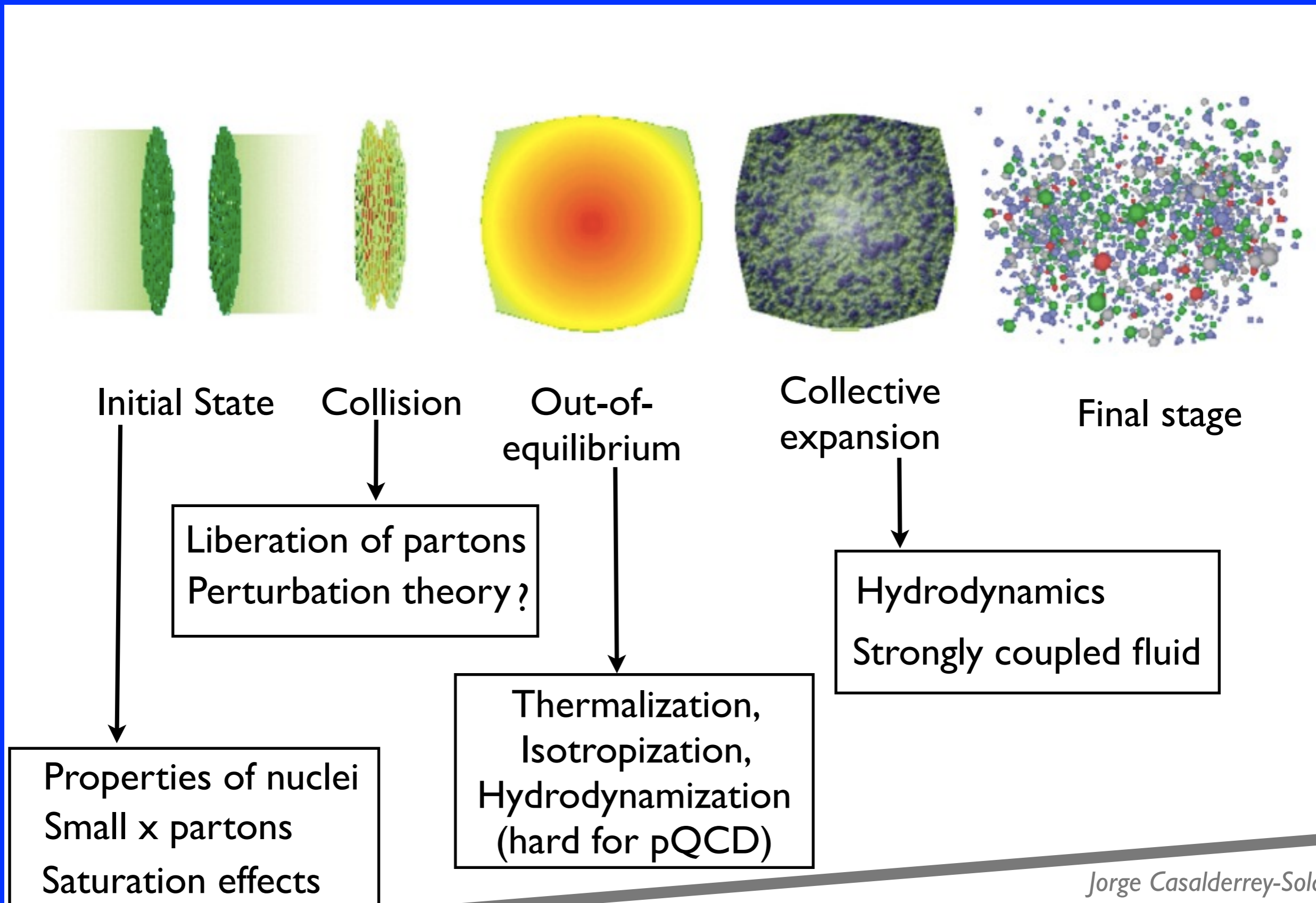
From Initial to Final State in Holography



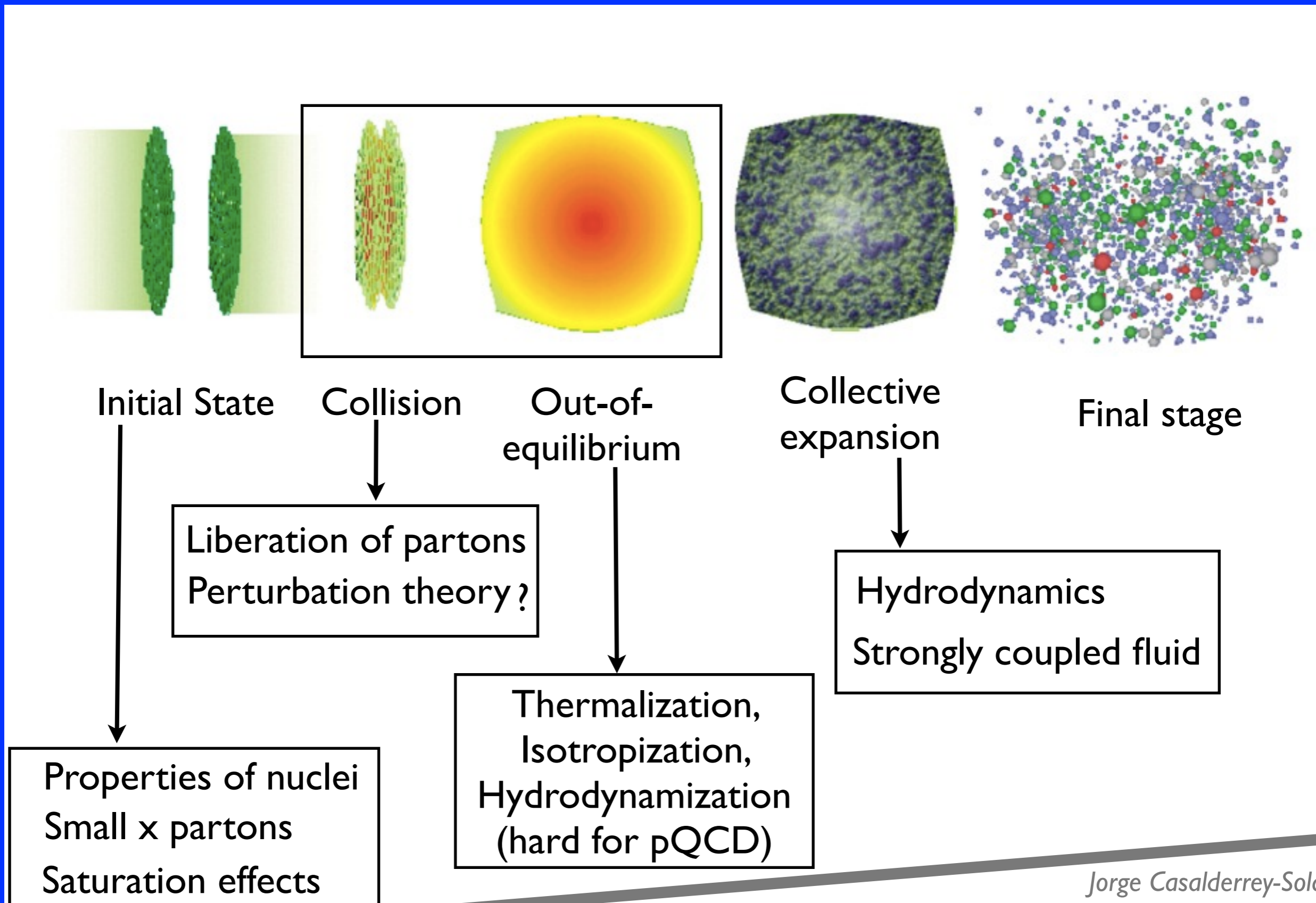
From Initial to Final State in Holography



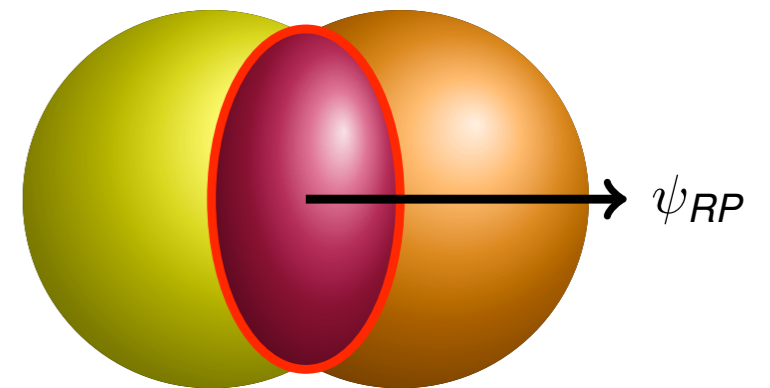
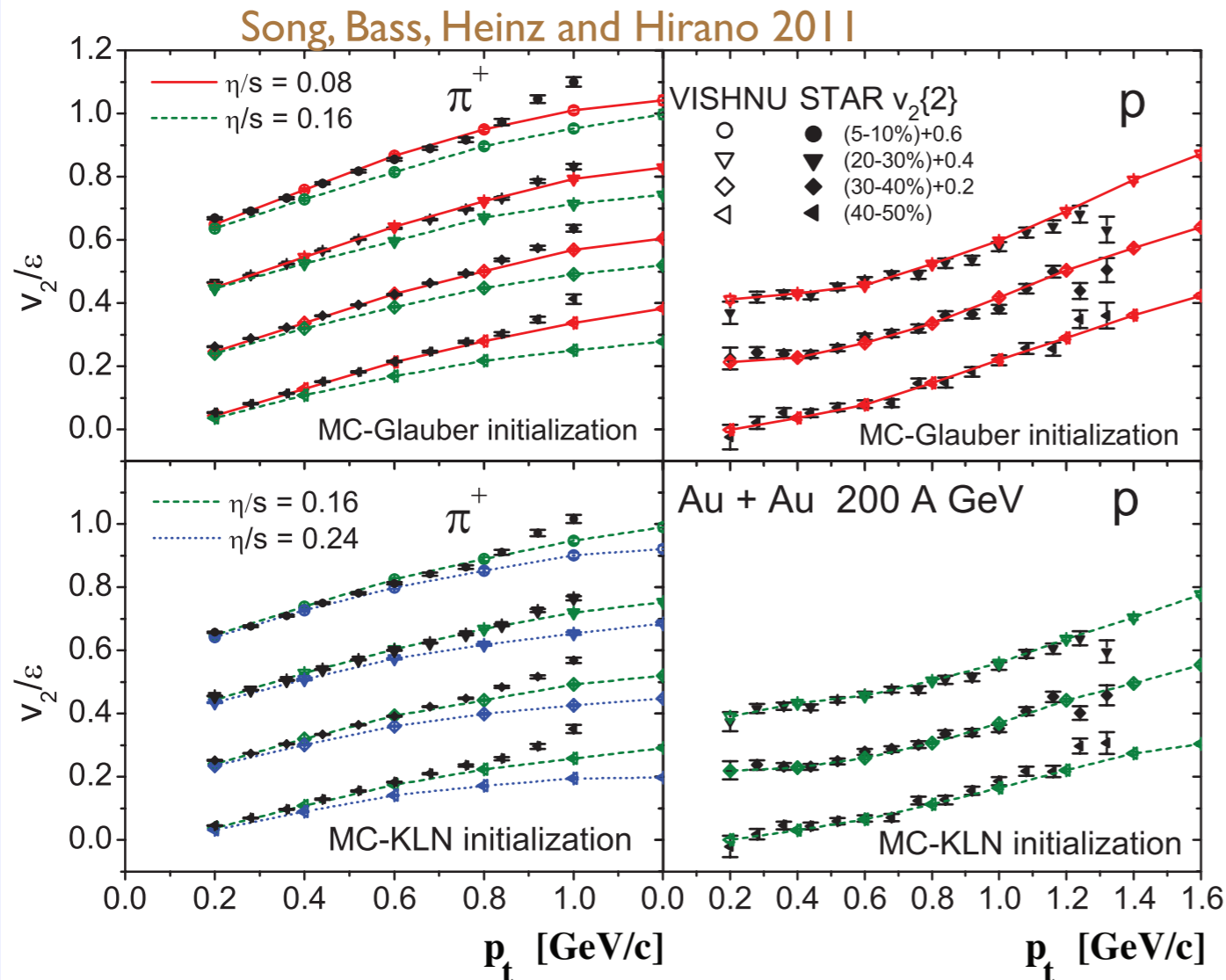
From Initial to Final State in Holography



From Initial to Final State in Holography



Hydrodynamics and Heavy Ion Collision



$$v_2 \equiv \langle \cos 2(\phi - \psi_{RP}) \rangle \propto \epsilon \equiv \frac{\langle y^2 \rangle - \langle x^2 \rangle}{\langle y^2 \rangle + \langle x^2 \rangle}$$

ϵ depends on the model for nuclear density

- Hydro describes well flow data for different species

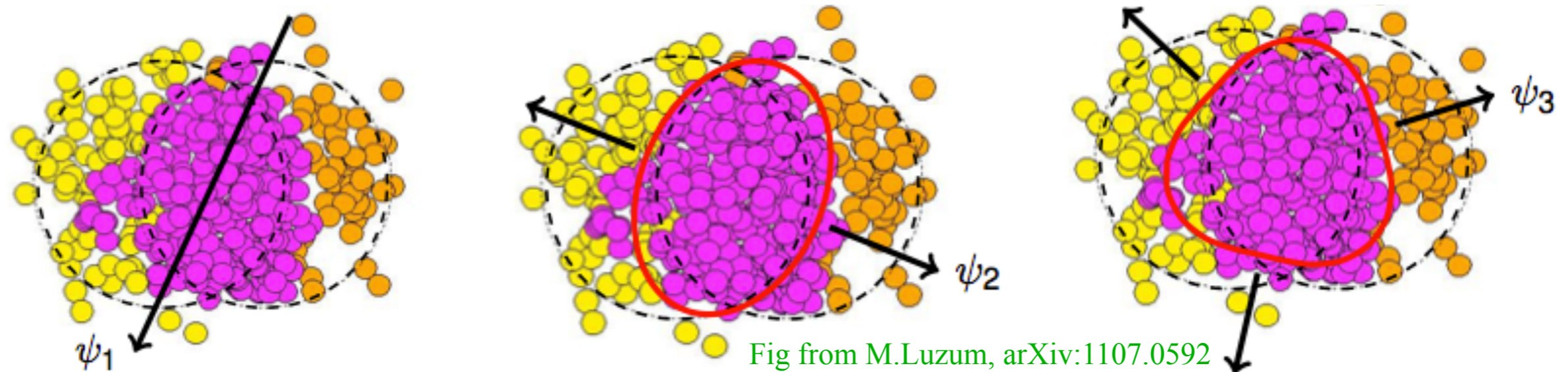
A very good liquid
strongly coupled!

$$1 < 4\pi \frac{\eta}{s} < 2.5$$

(Son, Bass, Heinz, Hirano, Shen 11)

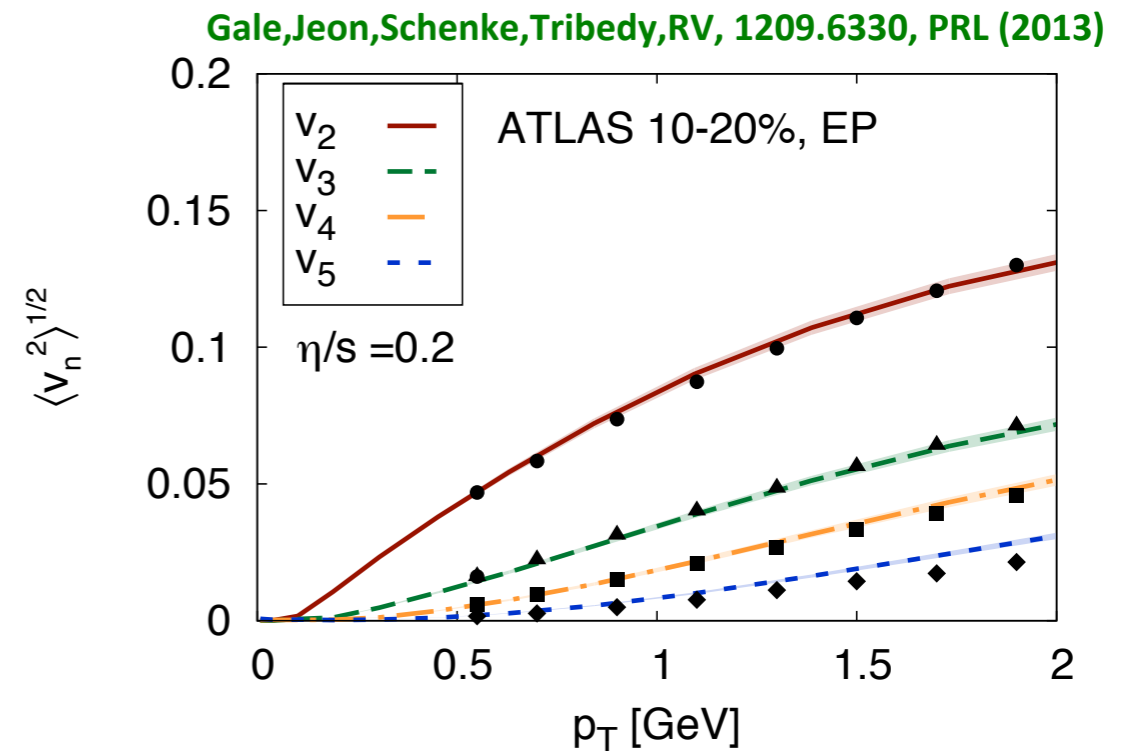
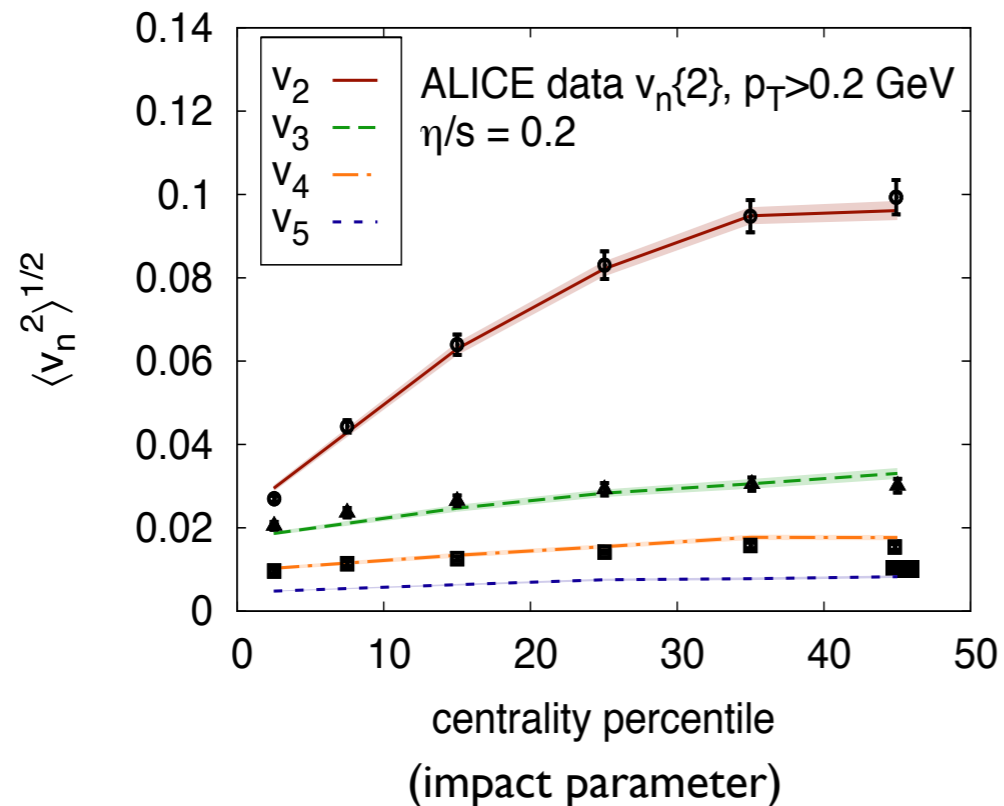
predicted by AdS/CFT!
(Policastro, Son, Starinets 01)

(Transverse) Fluctuations



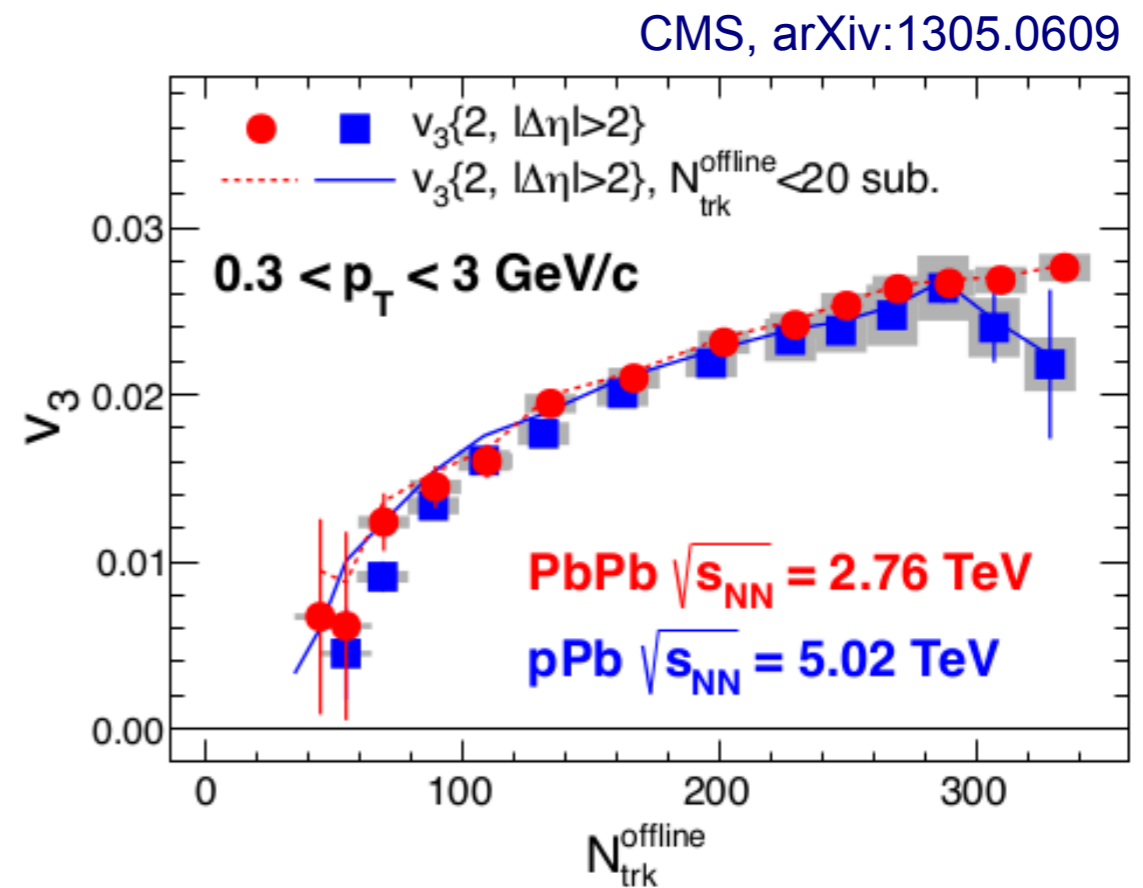
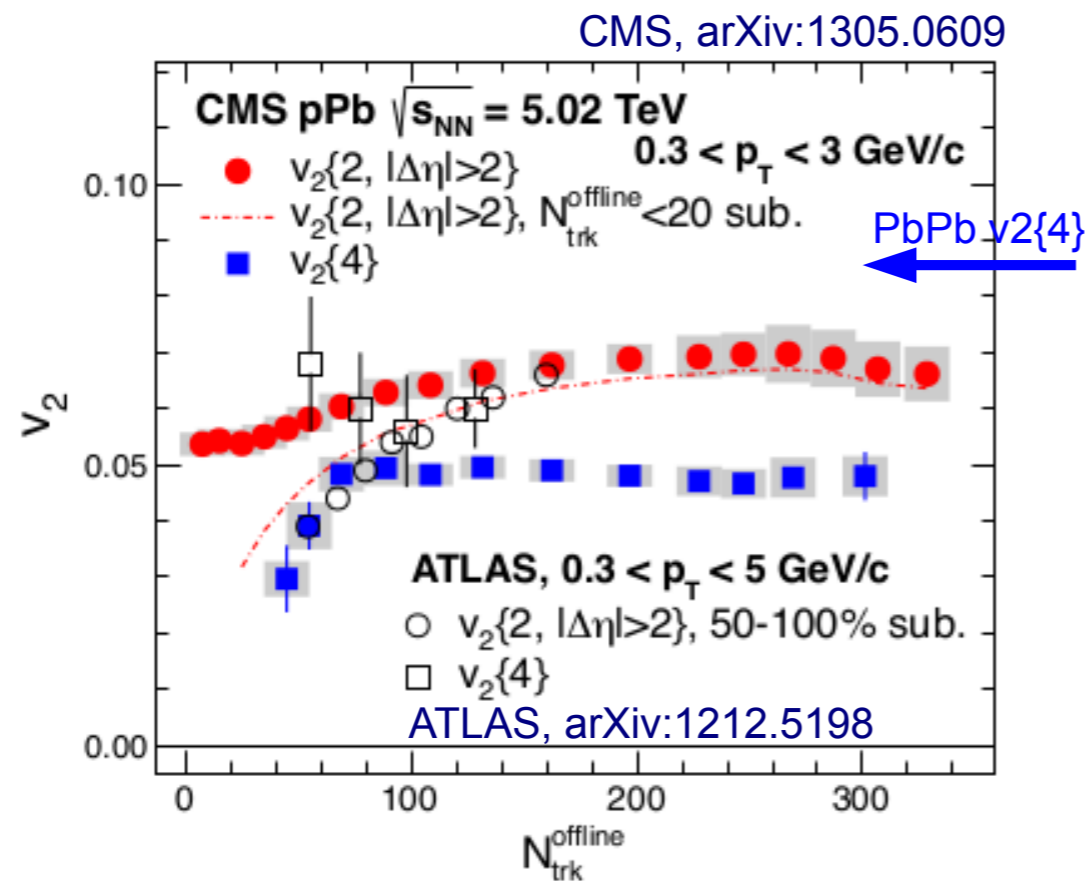
- Large event by event fluctuations \Rightarrow non-trivial hydro response

(Transverse) Fluctuations



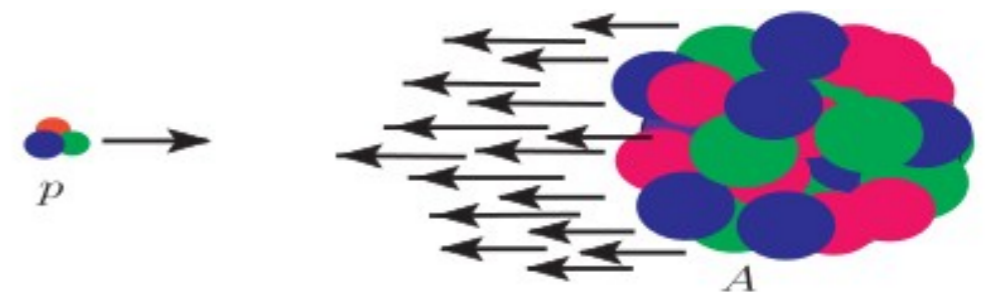
- Large event by event fluctuations \Rightarrow non-trivial hydro response
- Hydro + initial condition models can describe data
- New program to extract viscosities.
 - Focus on constraining initial state

Flow in p-Pb!



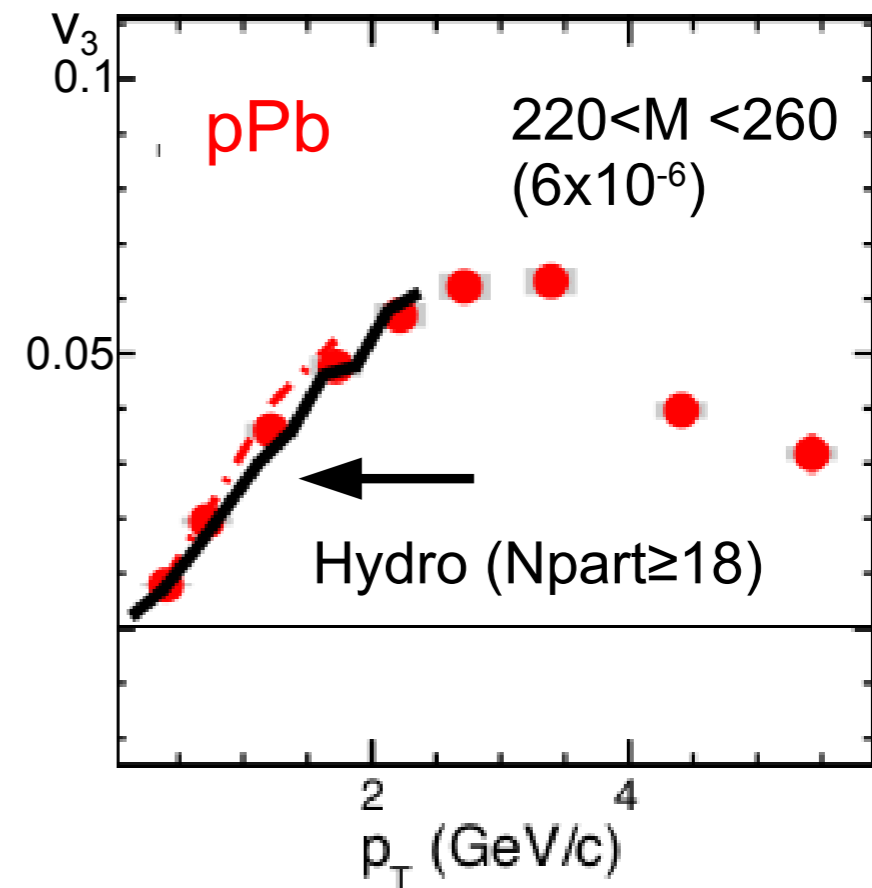
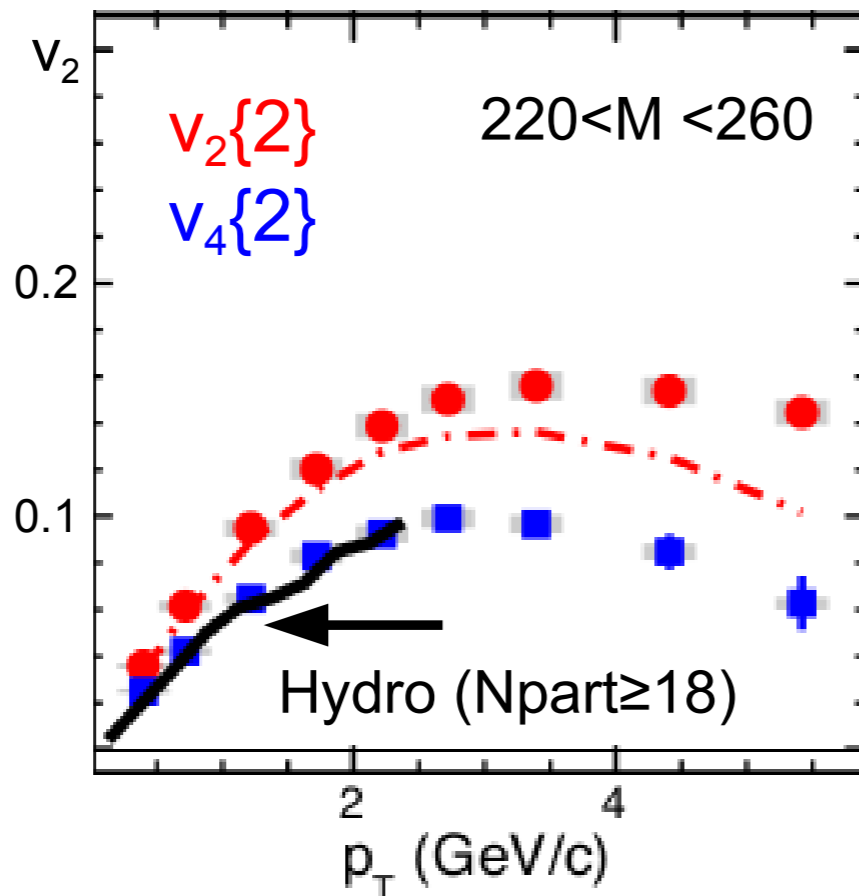
- Flow effects are large!

- v_2 smaller than PbPb (different shape)
- v_3 same as in PbPb (same fluctuations)



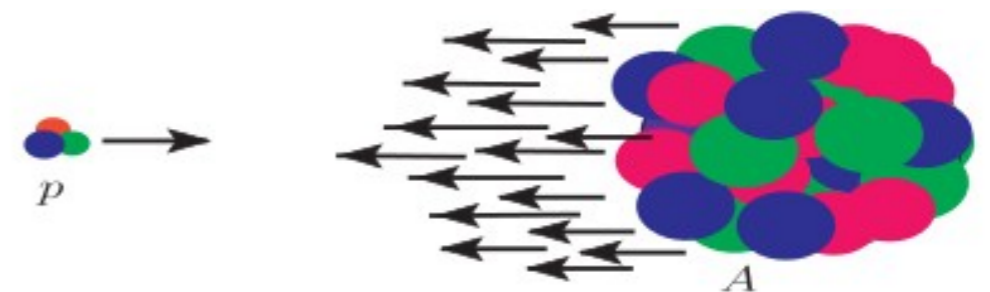
Flow in p-Pb!

Bozek arXiv:1112.0915

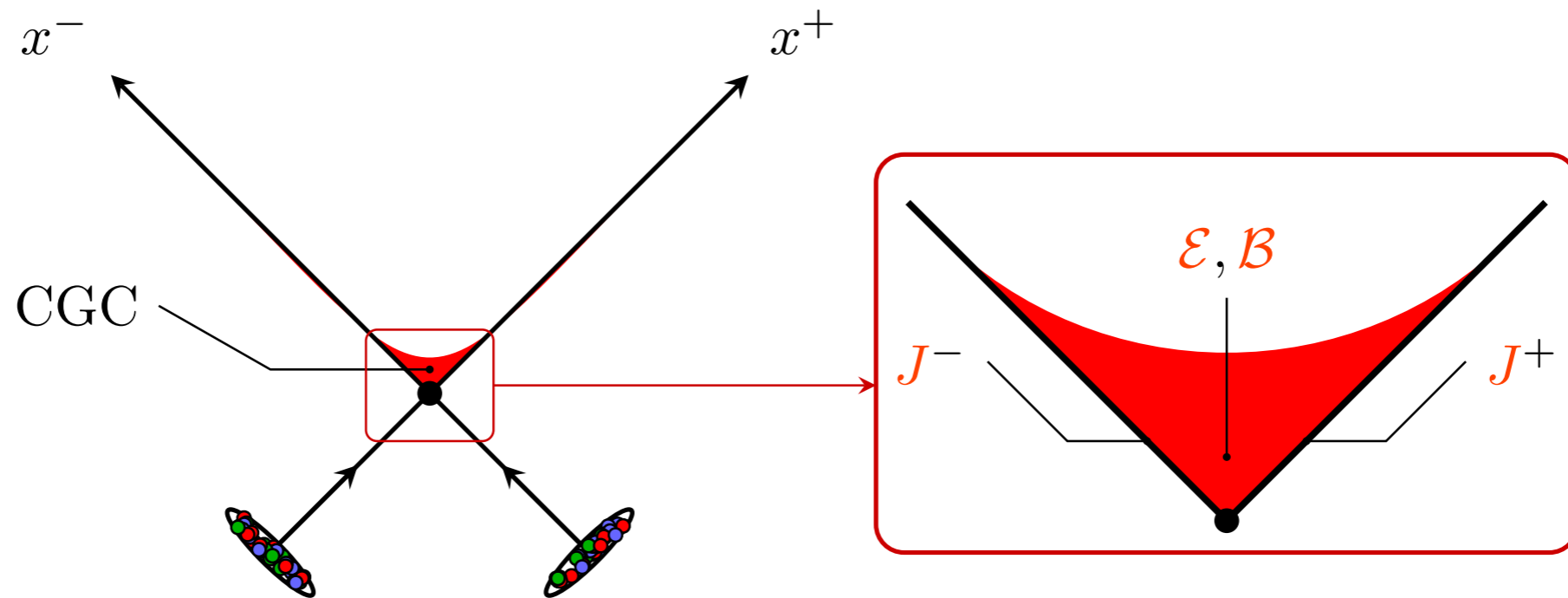


- Flow effects are large!
 - v_2 smaller than PbPb (different shape)
 - v_3 same as in PbPb (same fluctuations)
- Hydrodynamics predicts both these flows!

despite of smaller system \Rightarrow larger gradients



Initial Conditions and the CGC

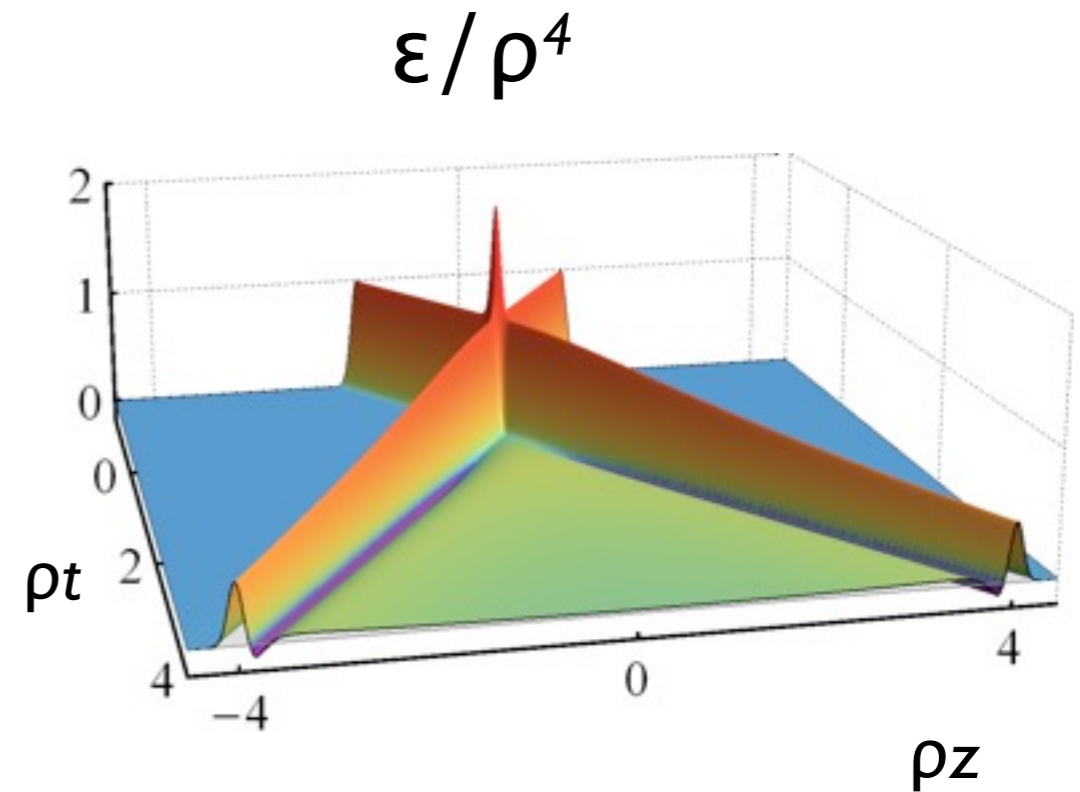
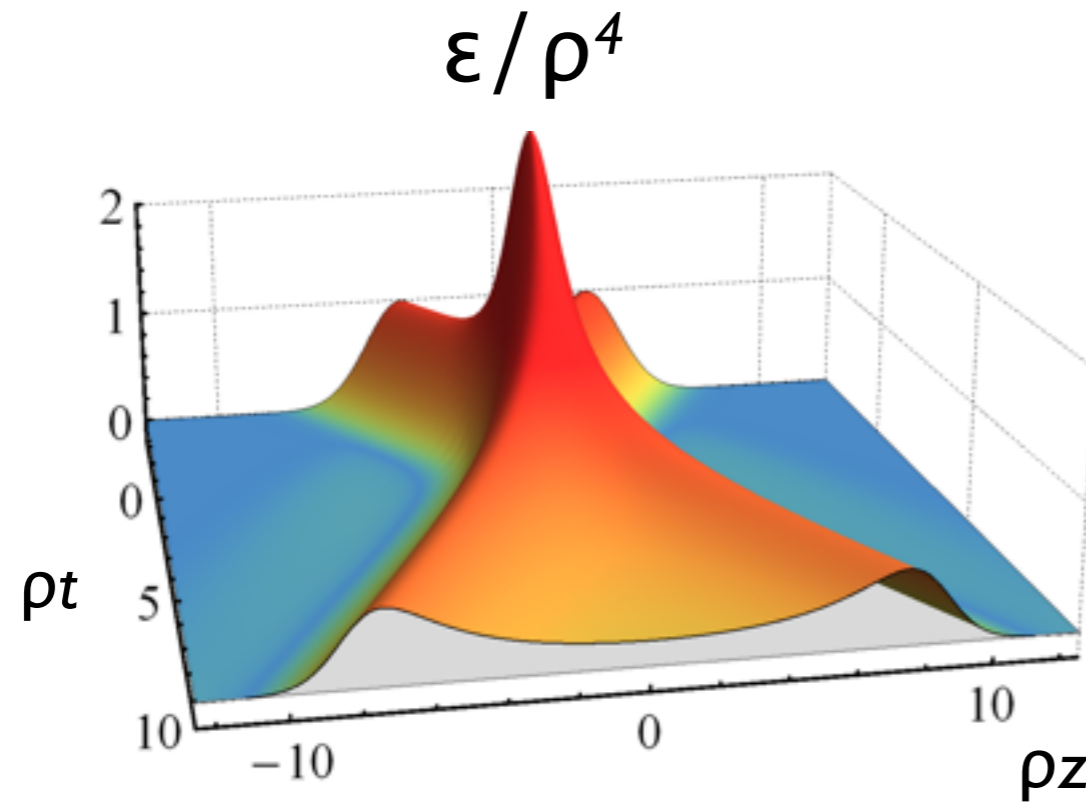


- Small x partons \Rightarrow Saturation physics
 - Typical size of partons $1/Q_s$ (perturbative scale)
 - Large occupation numbers $1/\alpha_s(Q_s) \Rightarrow$ classical fields
- Phenomenologically:

$$Q_s^{LHC} \sim 3 - 4 \text{ GeV} \quad \text{Albacete, Dumitru, Fujii, Nara 12}$$

- Not terribly perturbative, still room for strong coupling effects

Shock Collisions at Different Energies



Holographic collision of two gaussian shocks

ρ : maximum energy density

ω : width of the gaussian

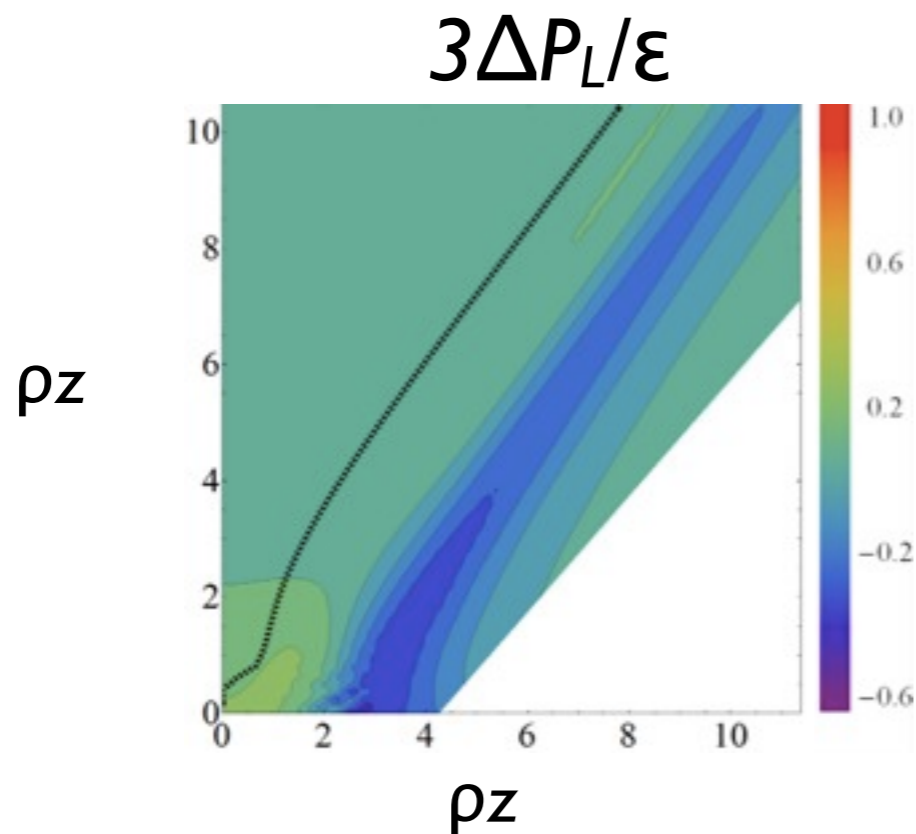
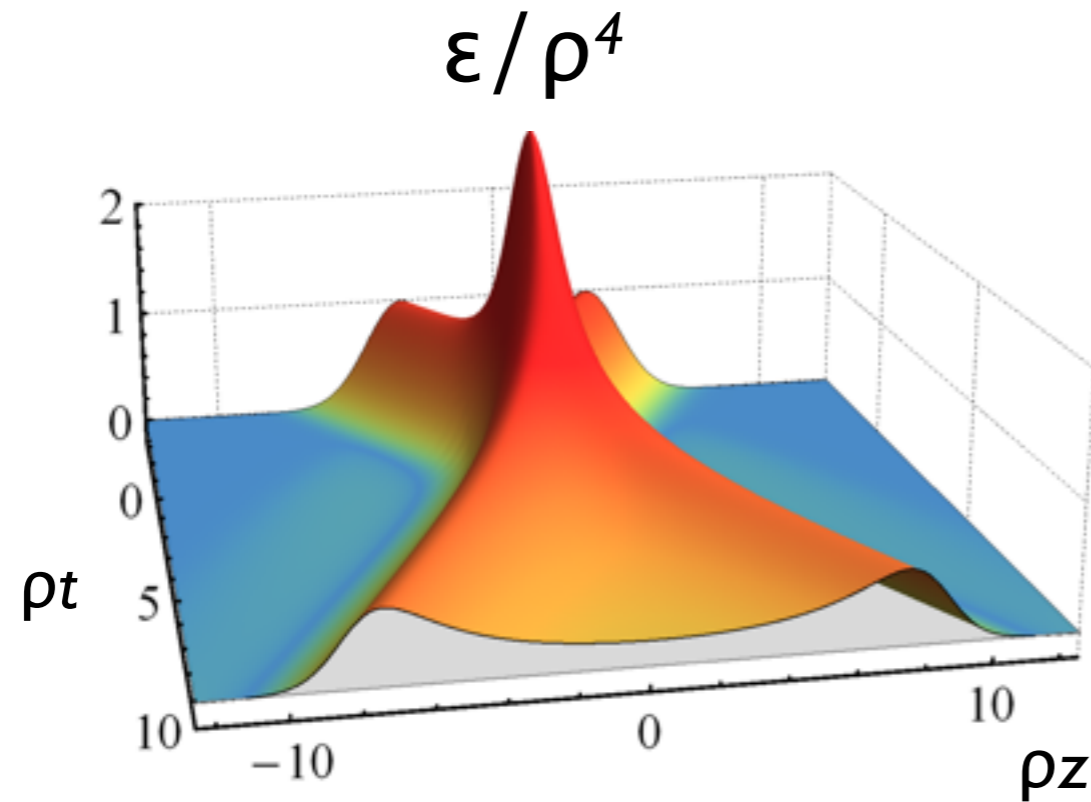
proxy for shock energy $\omega \sim 1/E$

μ : energy per transverse area

(Chesler & Yaffe 11)

*(Albacete, Kovchegov and Taliotis 08,
Grumiller & Romatschke 08)*

Low Energy Shocks



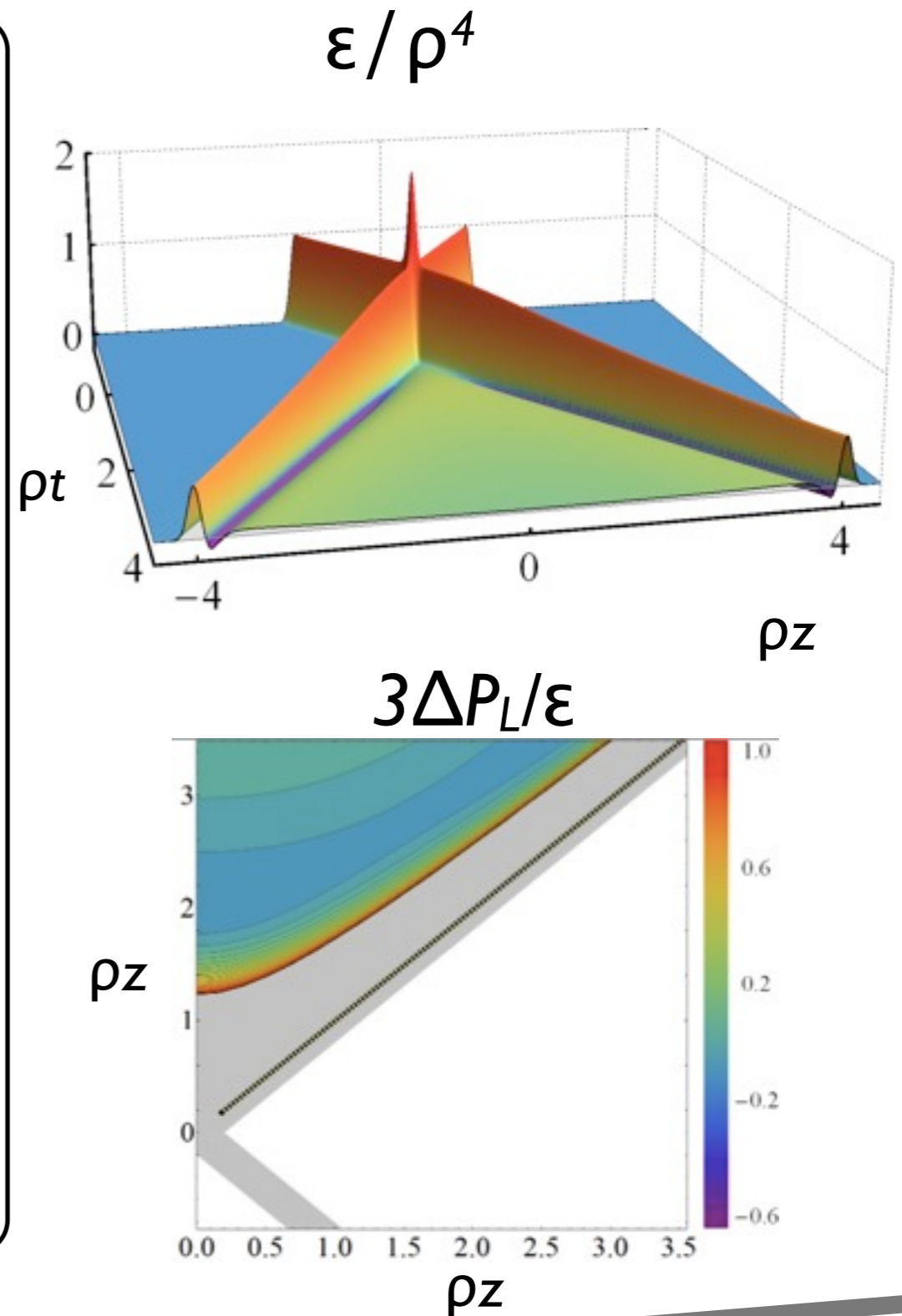
- The shocks merge and stall
- No propagation in light-cone
- Dynamics well approximated by viscous hydrodynamics

High energy shocks

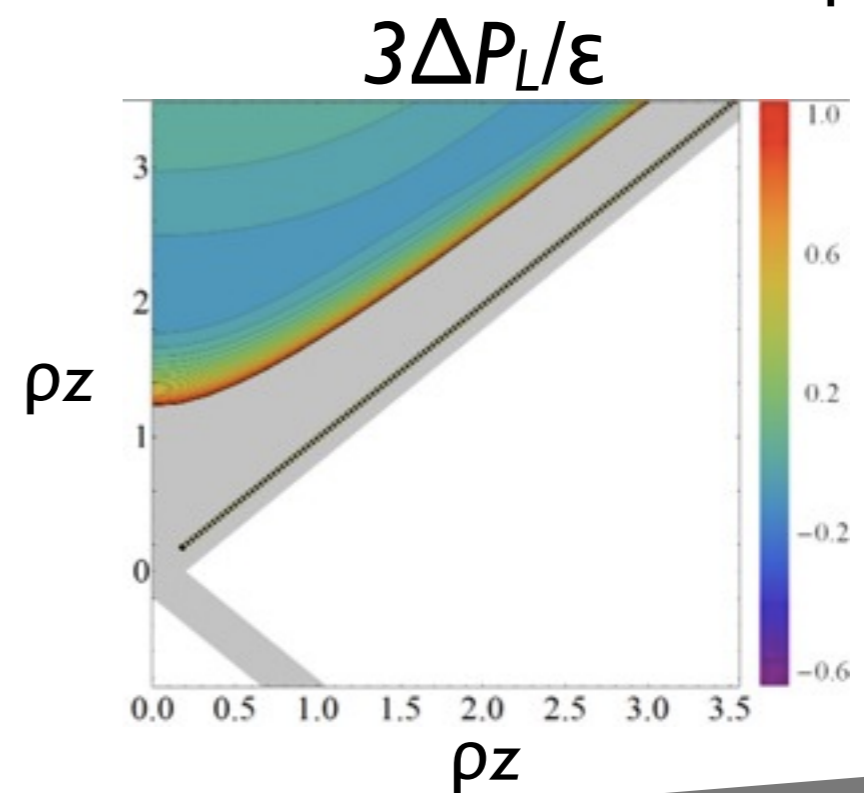
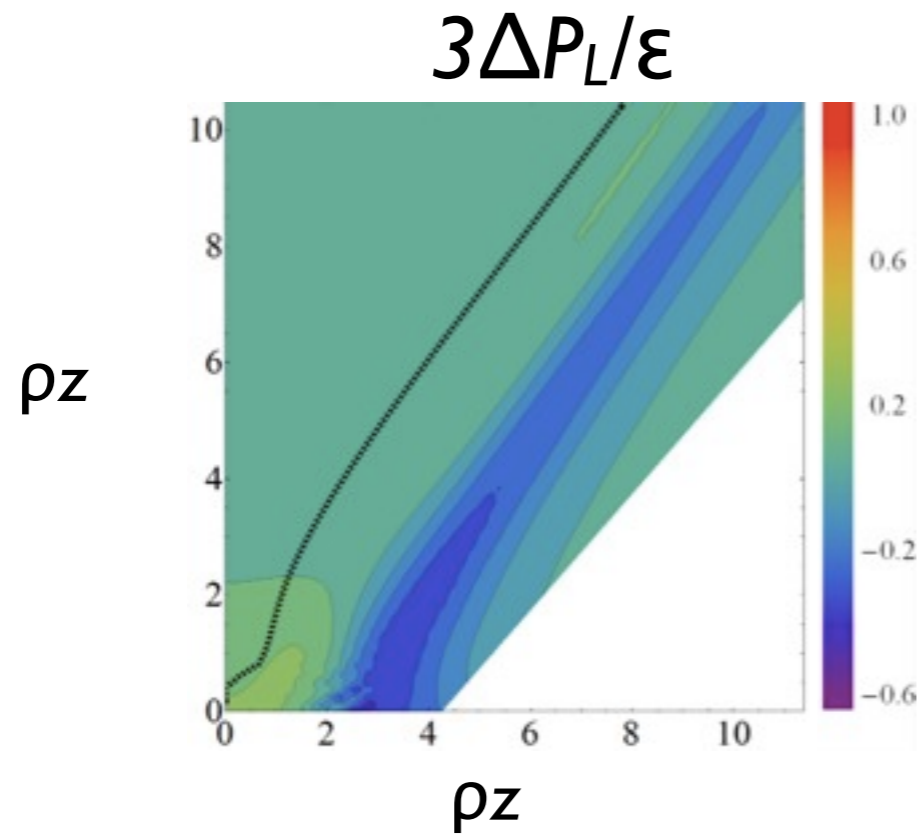
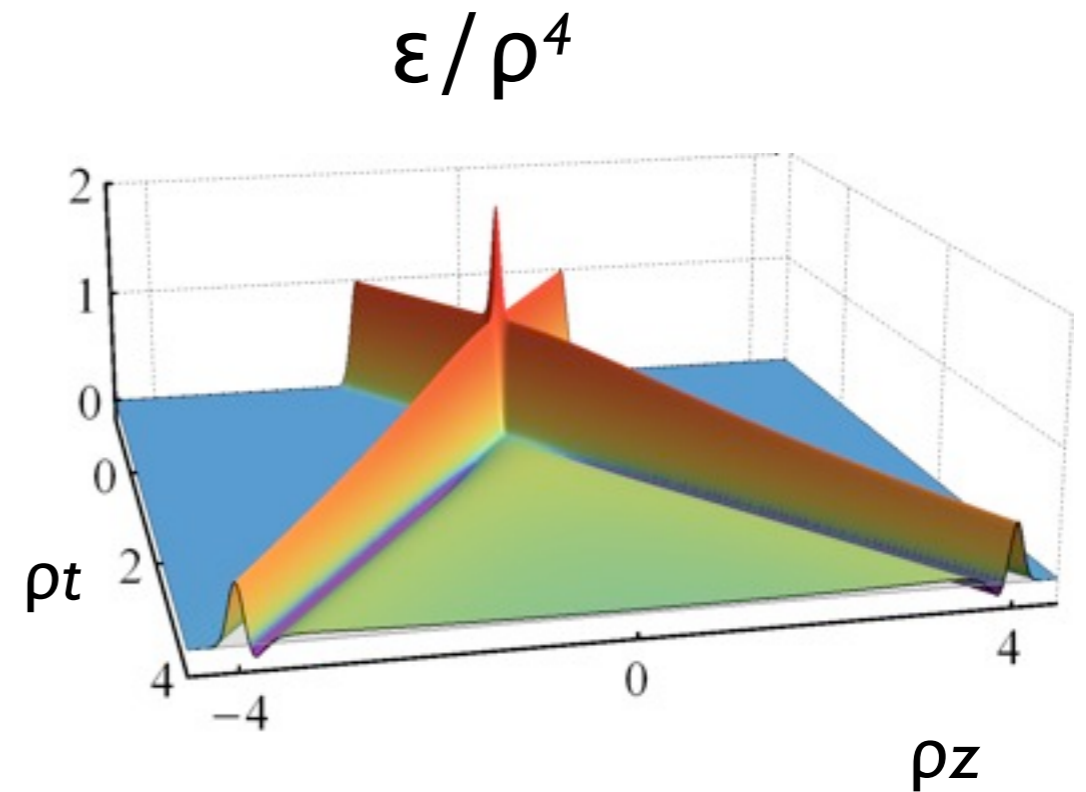
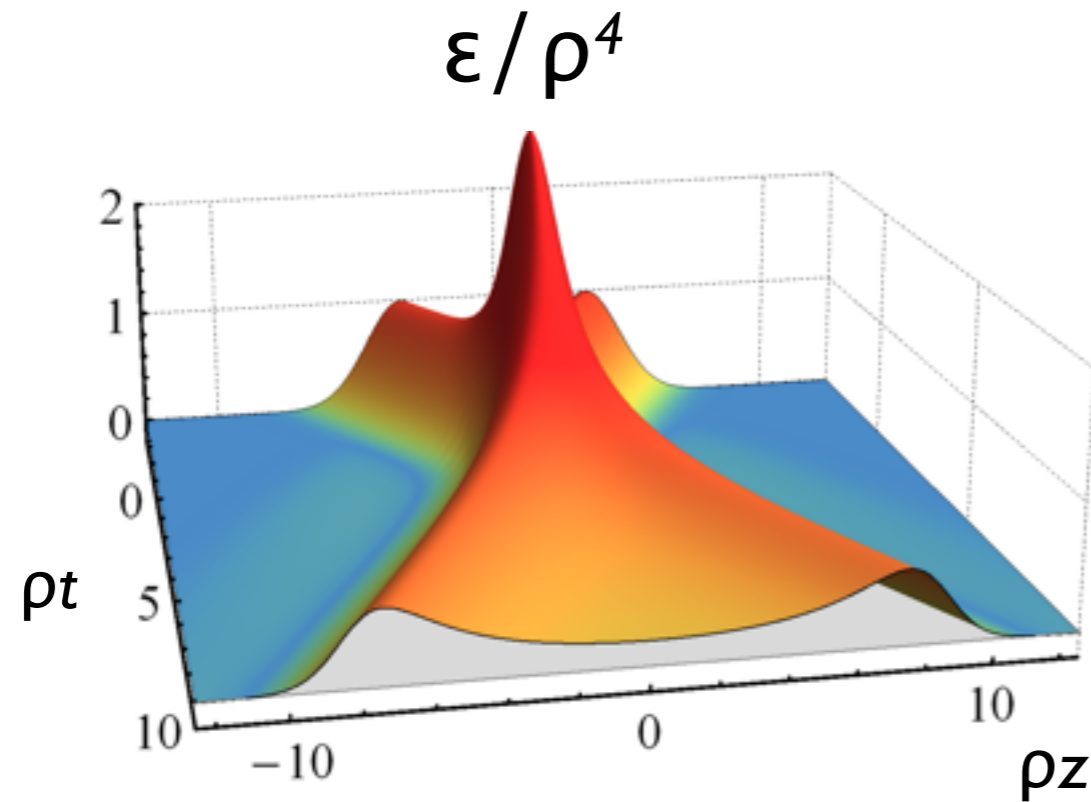
- Shocks pass through each other

Transparency

- Some energy remains in the light-cone
- Disturbed remnant shocks do not behave hydrodynamically
- They decay after the collision
- Hydrodynamic plasma is formed in the central rapidity region

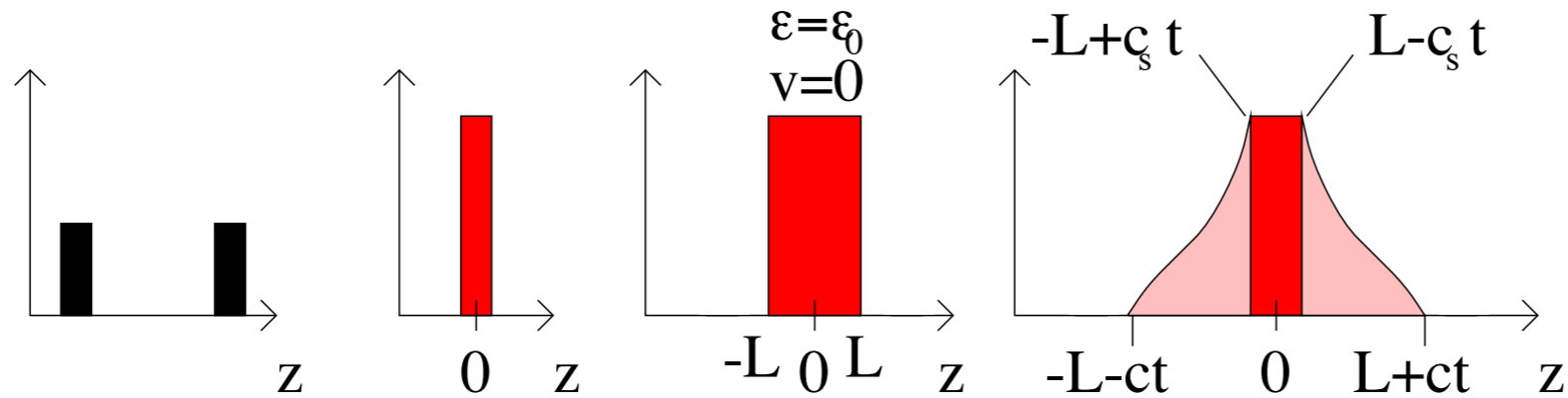


A Dynamical Cross Over



Landau vs Bjorken

- Landau model: All energy is stopped + hydro explosion

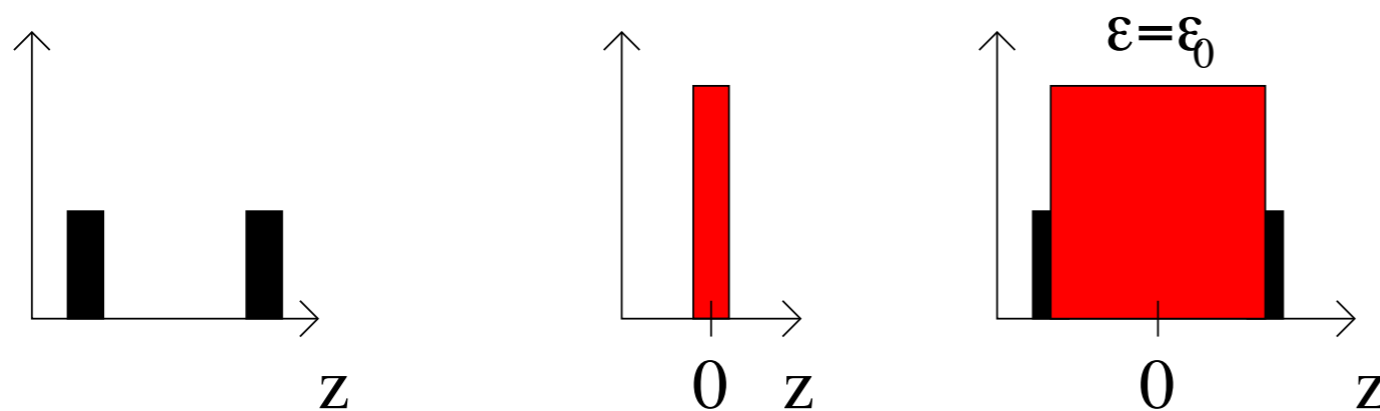


(predicts a gaussian distribution in rapidity)

$$t = \frac{z}{c} \text{Cosh}(\eta_{st})$$

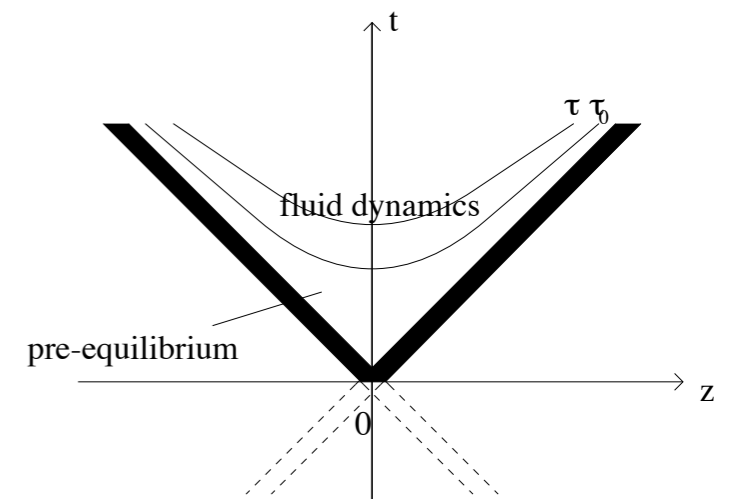
$$z = \frac{ct}{\text{Cosh}(\eta_{st})} \text{Sinh}(\eta_{st})$$

- Bjorken: matter produced at all rapidities

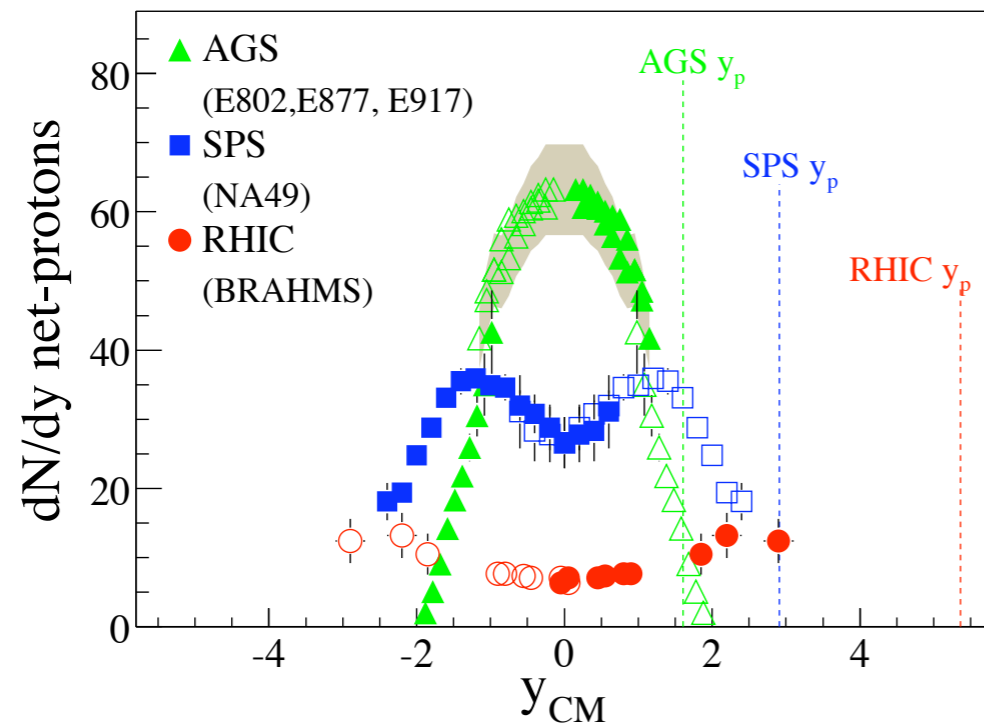


Boost invariant flow

Fluid rapidity = space time rapidity

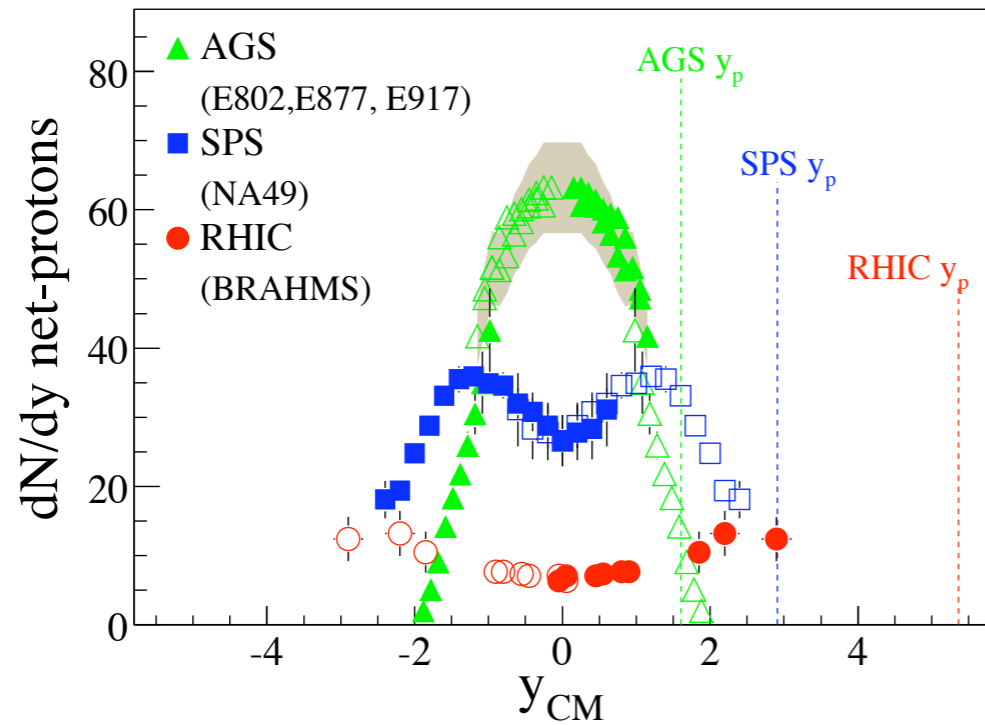


Landau vs Bjorken



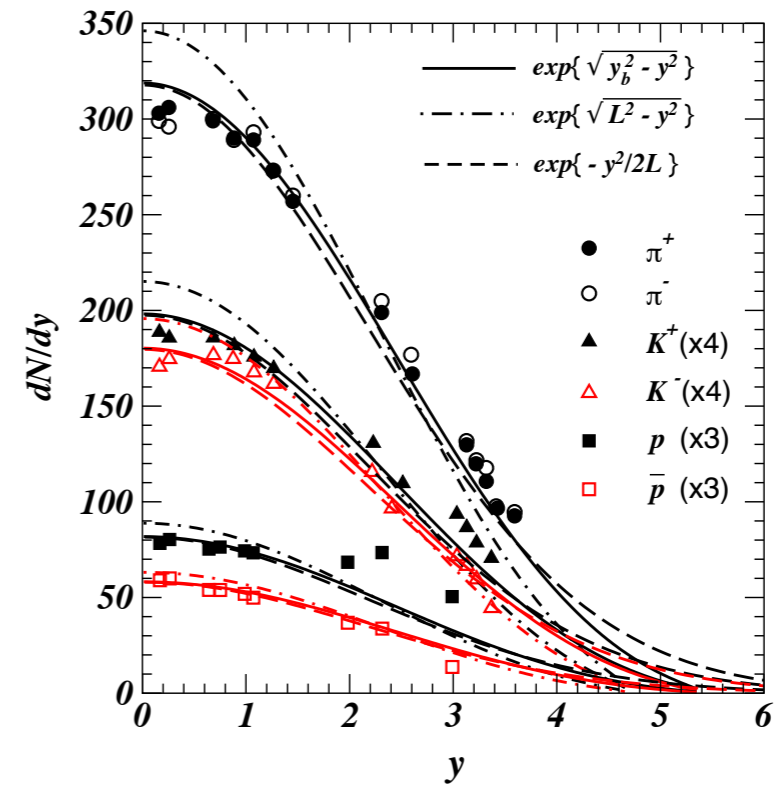
Baryons are not stopped
(Bjorken)

Landau vs Bjorken



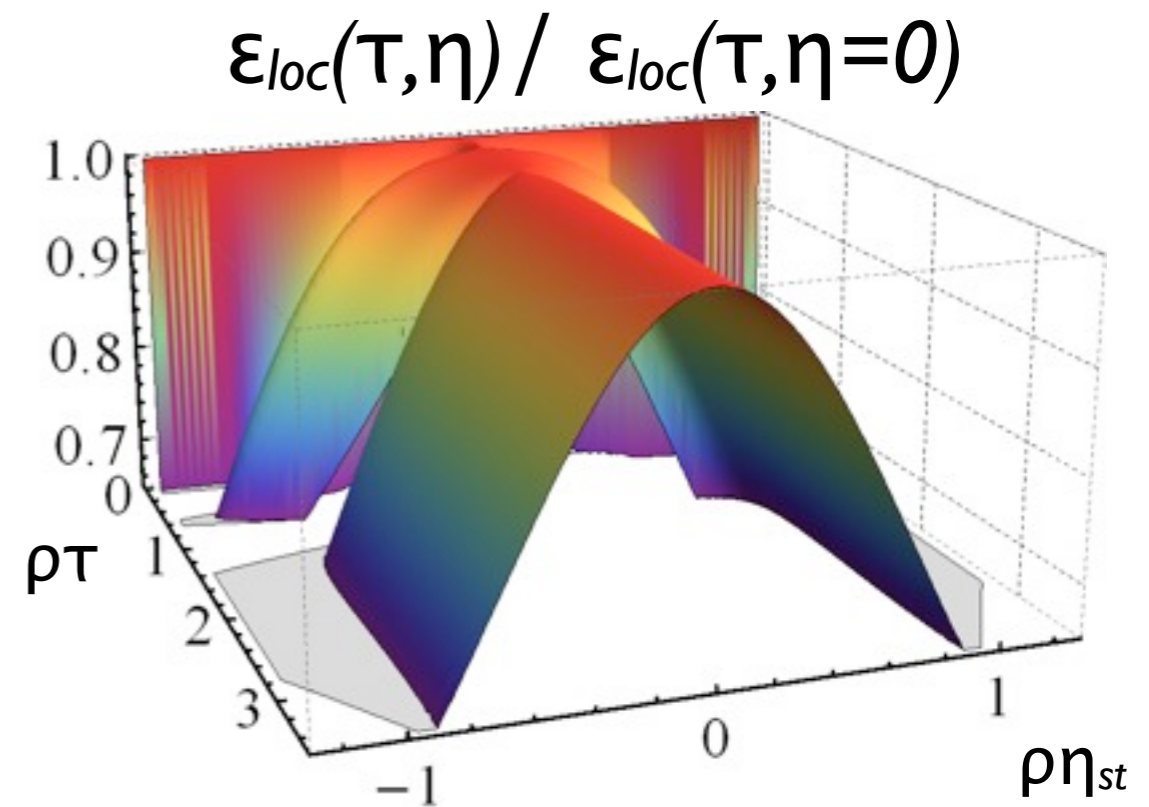
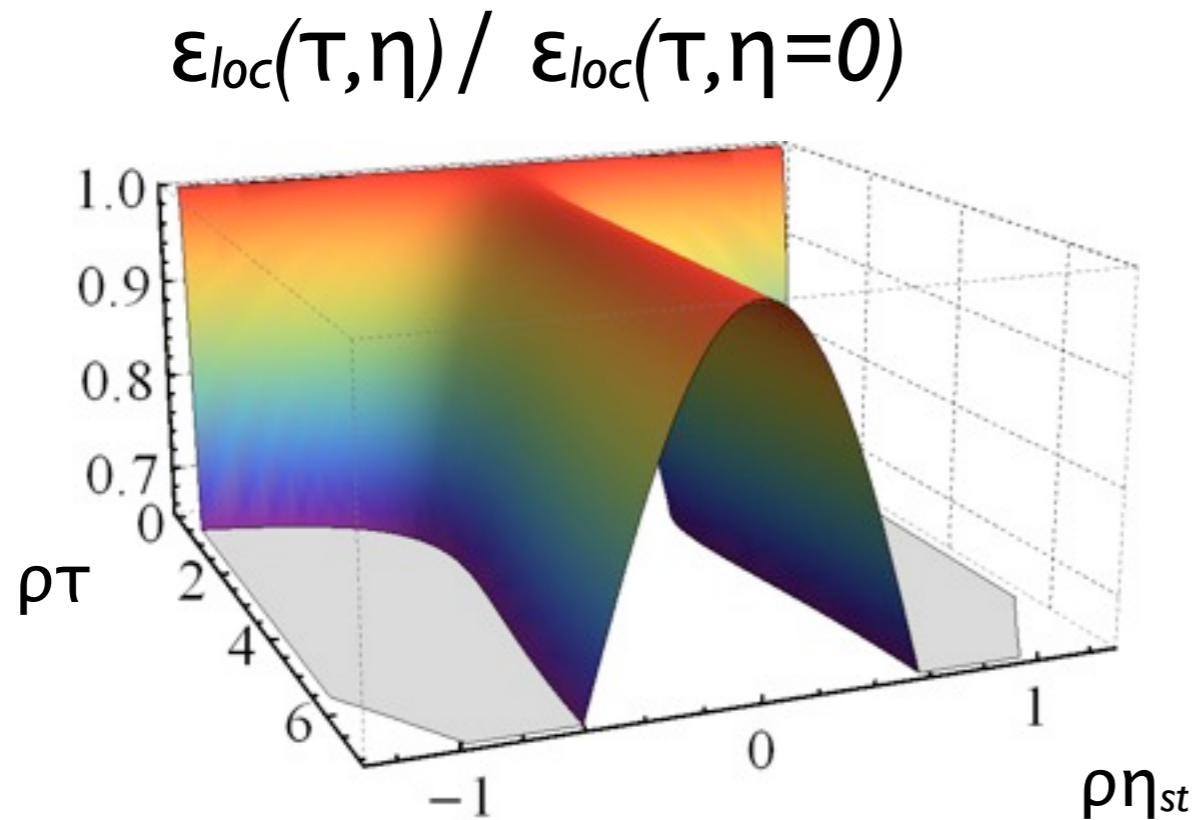
Baryons are not stopped
(Bjorken)

(C-Y Wong 09)



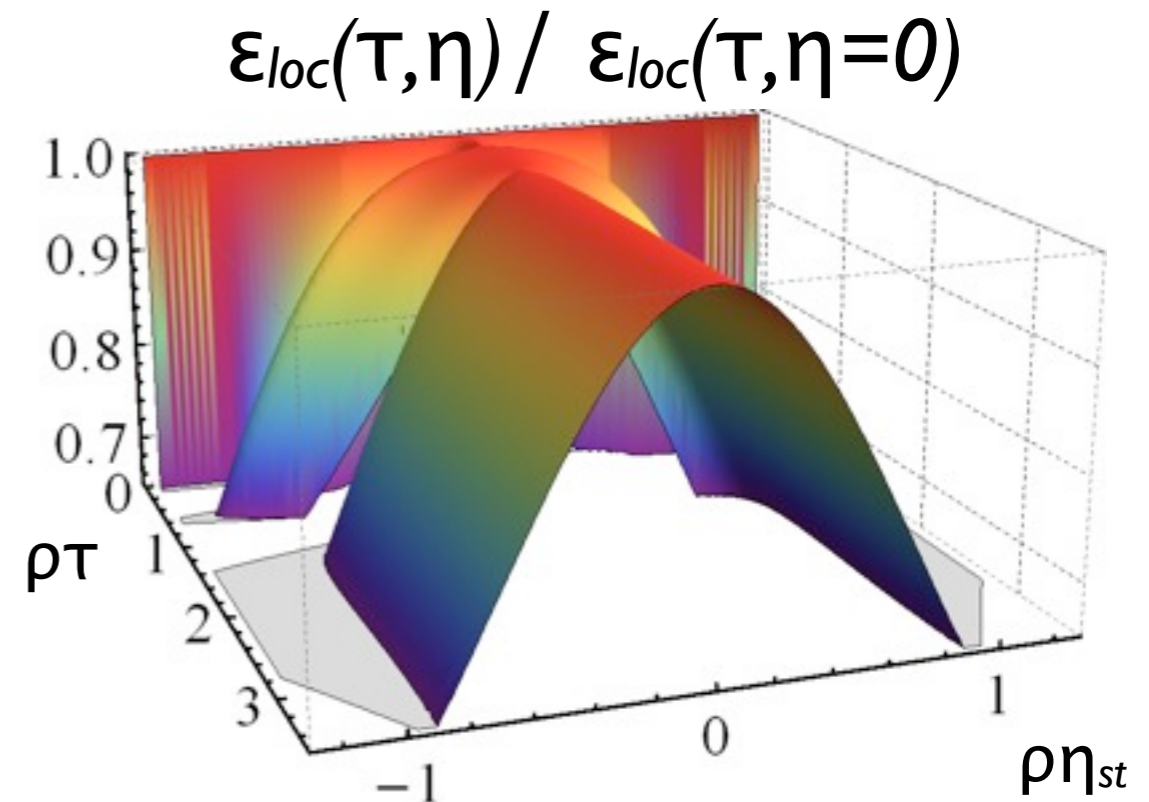
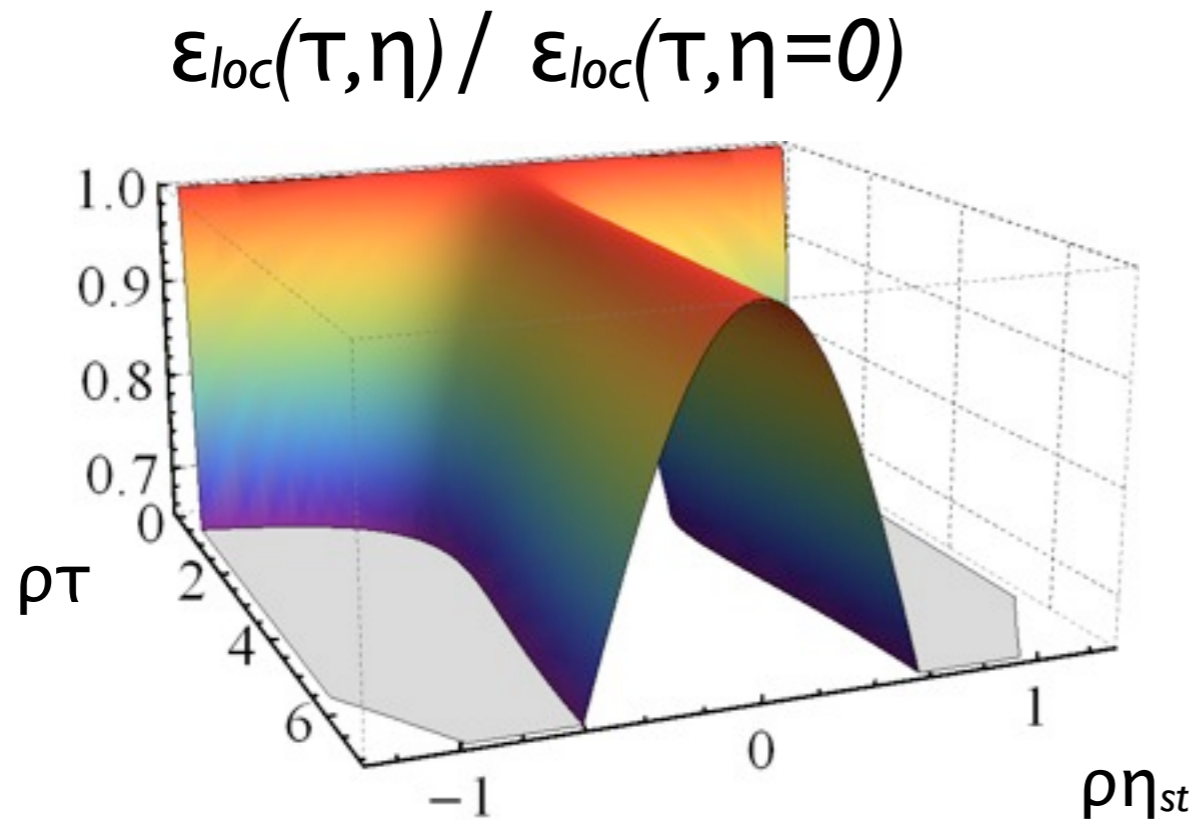
Gaussian profile in rapidity
(Landau)

Non-Boost Invariant Initial Conditions



- Gaussian rapidity profile
 - Low energies: expected from Landau hydrodynamics
 - High energies: relatively mild increase of width

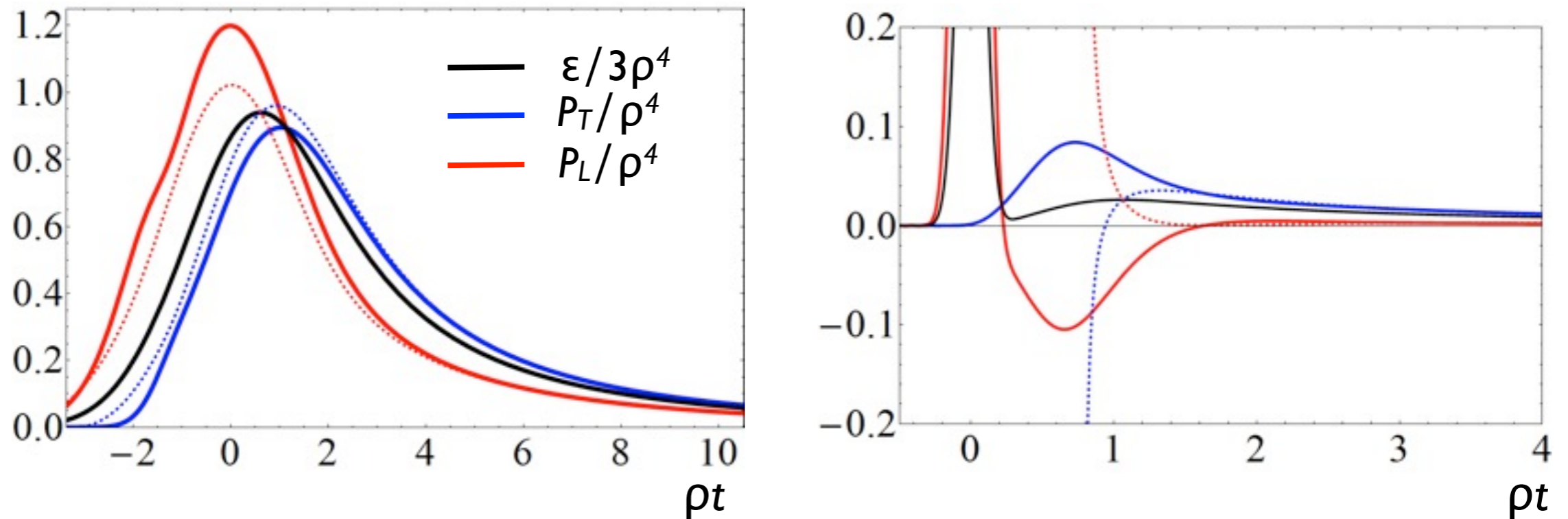
Non-Boost Invariant Initial Conditions



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 - Low energies: expected from Landau hydrodynamics
 - High energies: relatively mild increase of width

(subsequent time evolution well described by Bjorken like flow)

Surprisingly Hydrodynamic

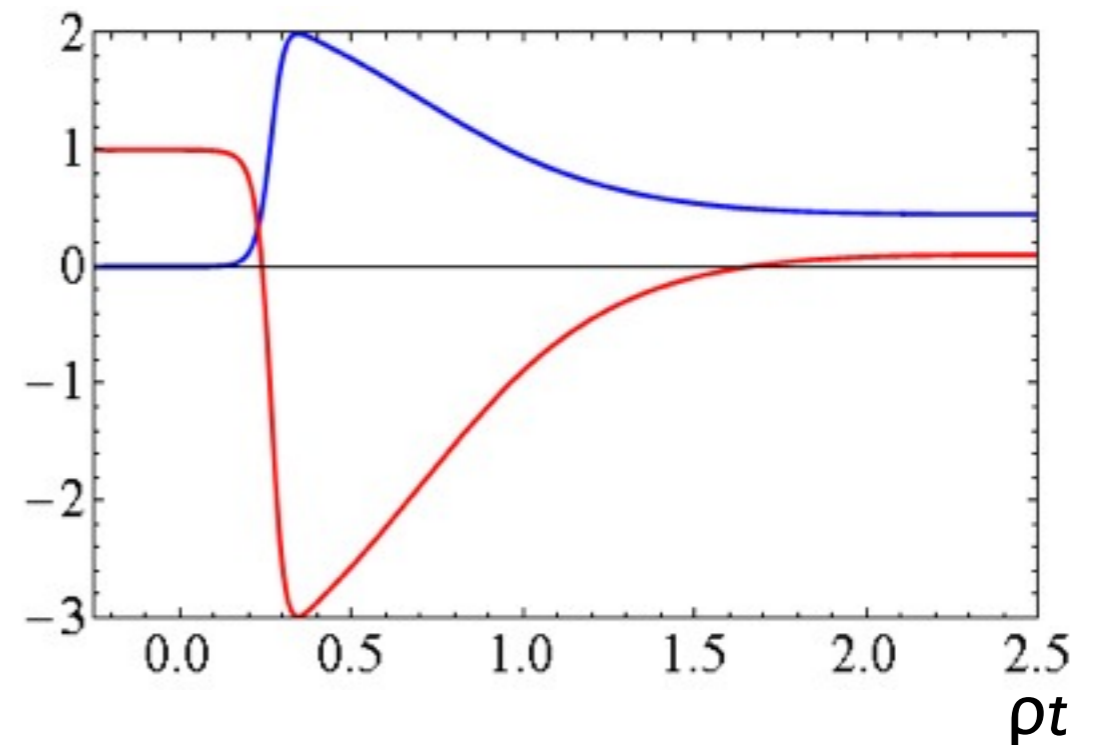
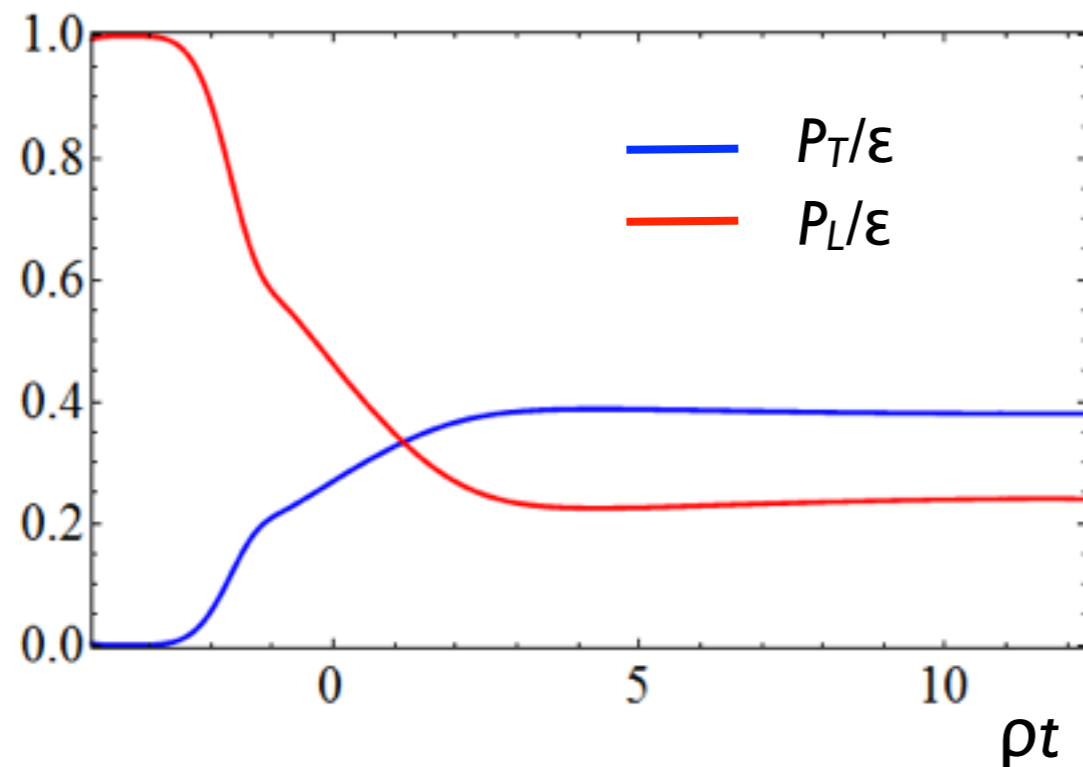


- Good hydrodynamic behavior from very early on
- Energetic shocks: Plasma develops after $t_{\text{hyd}} = 1/\pi T_{\text{hyd}} = 0.87/\mu$
- Very large viscous corrections! *Hydrodynamization*

Chesler & Yaffe, Wu & Romatschke, Heller, Janik & Witaszczyk,
Heller, Mateos, van der Schee, Trancanelli

- Early behavior of pressures due to vanishing initial ϵ

Surprisingly Hydrodynamic

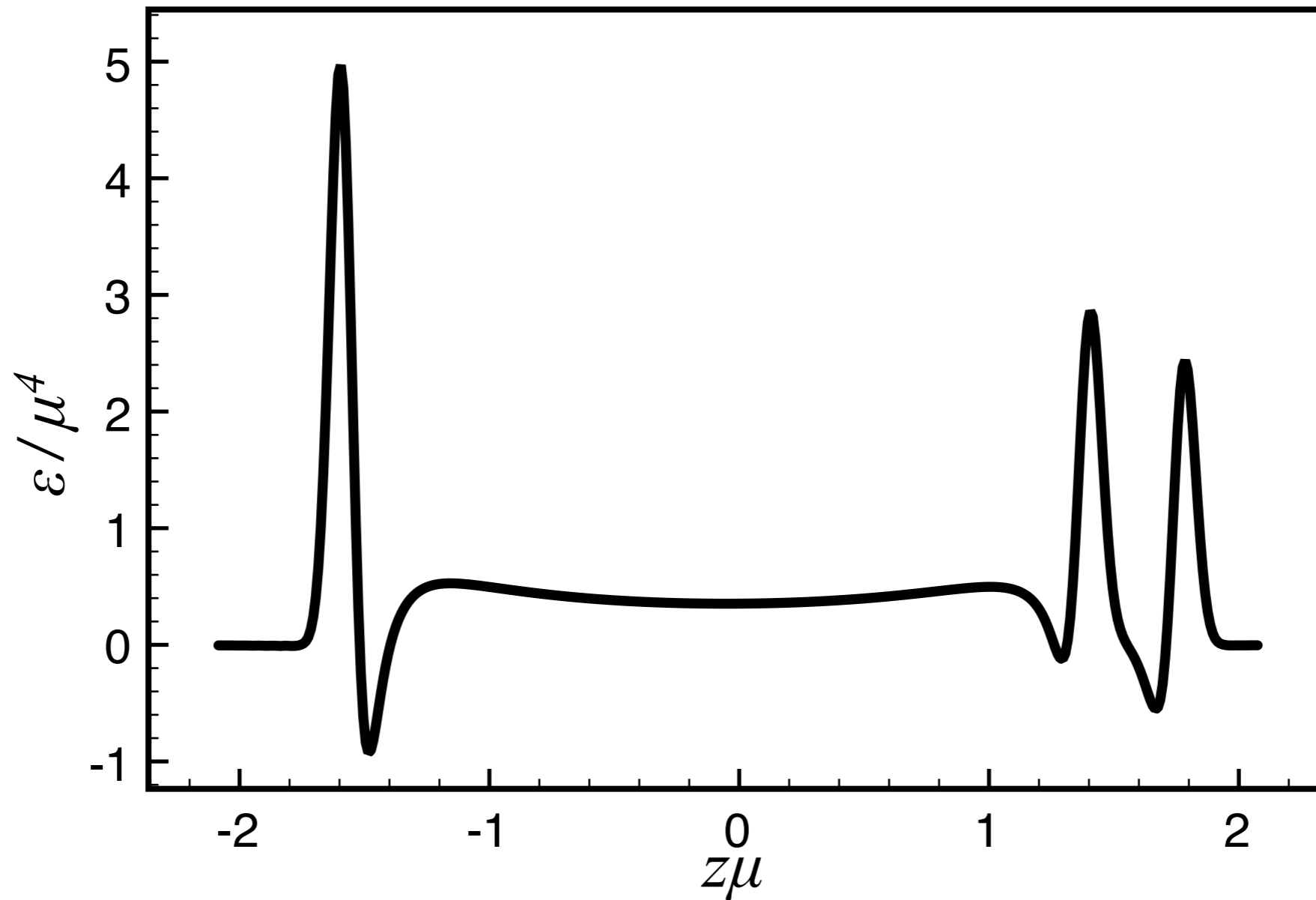


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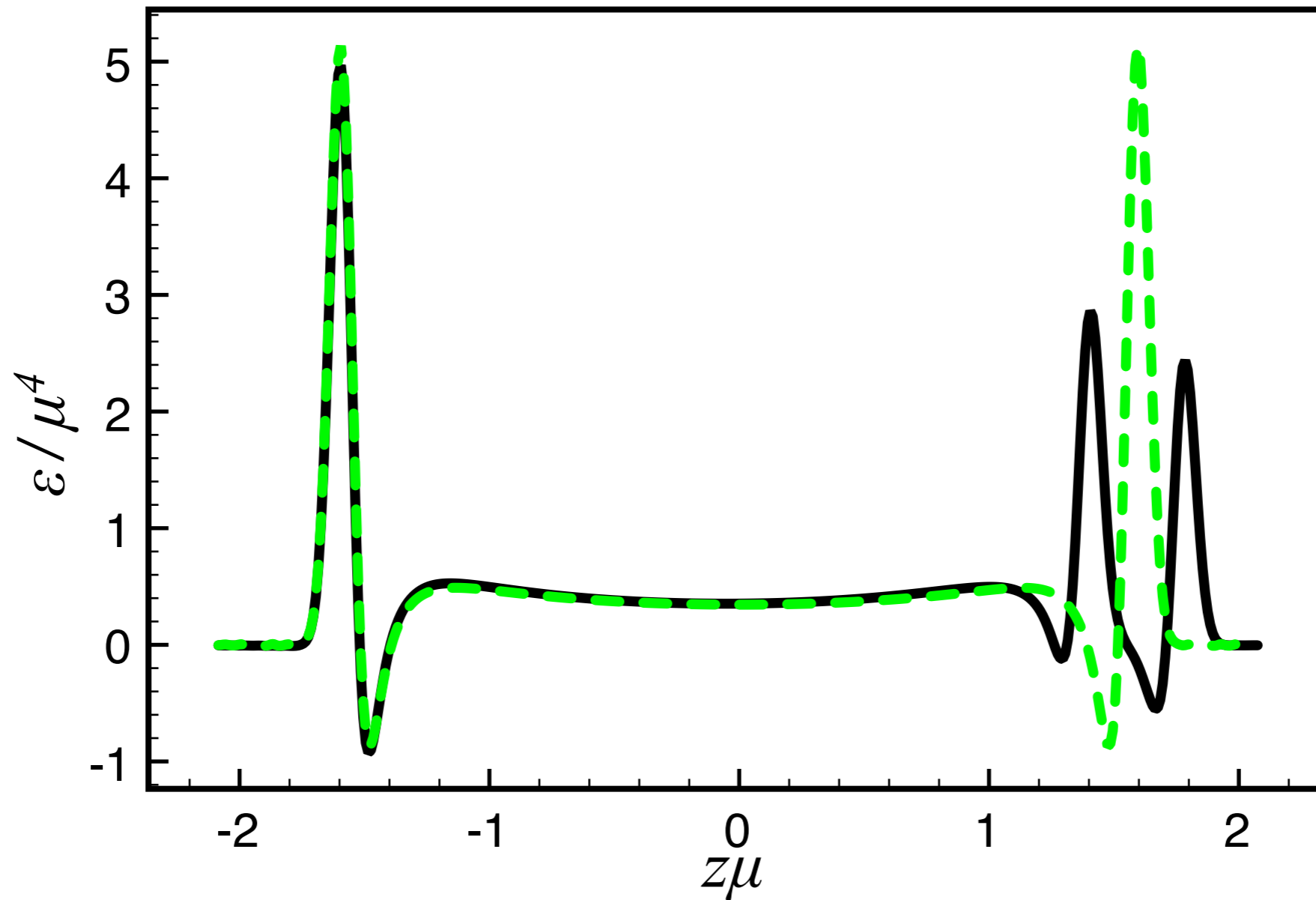
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Towards p - A : Longitudinal Coherence



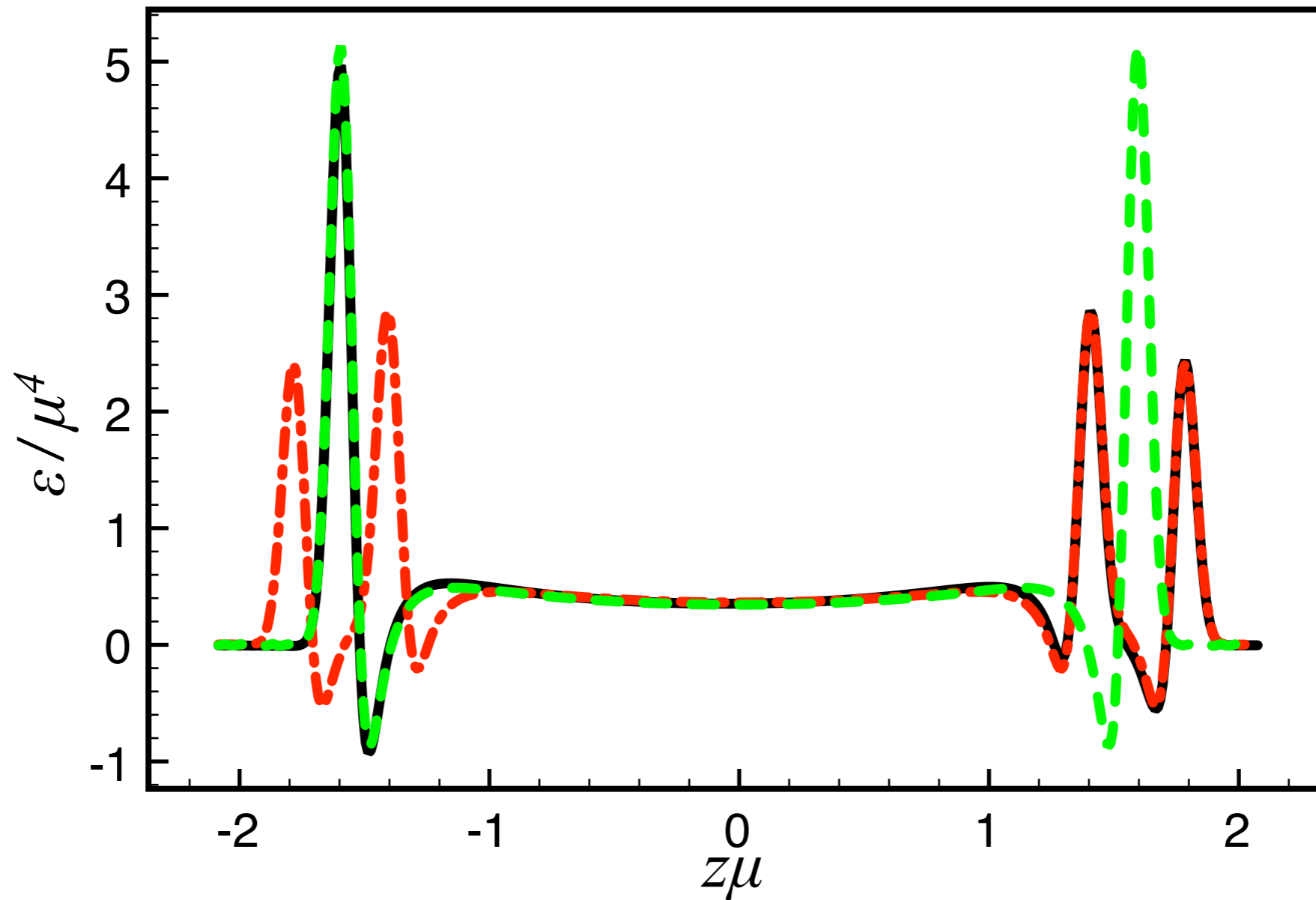
- In the center of mass of the “nucleus-nucleon” collision

Towards p - A : Longitudinal Coherence



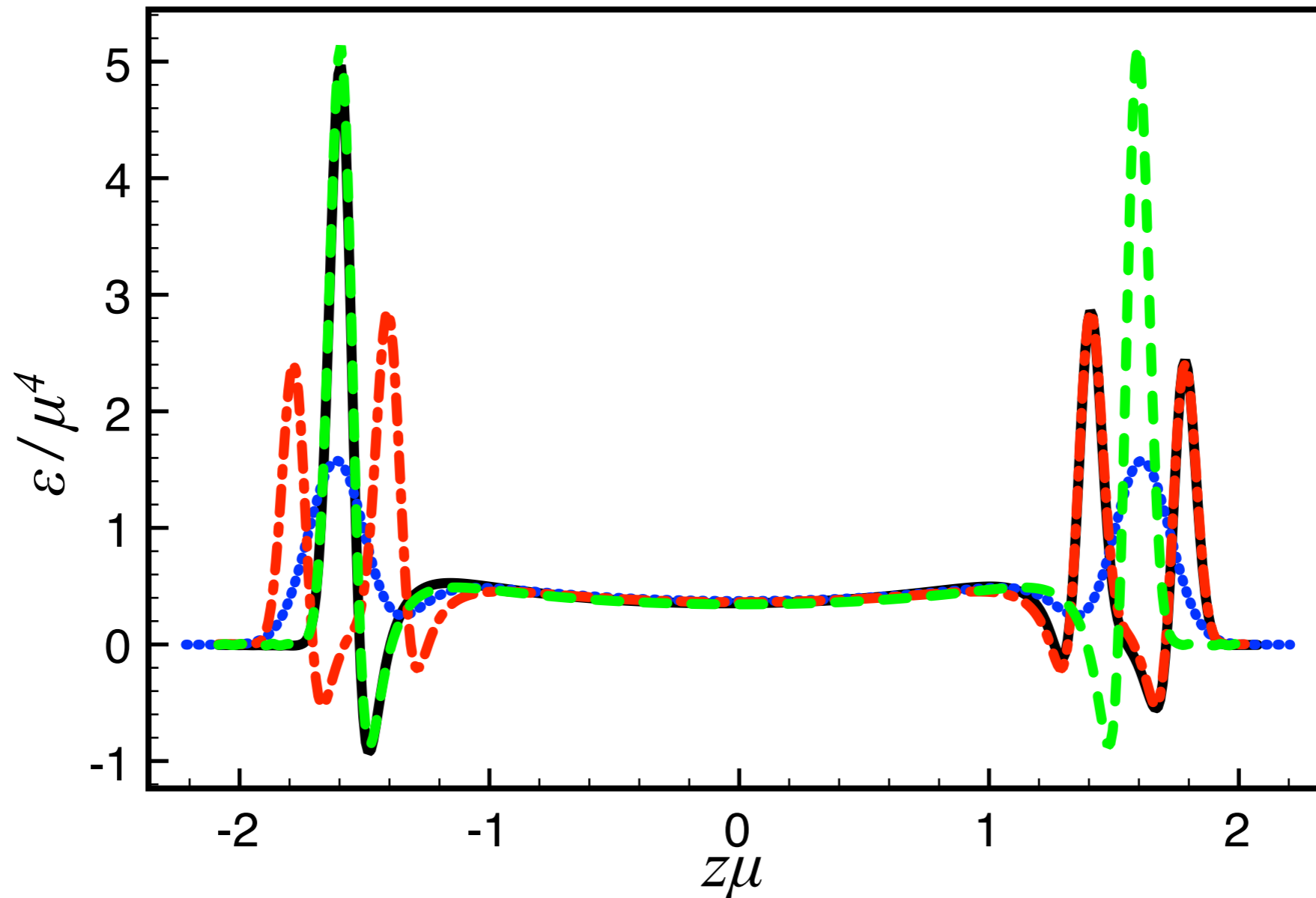
- In the center of mass of the “nucleus-nucleon” collision

Towards p -A: Longitudinal Coherence



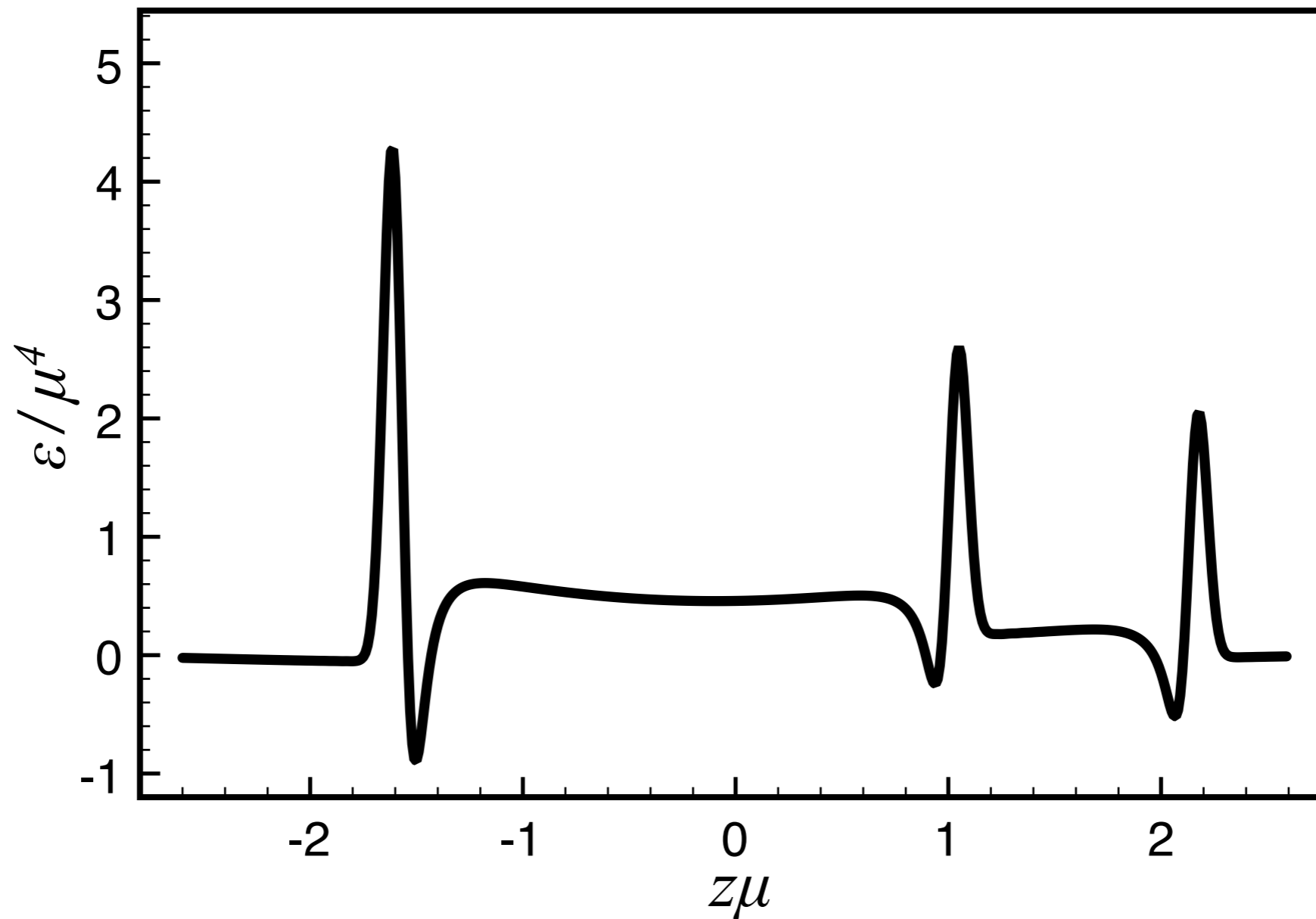
➤ In the center of mass of the “nucleus-nucleon” collision

Towards p -A: Longitudinal Coherence



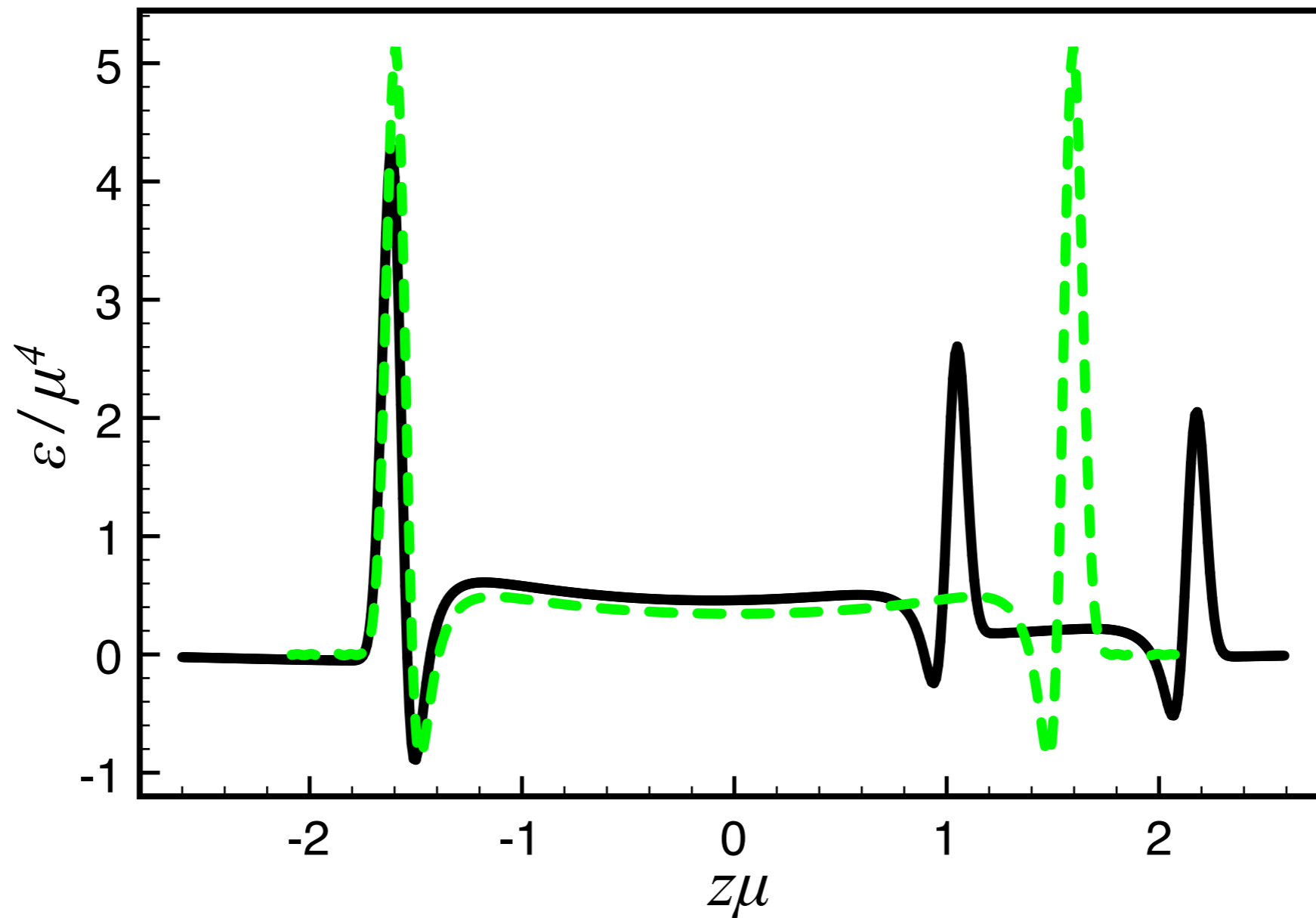
- In the center of mass of the “nucleus-nucleon” collision
- Midd rapidity region independent of collision system
- Maximum at $y=0$ and symmetric w.r.t center of mass

Absence of Longitudinal Coherence



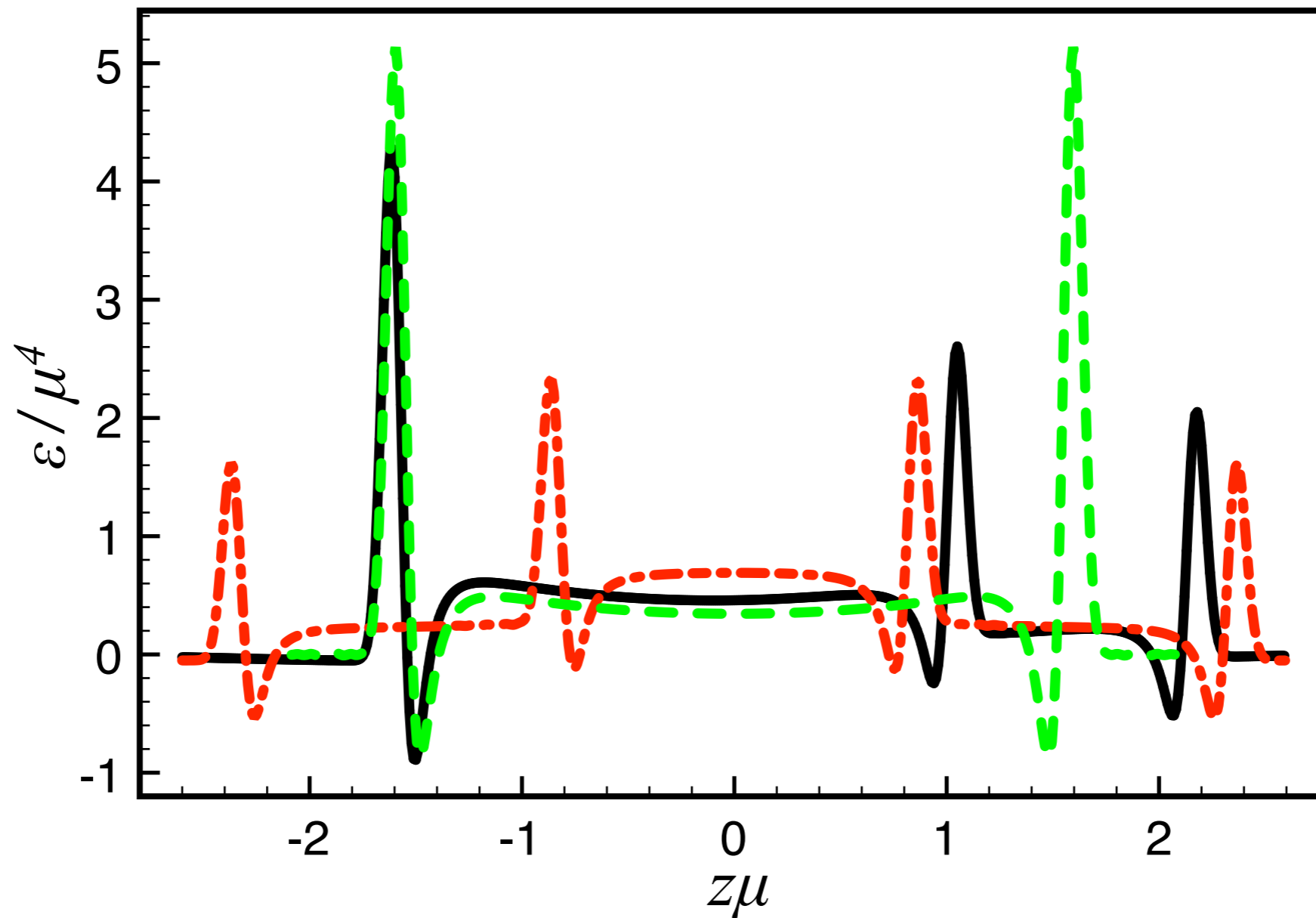
➤ Larger longitudinal structures

Absence of Longitudinal Coherence



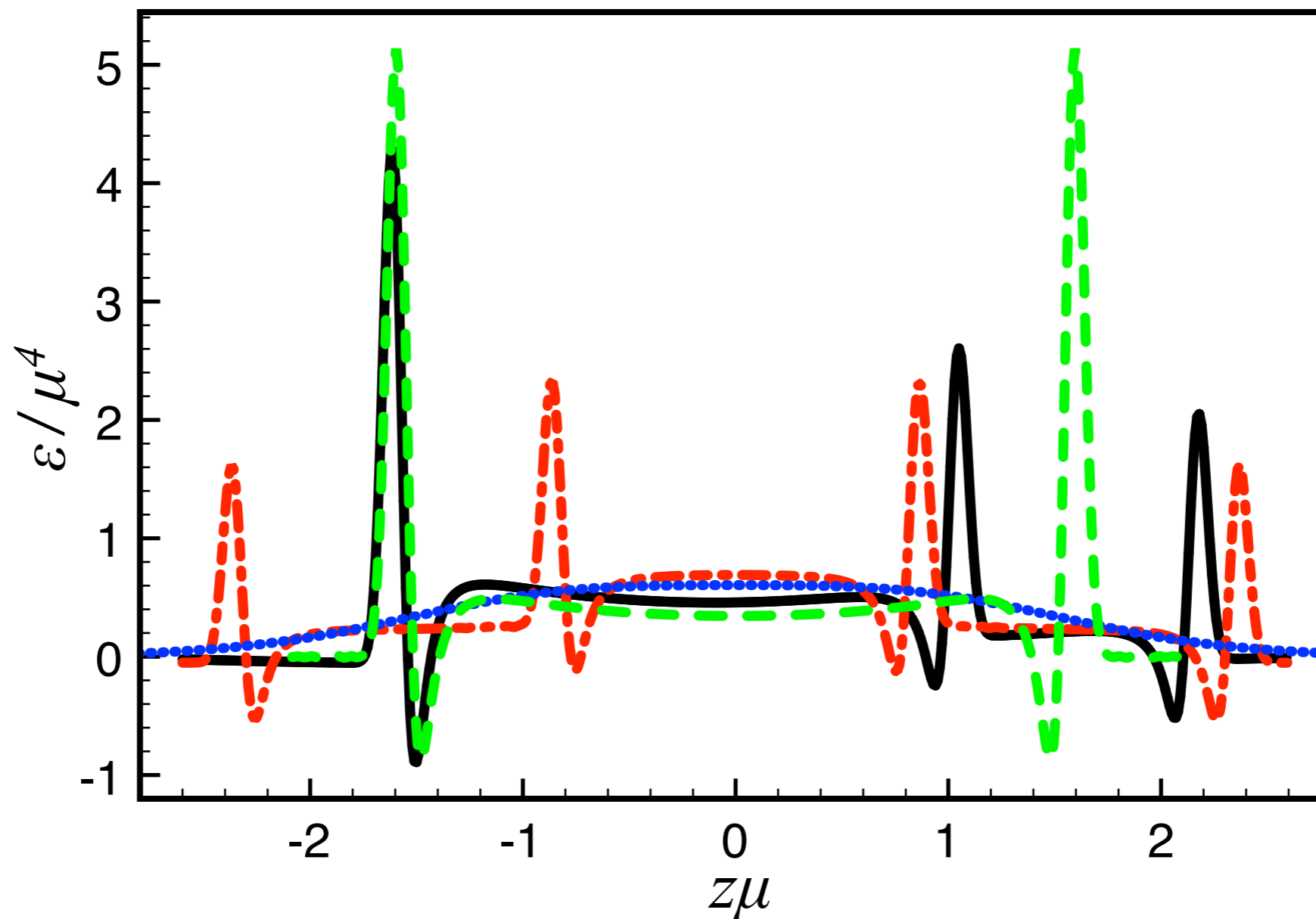
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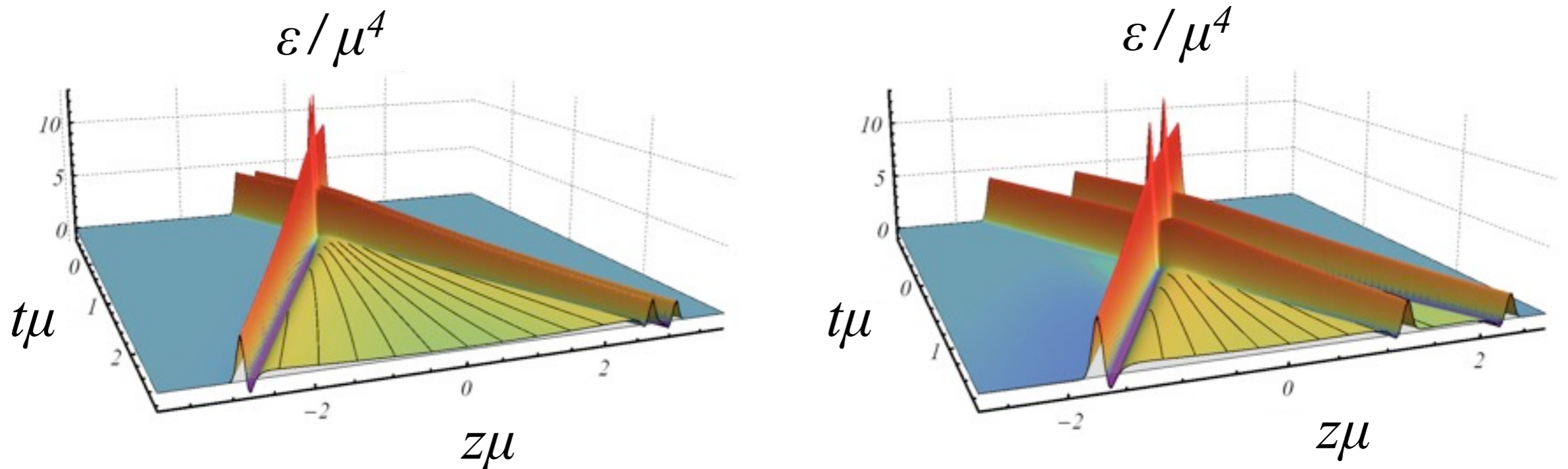
➤ Larger longitudinal structures

Absence of Longitudinal Coherence



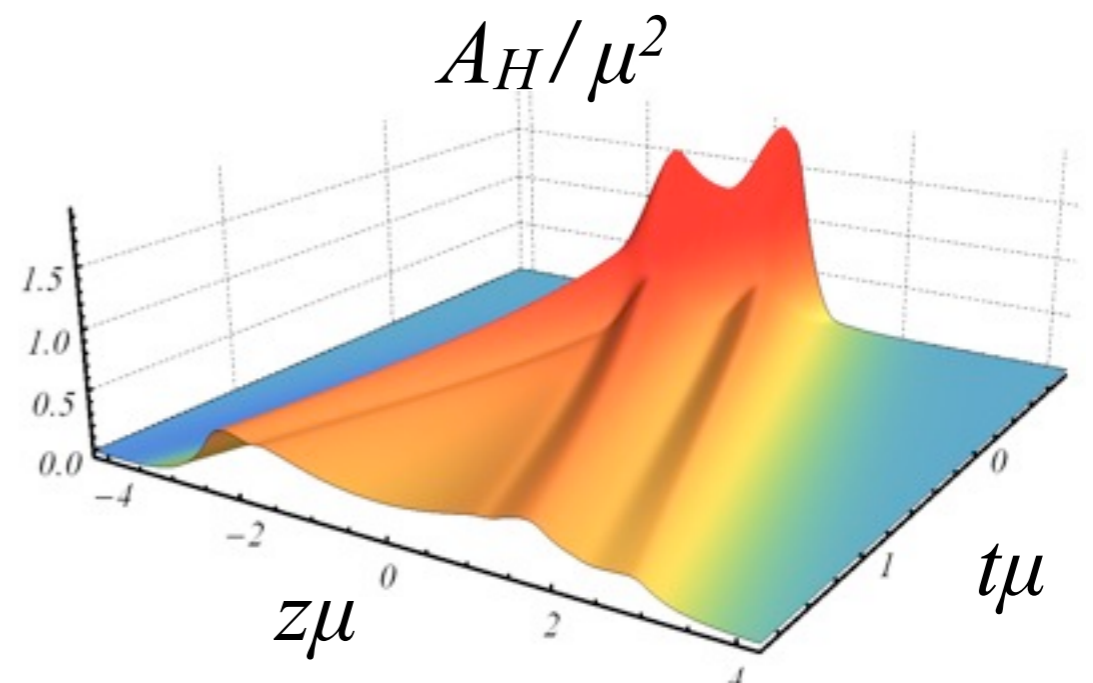
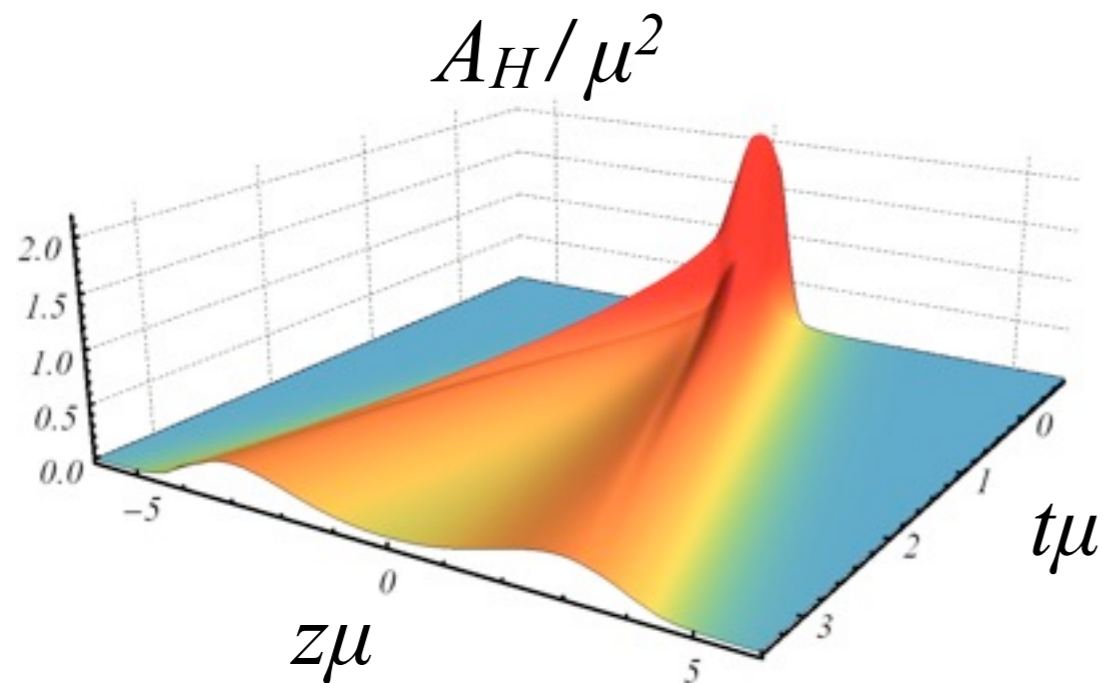
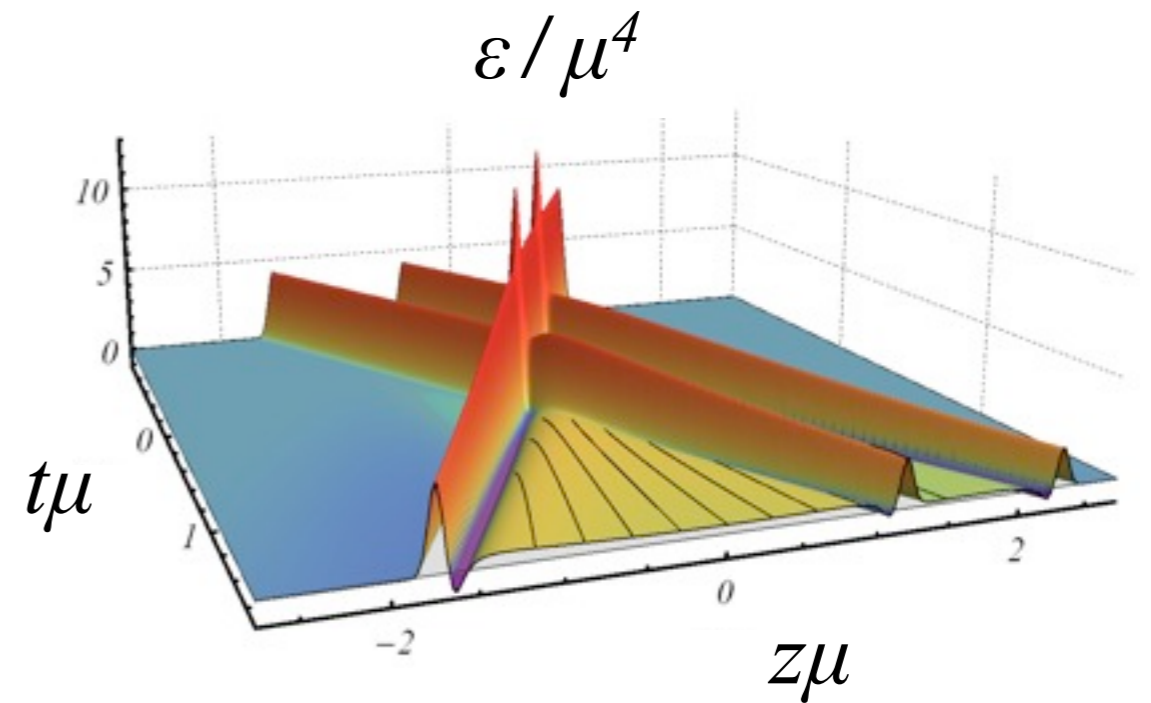
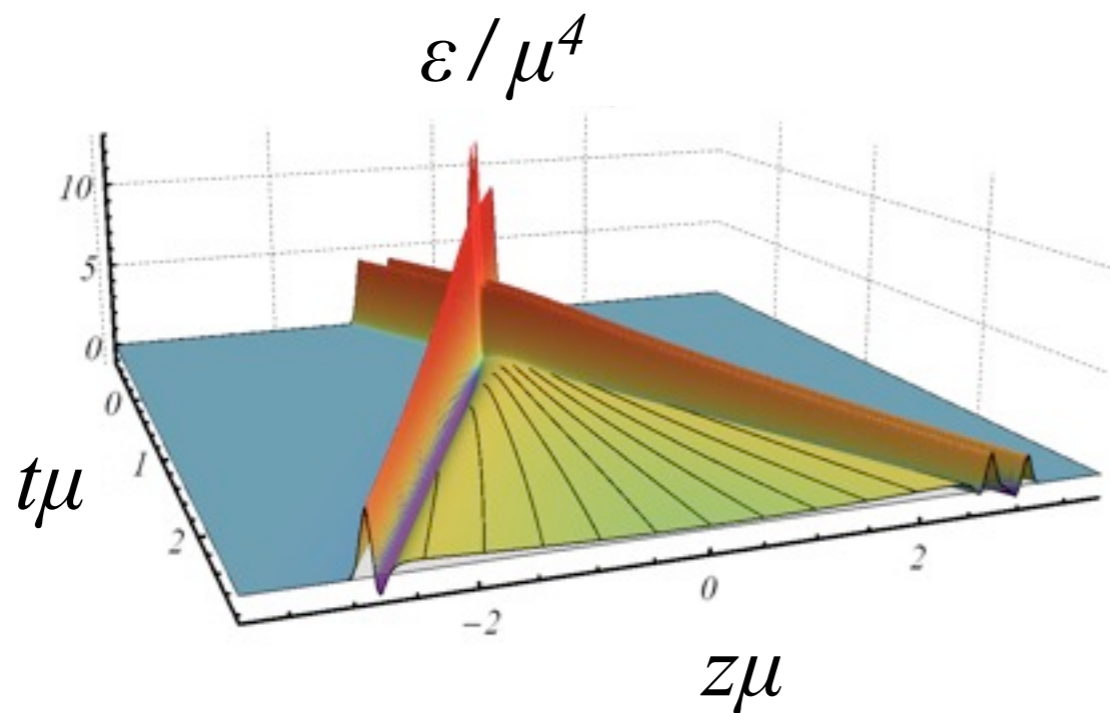
- Larger longitudinal structures
- Sensitivity to the colliding system
- Asymmetric distribution for asymmetric systems

Holography and Physics at T-Scale



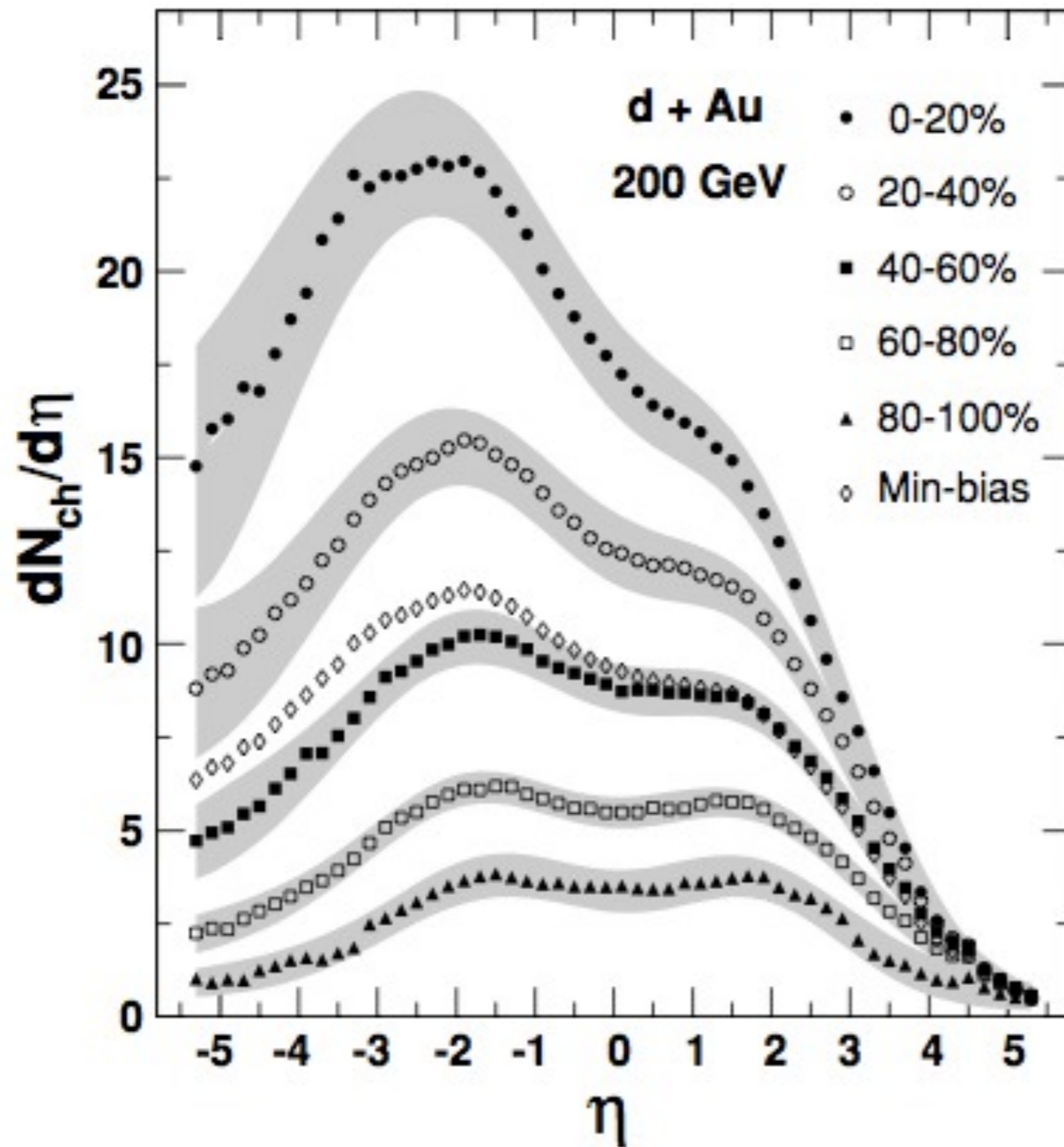
- Coherent response depends on longitudinal structure
 - Structures of size $< 1/\pi T_{\text{hyd}}$ are not resolved by the collision dynamics
 - Structures of size $< 1/\pi T_{\text{hyd}}$ act incoherently

Holography and Physics at T-Scale



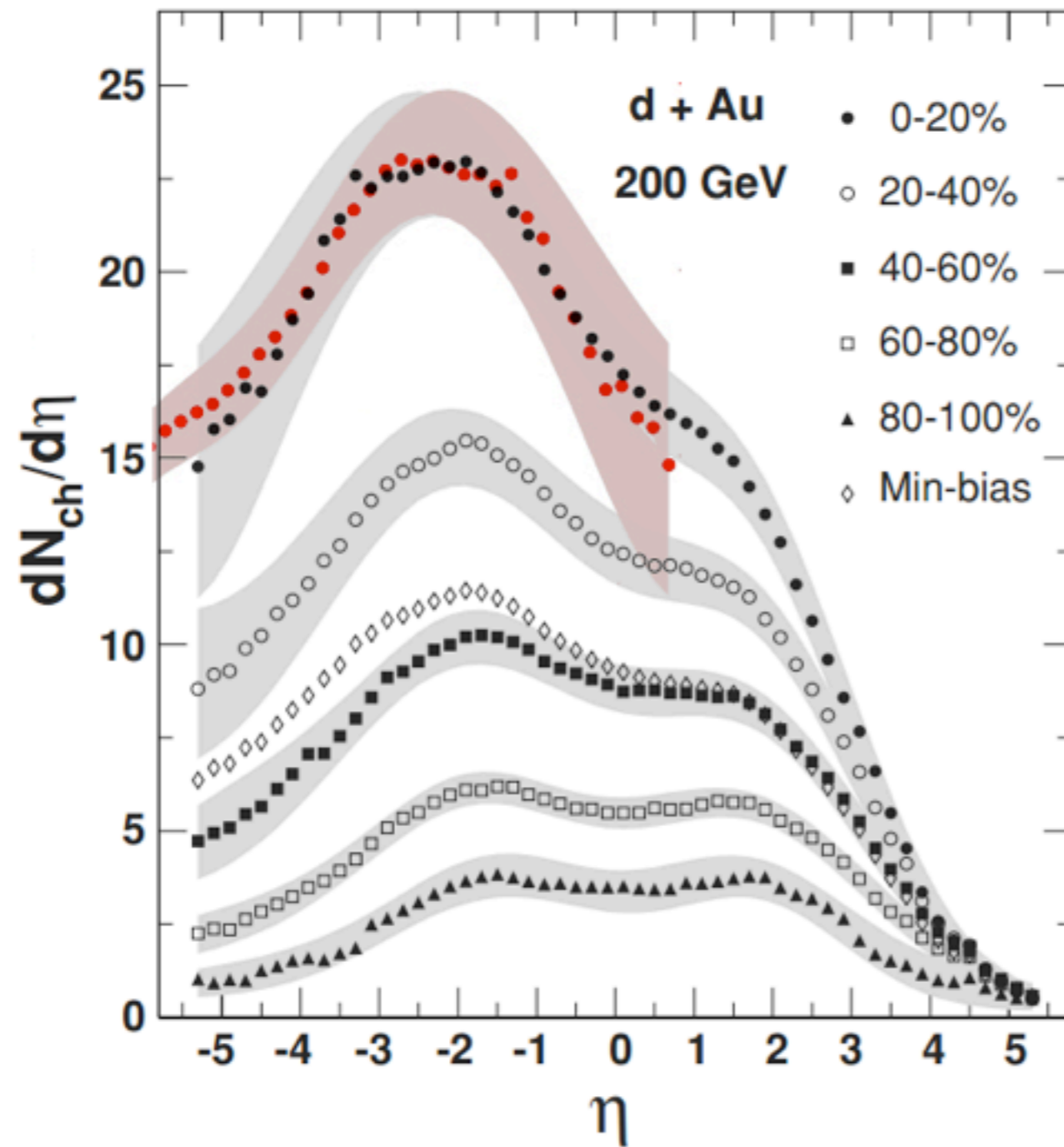
➤ Holography provides a simple picture!

Qualitative Expectation



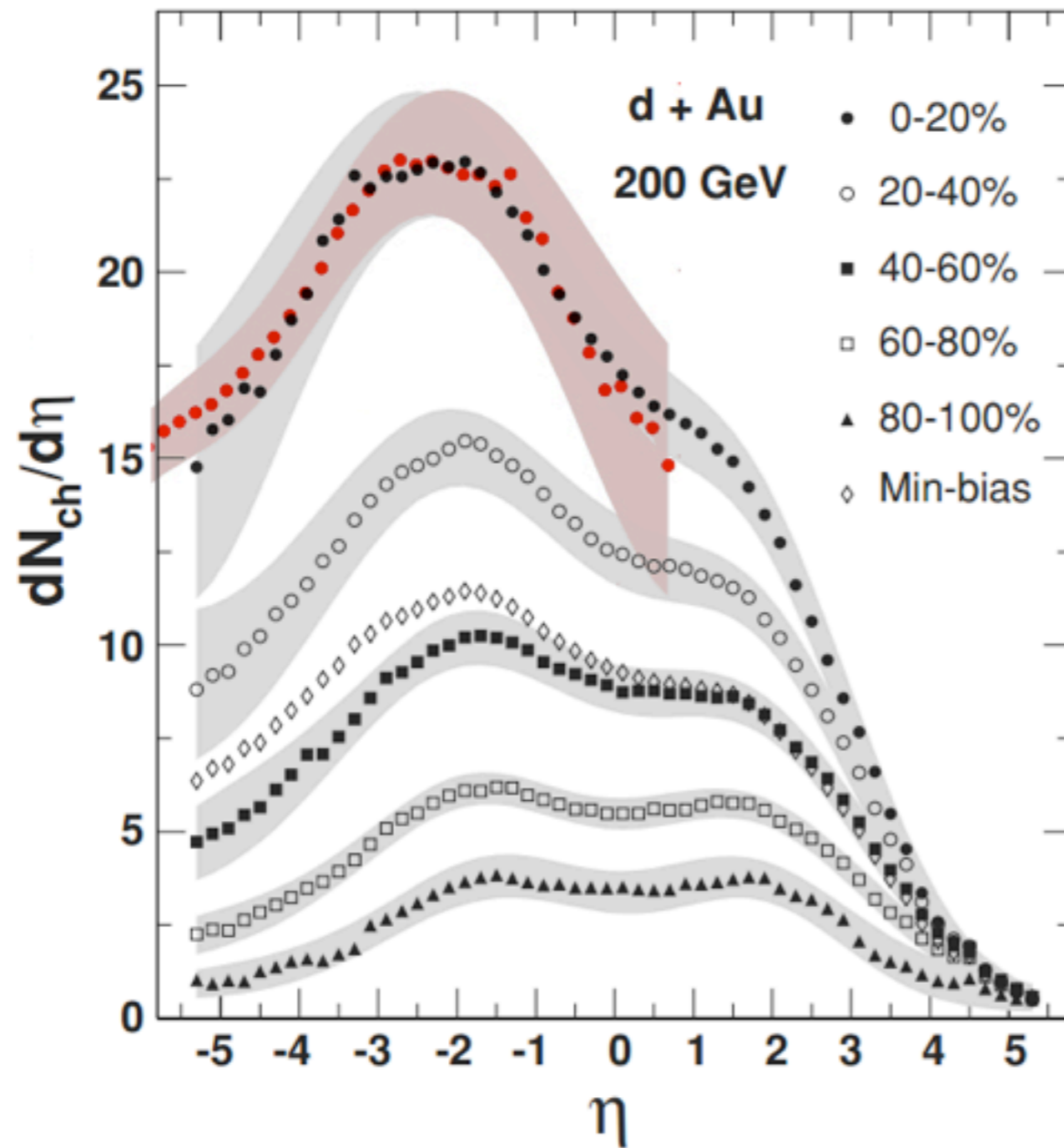
(Steinberg 07)

Qualitative Expectation



(Steinberg 07)

Qualitative Expectation



➤ Unfair comparison:
pseudo-rapidity

$$\eta = \frac{1}{2} \ln \left(\frac{|\mathbf{P}| + p_L}{|\mathbf{P}| - p_L} \right)$$

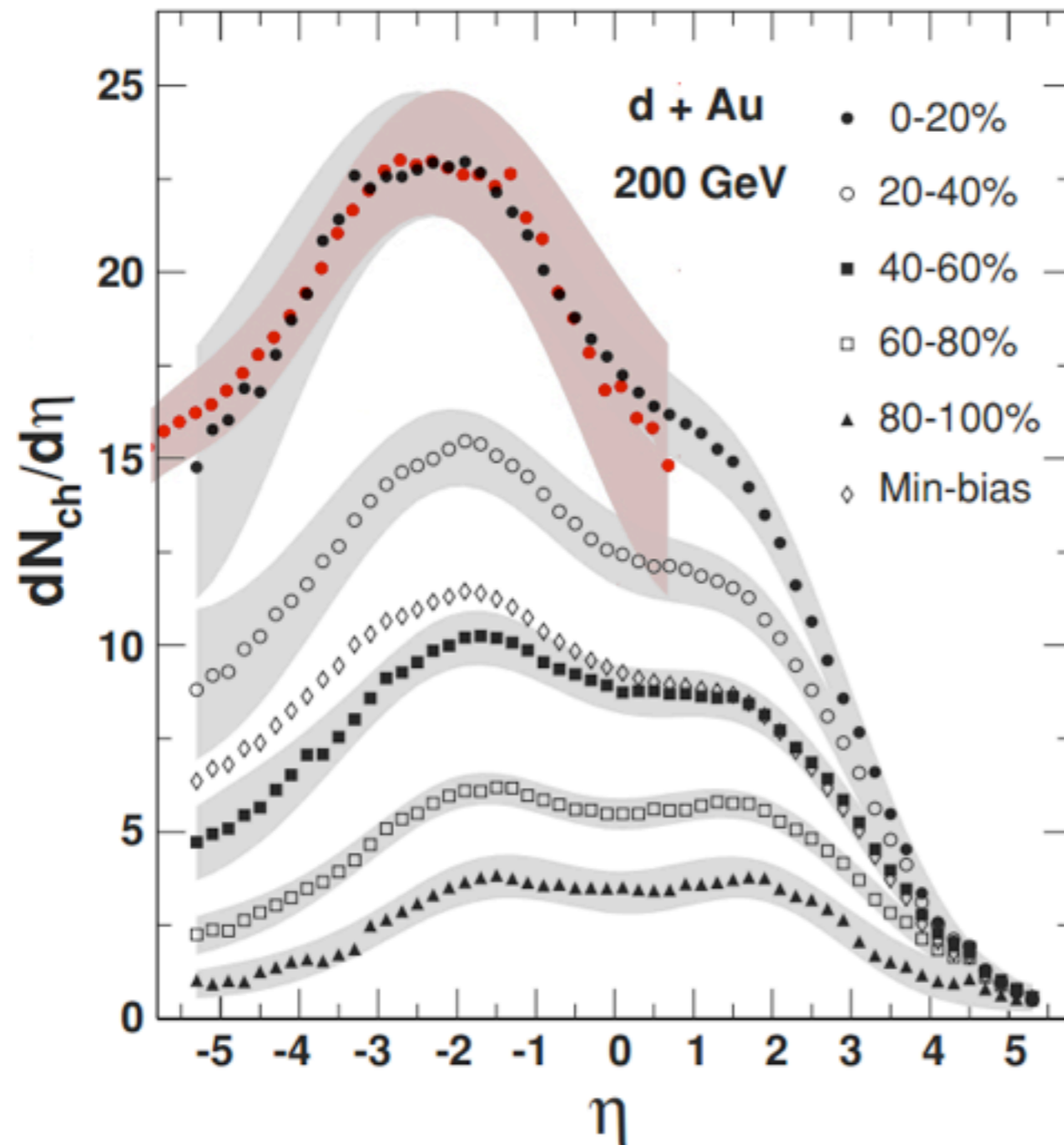
rapidity

$$y = \frac{1}{2} \ln \left(\frac{E + p_L}{E - p_L} \right)$$

most non-trivial structure
due to the transverse mass

(Steinberg 07)

Qualitative Expectation



► Unfair comparison:
pseudo-rapidity

$$\eta = \frac{1}{2} \ln \left(\frac{|\mathbf{P}| + p_L}{|\mathbf{P}| - p_L} \right)$$

rapidity

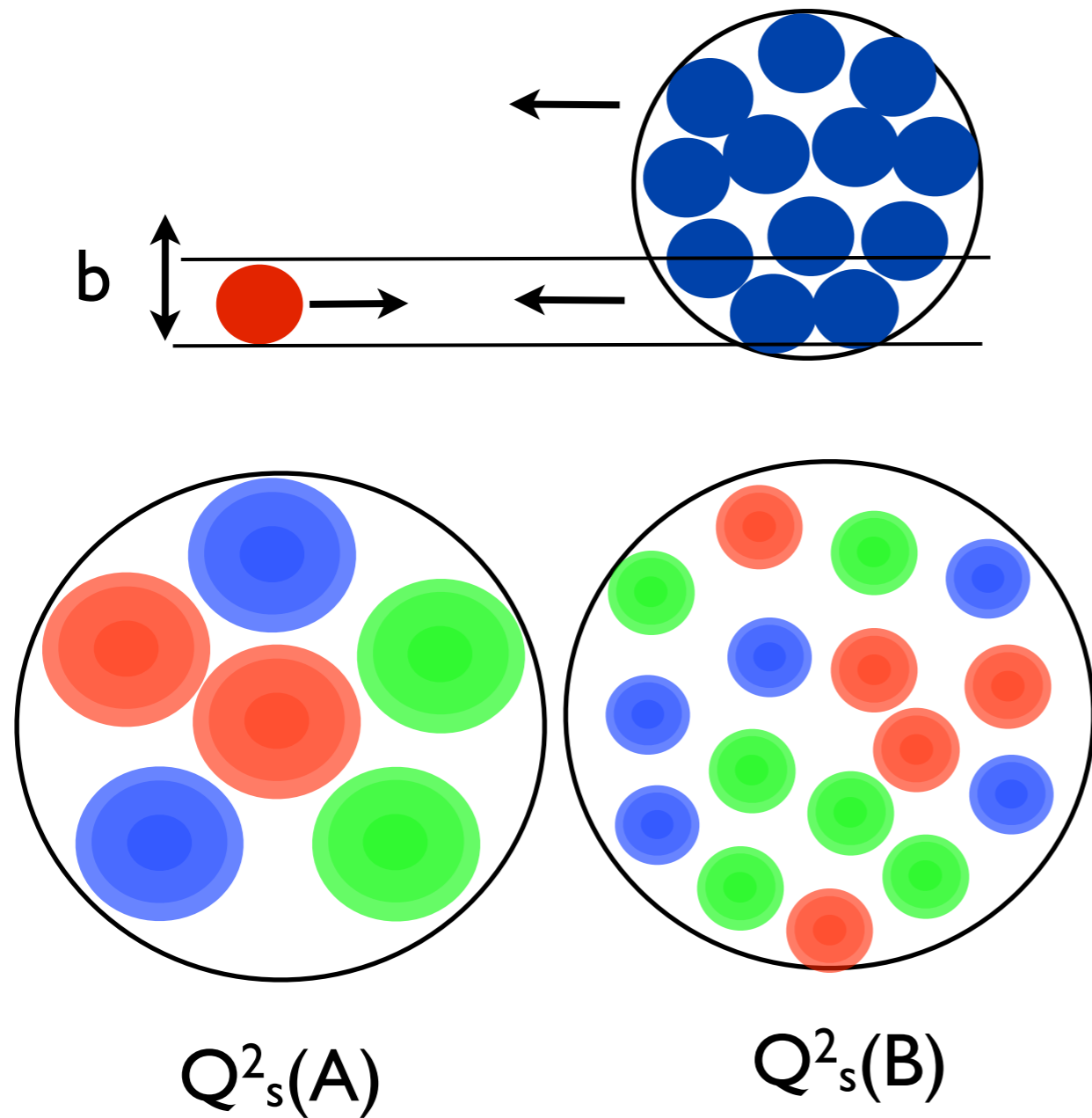
$$y = \frac{1}{2} \ln \left(\frac{E + p_L}{E - p_L} \right)$$

most non-trivial structure
due to the transverse mass

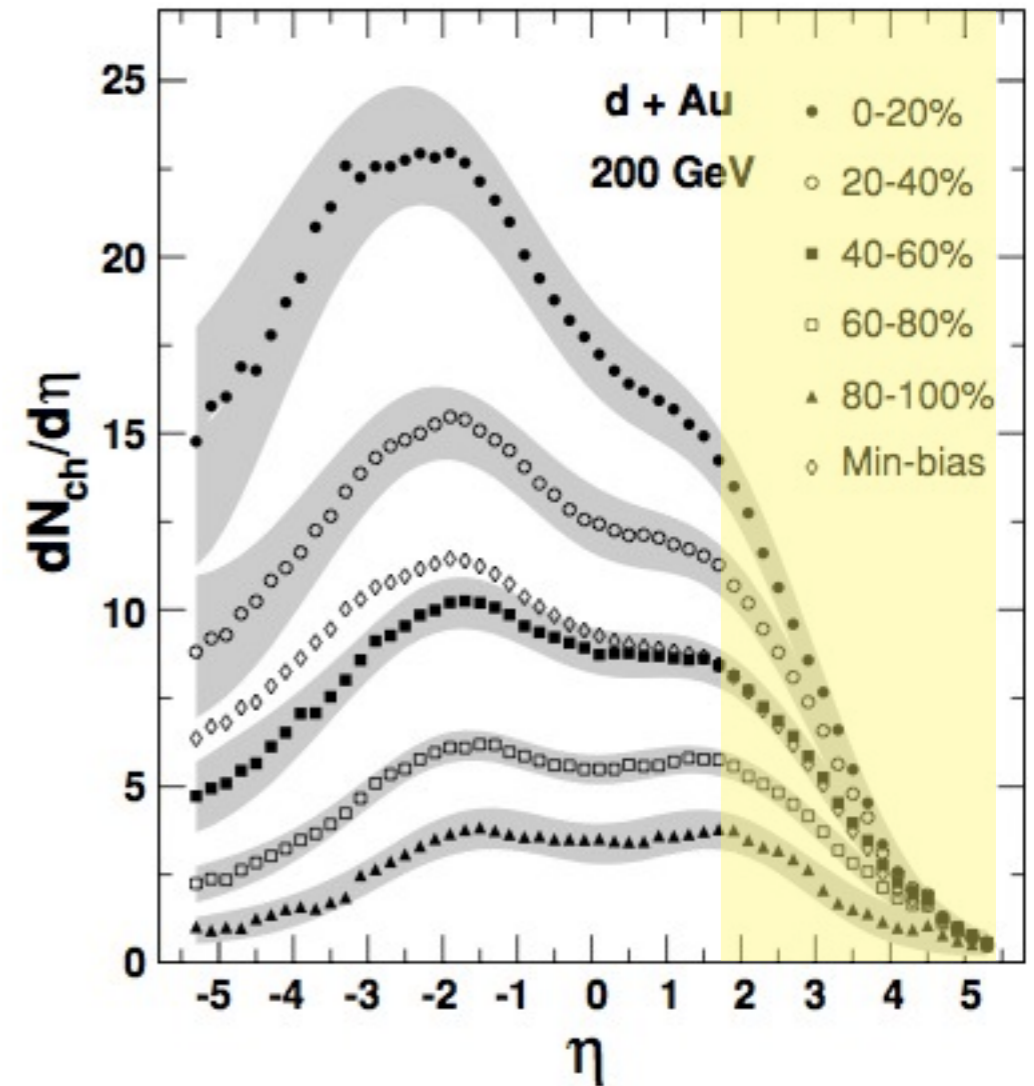
Careful analysis of LE data: consistent with a rapidity shift

$$y_s = y_{\text{c.o.m}} + 0.3 \quad (\text{Steinberg 07})$$

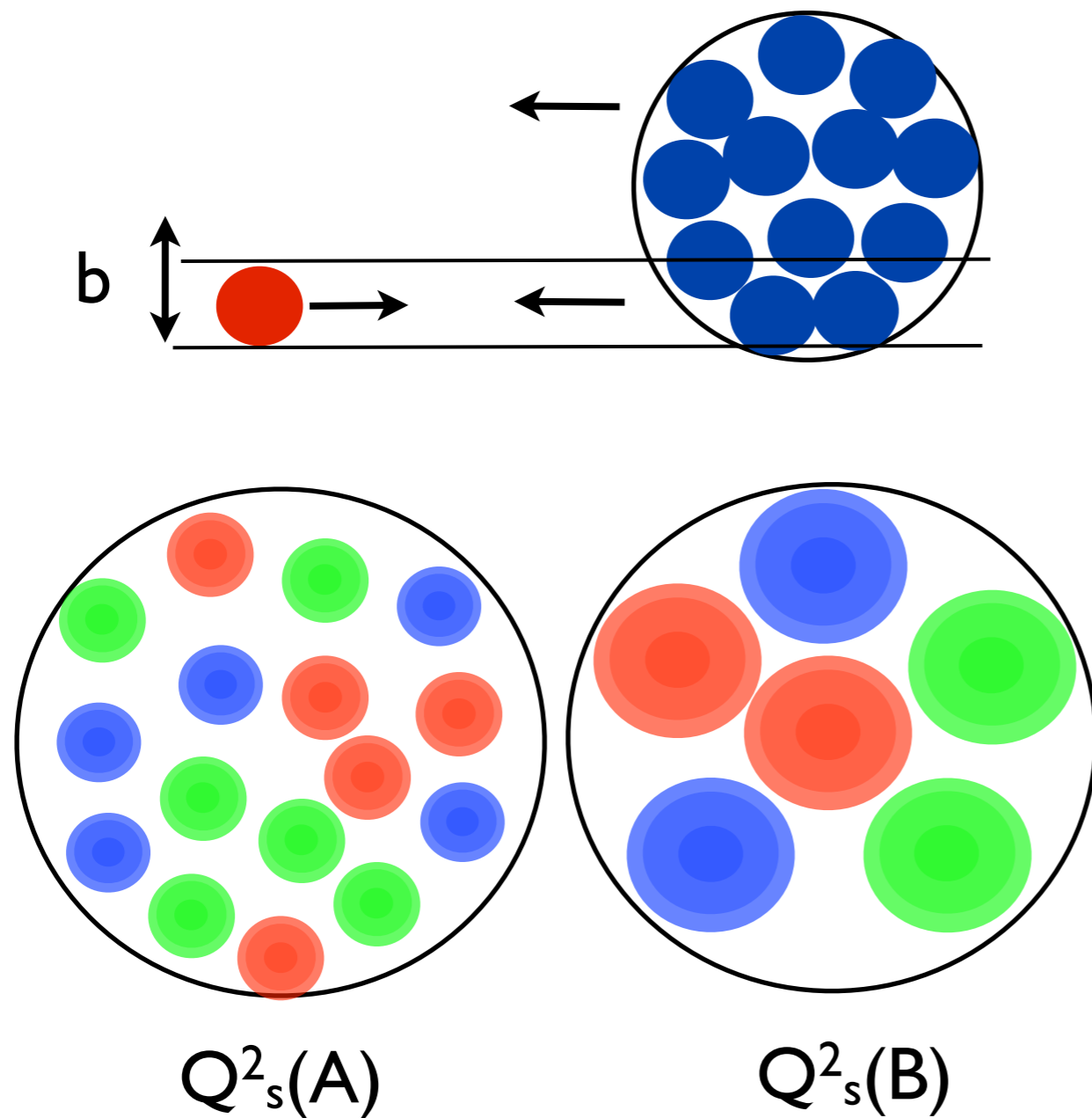
The Weak Coupling Picture



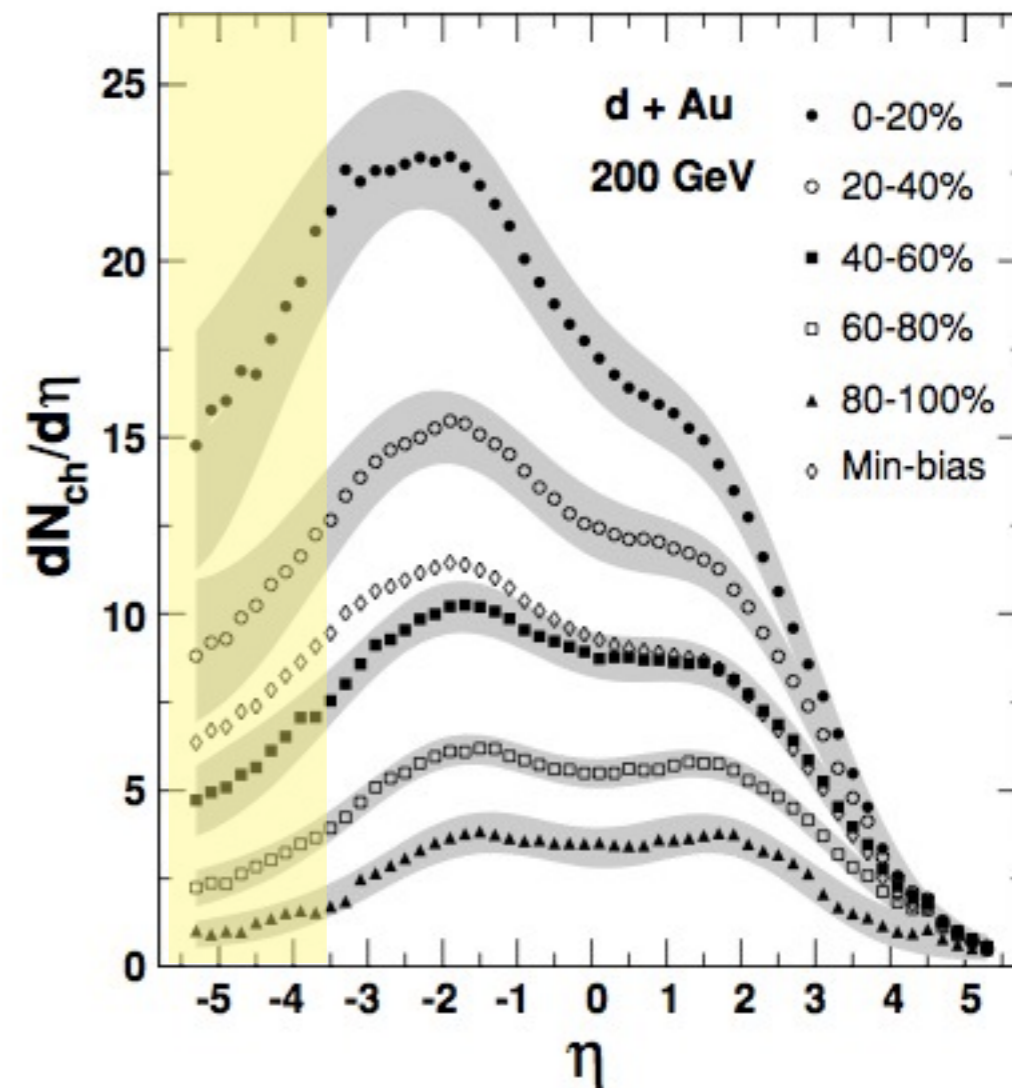
$$\frac{dN}{dy} \propto \pi R^2 Q_s^2(A)$$



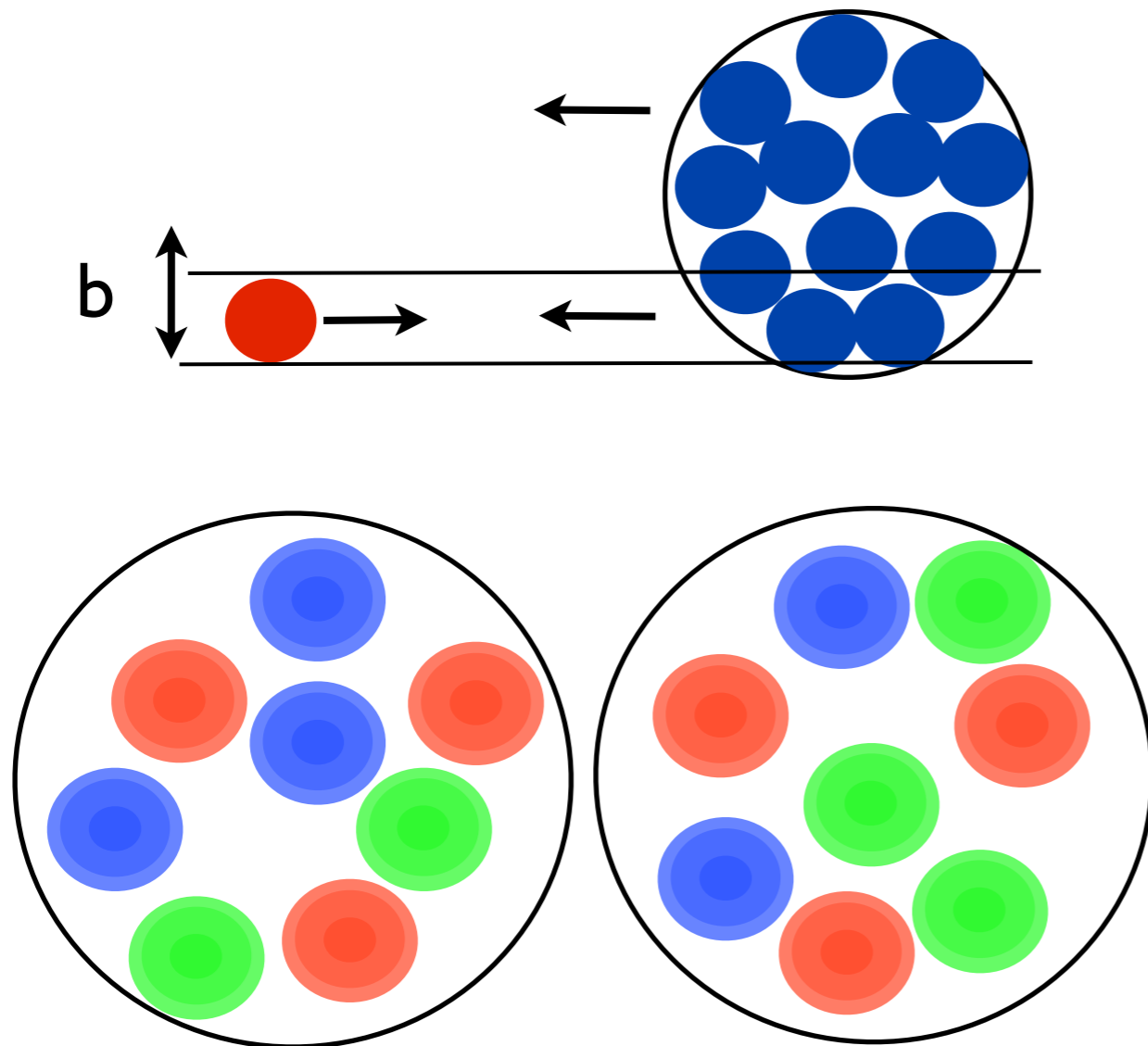
The Weak Coupling Picture



$$\frac{dN}{dy} \propto \pi R^2 Q^2_s(B)$$



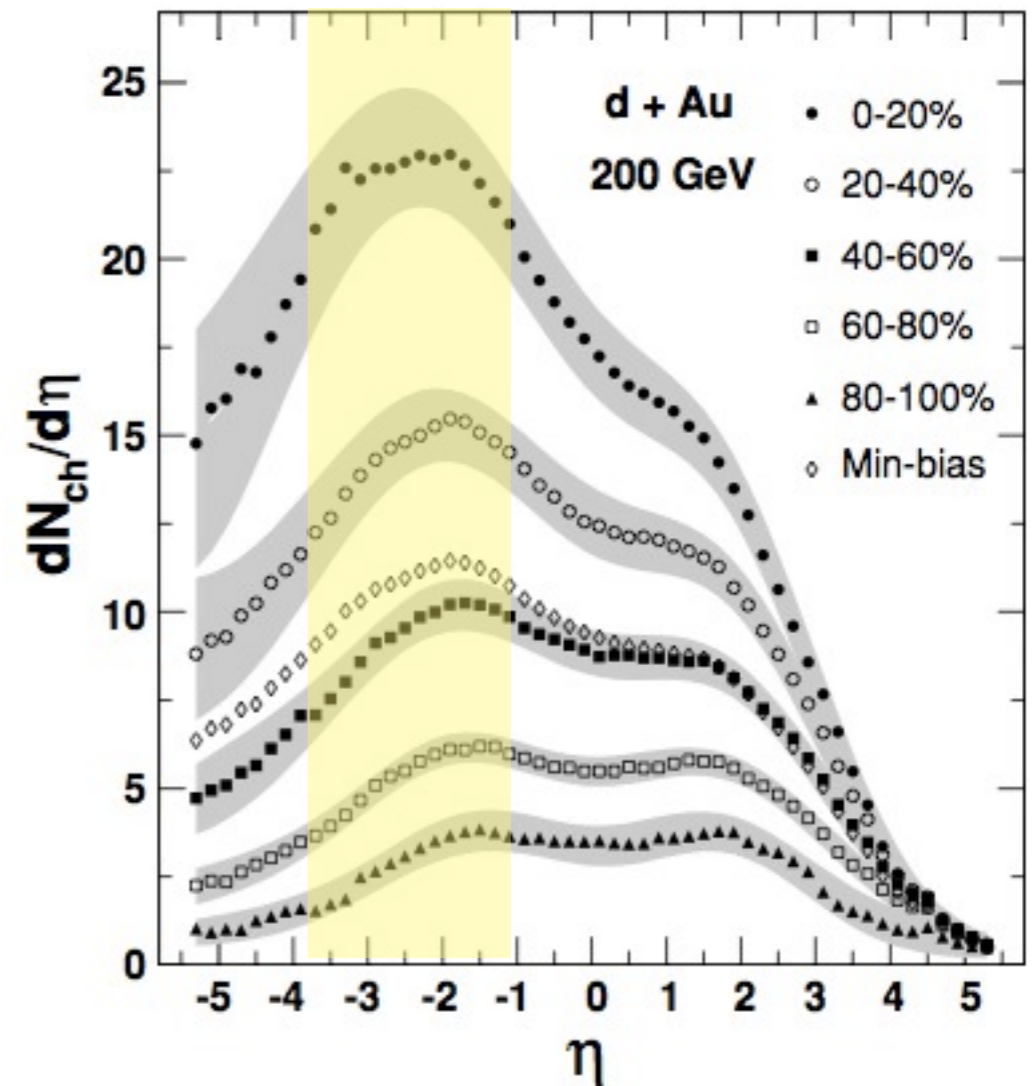
The Weak Coupling Picture



$$Q_s^2(A) = Q_s^2(B)$$

$$Q_s^2(A) = Q_0^2 A e^{\lambda(Y_a - y)}$$

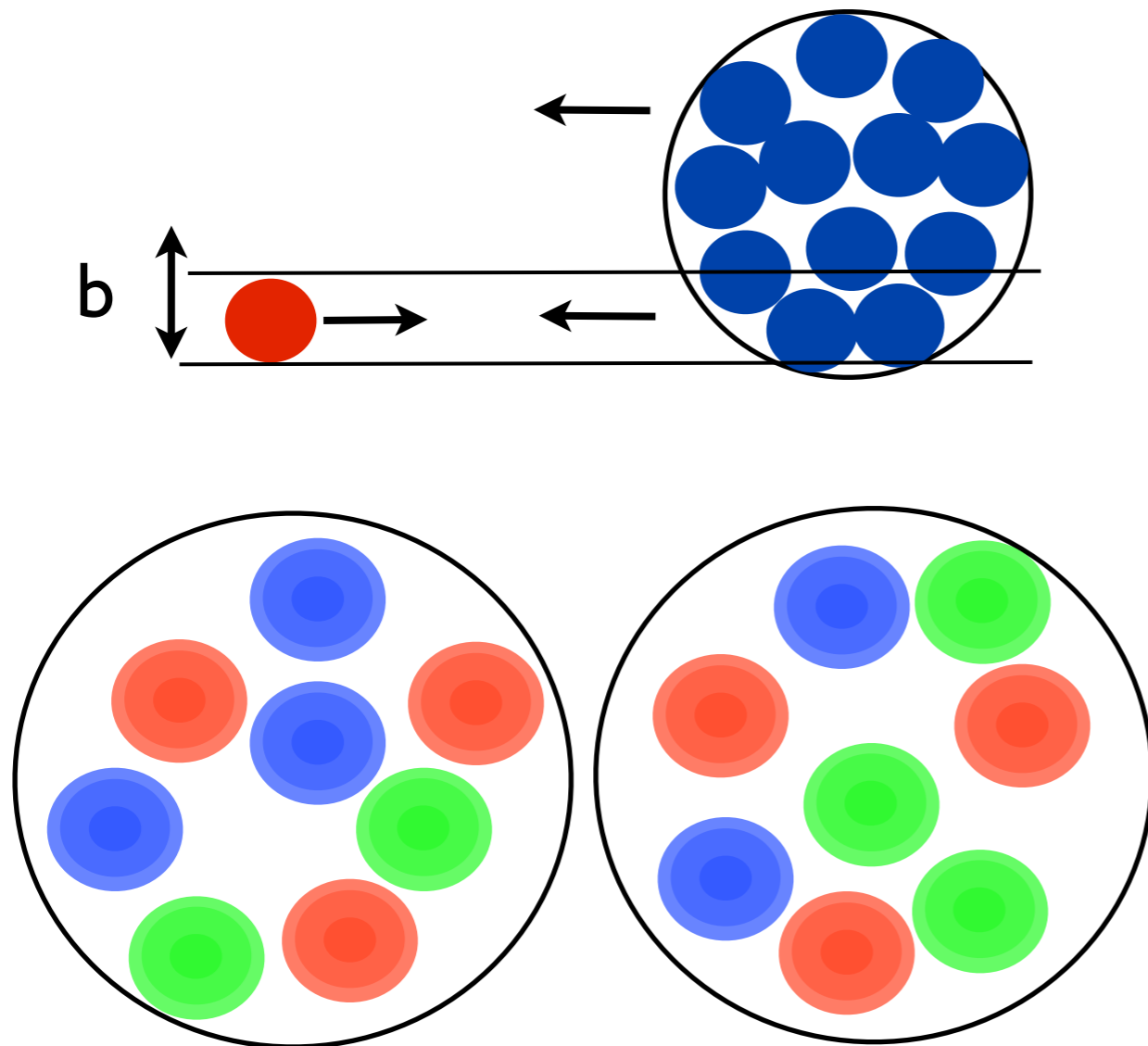
$$Q_s^2(B) = Q_0^2 B e^{\lambda(Y_b - y)}$$



Choosing $y_{lab} = (Y_b - Y_a)/2$

$$y_{max} = \frac{Y_{c.o.m}}{\lambda} \quad \lambda \approx 0.5$$

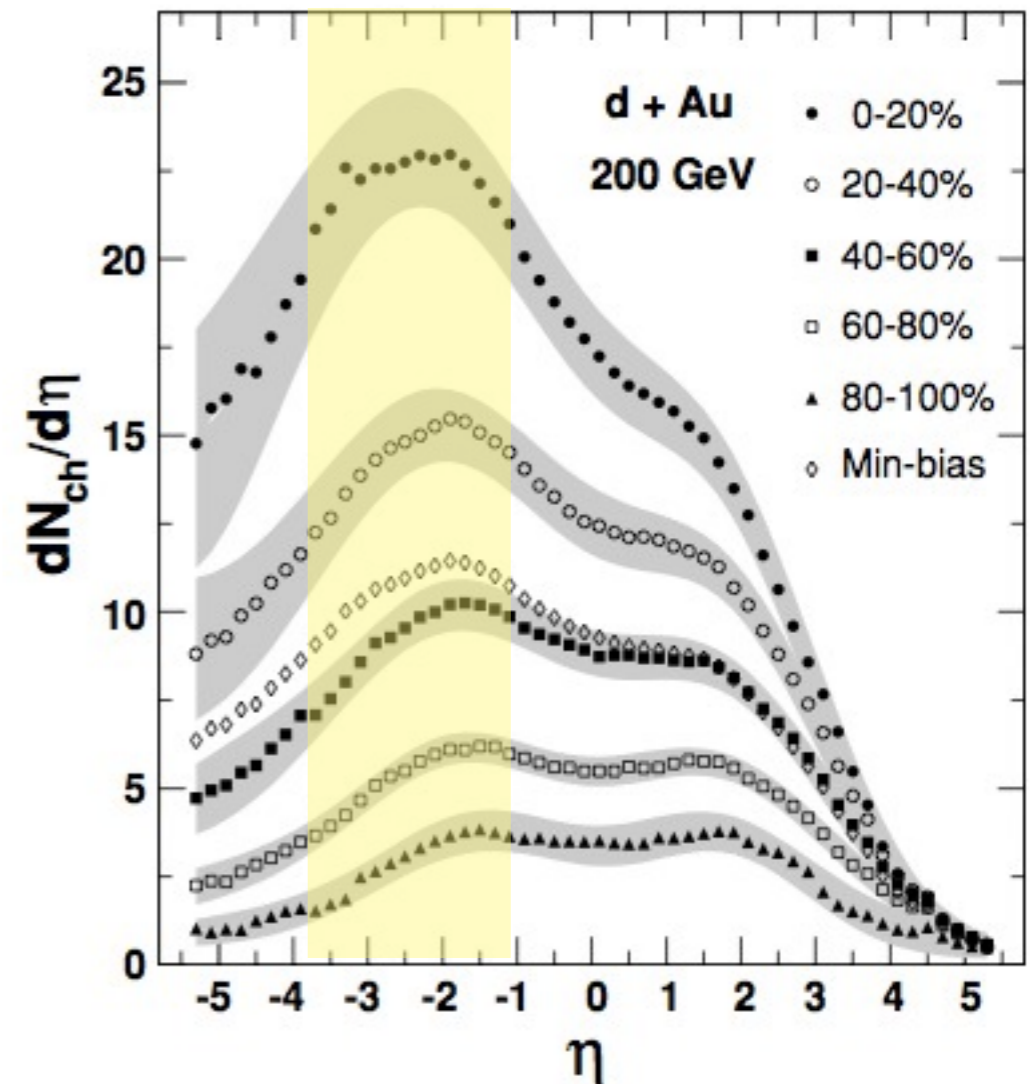
The Weak Coupling Picture



$$Q_s^2(A) = Q_s^2(B)$$

$$Q_s^2(A) = Q_0^2 A e^{\lambda(Y_a - y)}$$

$$Q_s^2(B) = Q_0^2 B e^{\lambda(Y_b - y)}$$



Choosing $y_{\text{lab}} = (Y_b - Y_a)/2$

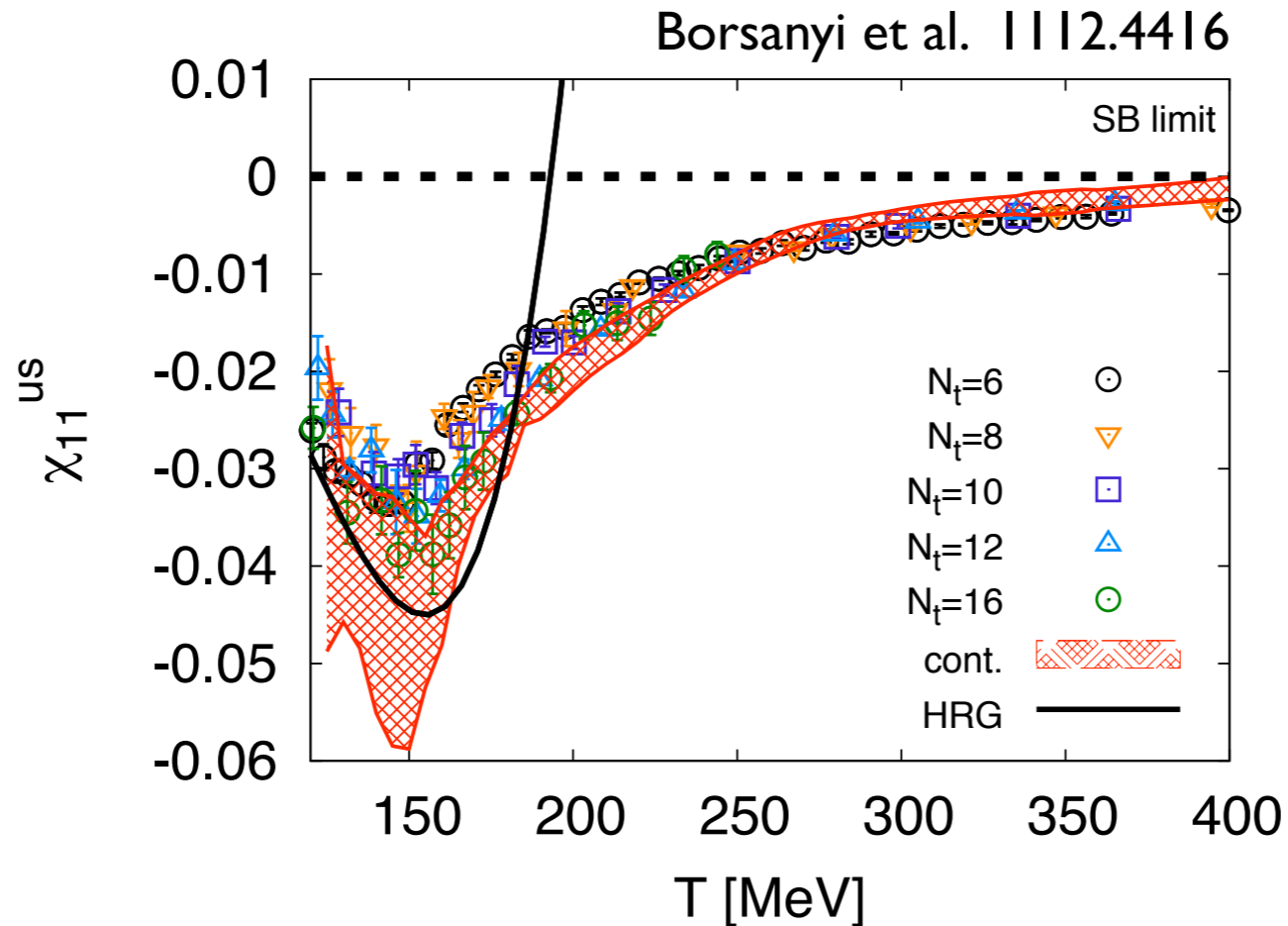
$$y_{\text{max}} = \frac{y_{\text{c.o.m}}}{\lambda} \quad \lambda \approx 0.5$$

LHC most central pPB $y_{\text{c.o.m}} = 1.7$

Conclusions

- Shock wave collisions exhibit a dynamical cross over
 - Low energy: full stopping \Rightarrow Landau hydrodynamics
 - High energy:
 - Transparency: energy propagation in the lightcone
 - Not-boost invariant initial conditions
- Longitudinal coherence on shock components
 - Fluid c.o.m = collision c.o.m.
 - Reflexion symmetric matter around the collision point
(near mid rapidity)

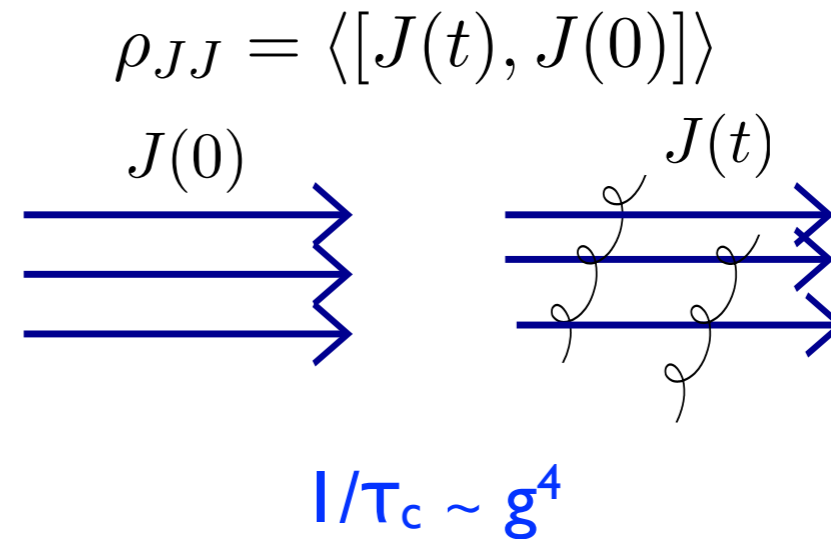
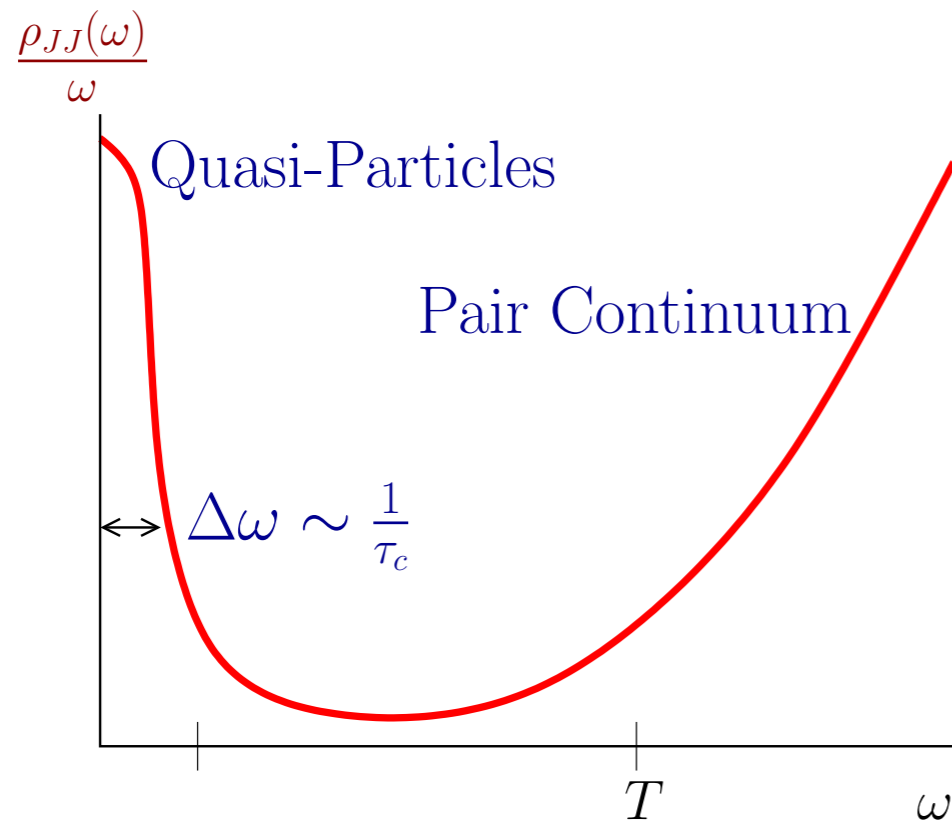
Correlations



$$\chi_{11}^{us} \propto \langle n_u n_s \rangle$$

- Most thermodynamic properties are insensitive to the nature of the degrees of freedom
- However: strong correlations among flavor
 - not compatible with an almost free gas of quarks and gluons
 - consistent with strongly correlated flavor (JCS and D. Mateos 12)

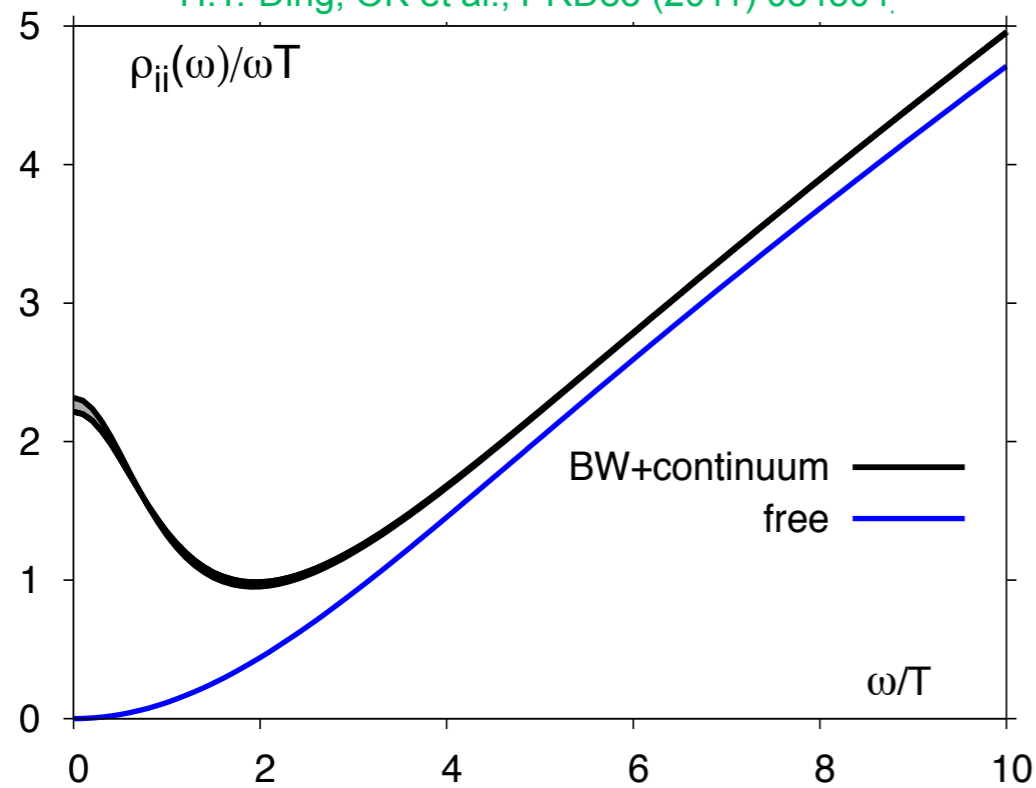
Quasi-Particles



- Fishing for quasi-particles: conserved current correlator narrow structures?

Quasi-Particles

H.T.-Ding, ÖK et al., PRD83 (2011) 034504

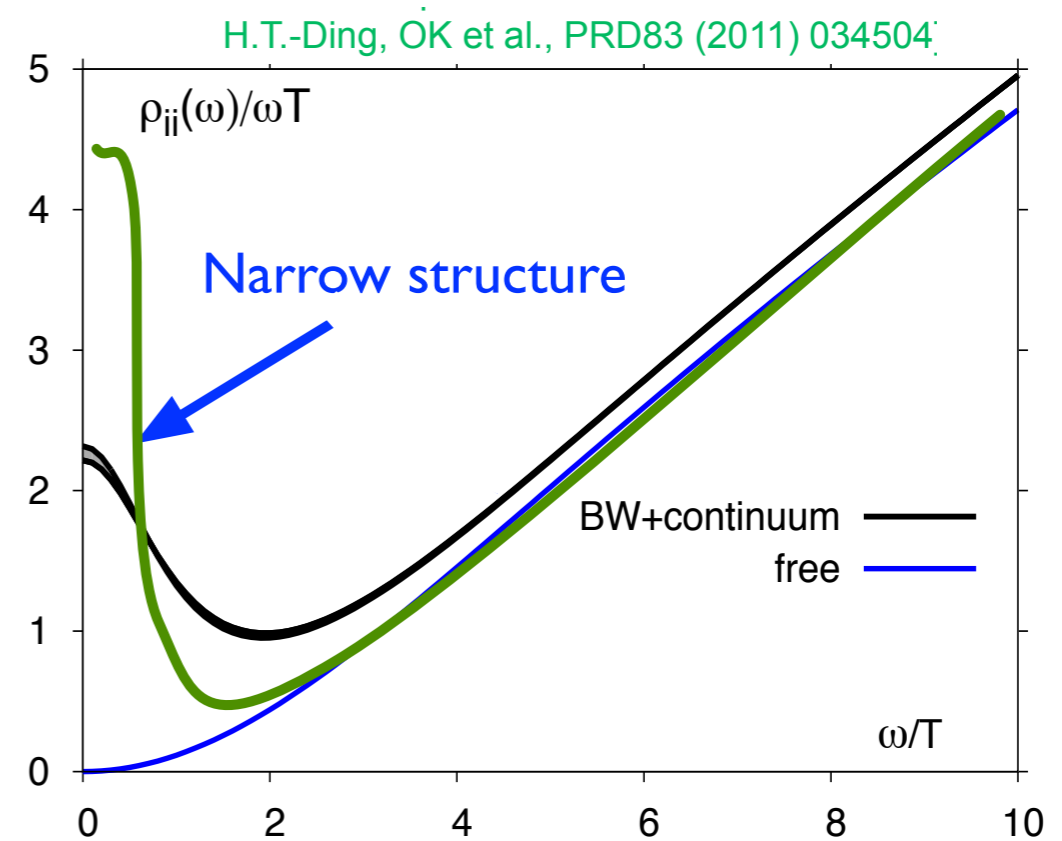
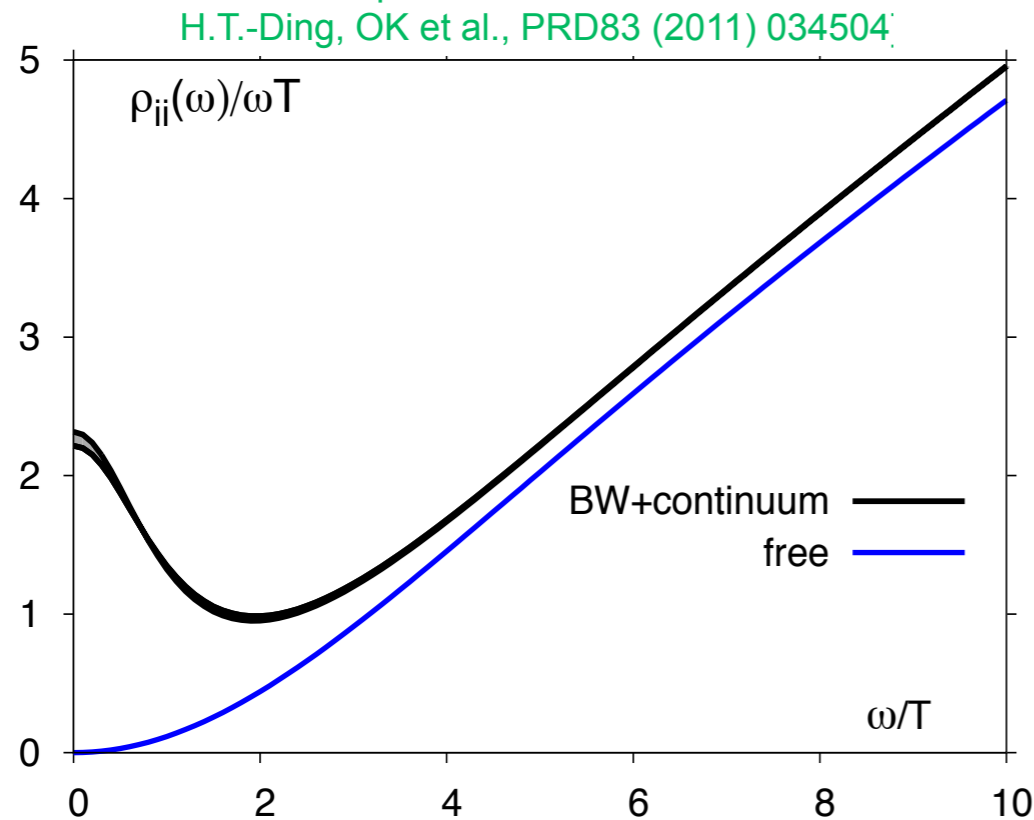


$$\rho_{JJ} = \langle [J(t), J(0)] \rangle$$

$1/\tau_c \sim g^4$

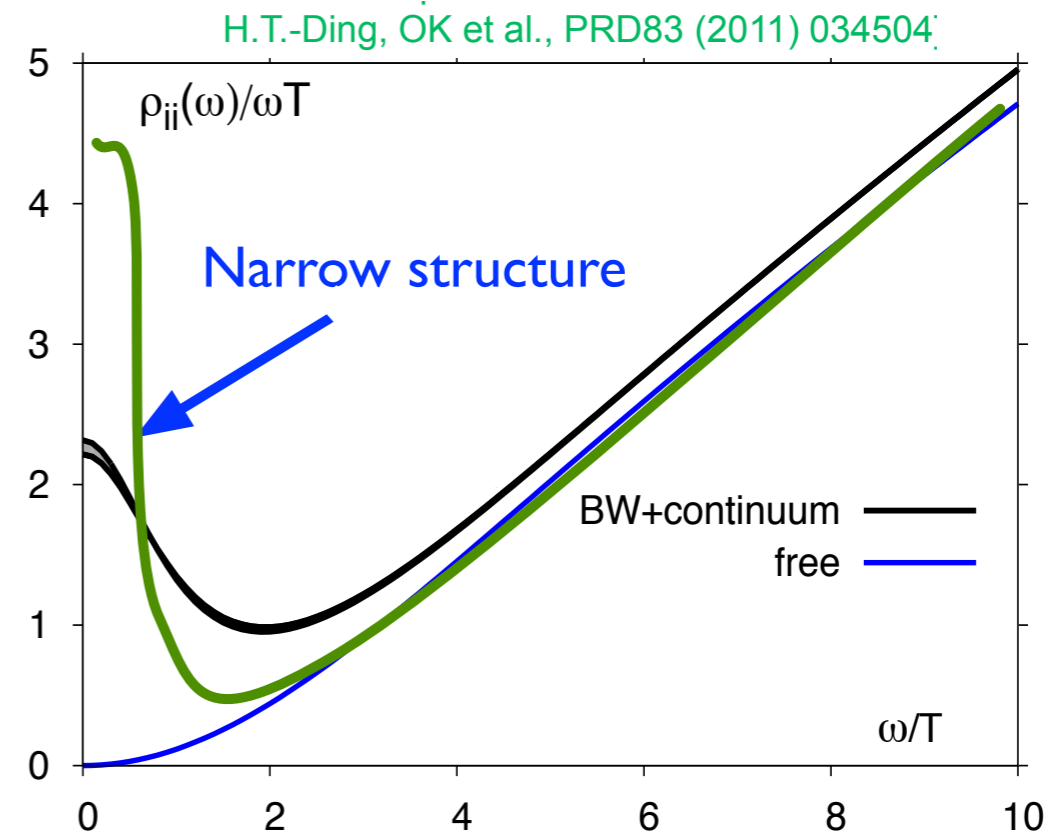
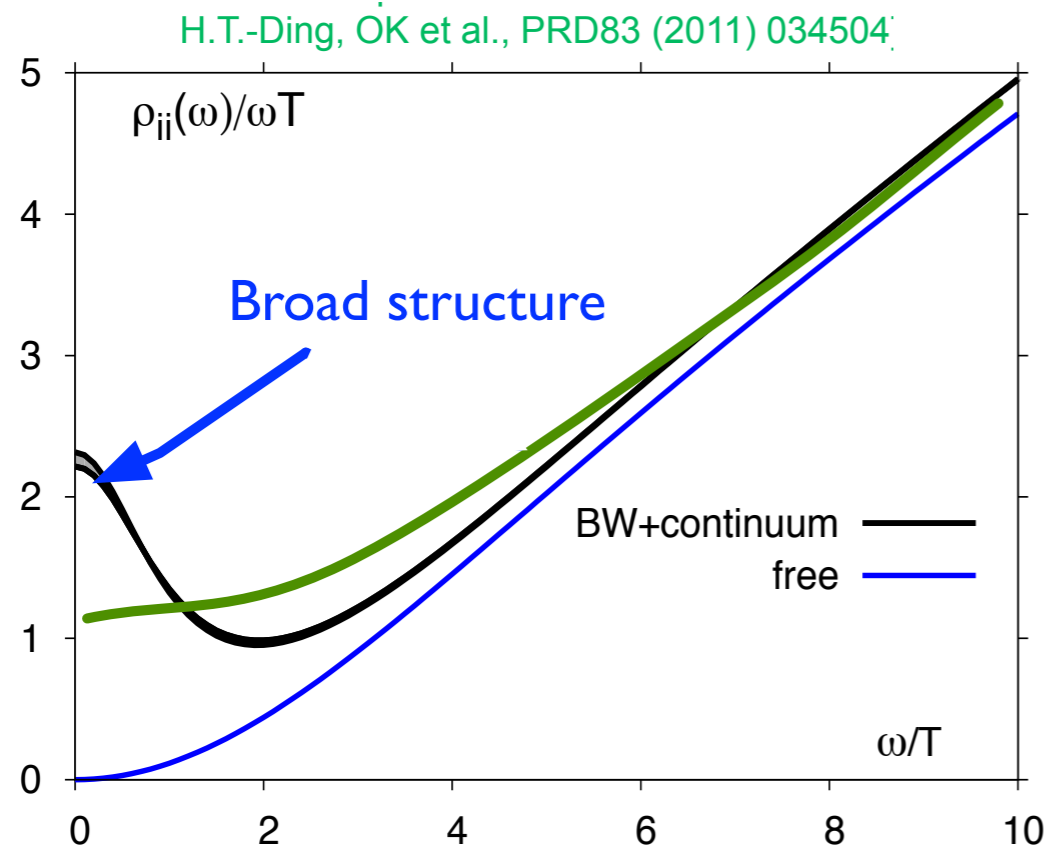
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Quasi-Particles



- Fishing for quasi-particles: conserved current correlator narrow structures?
- Lattice results (hard)
 - no clear quasi-particle peak (unlike pQCD)
 - some broad structure remainscomparable to $N_c g^2 \rightarrow \infty$ for SYM via AdS/CFT

Teaney 06

