

Noether theorem

(global transform. of fields)

$$\mathcal{L} = \mathcal{L}(\phi^I, \partial_\mu \phi^I)$$

I : collective index for indep. fields,

e.g. if $\mathcal{L} = \mathcal{L}(\psi, \psi^*, \partial_\mu \psi, \partial_\mu \psi^*)$,

then $\phi^1 = \psi$, $\phi^2 = \psi^*$.

$$\phi^I \rightarrow \phi'^I = \left(\delta^{IJ} + \varepsilon^a T_a^{IJ} \right) \phi^J -$$

- infinitesimal transf. of fields

parametrised by ε^a (ε^a are x-indep.)

Example: $\mathcal{L} = \partial_\mu \psi \partial^\mu \psi^* - V(\psi^* \psi)$

$$\left\{ \begin{array}{l} \psi \rightarrow \psi' = e^{i\alpha} \psi = (1 + i\alpha) \psi + O(\alpha^2) \\ \psi^* \rightarrow \psi^{*'} = e^{-i\alpha} \psi^* = (1 - i\alpha + O(\alpha^2)) \psi^* \end{array} \right.$$

Here $\varepsilon^a = \alpha$, $a = 1$.

$$T^{11} = i, \quad T^{22} = -i, \quad T^{12} = T^{21} = 0$$

$$T^{IJ} = i \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$$

Suppose \mathcal{L} (and thus S) is

invar. under $\delta\phi^I = \epsilon^a T_a^{IJ} \phi^J$;

$$\delta\mathcal{L} = \mathcal{L}(\phi + \delta\phi, \partial_\mu\phi + \delta\partial_\mu\phi) - \mathcal{L}(\phi, \partial_\mu\phi) = 0$$

$$\begin{aligned} \delta\mathcal{L} &= \frac{\partial\mathcal{L}}{\partial\phi^I} \delta\phi^I + \frac{\partial\mathcal{L}}{\partial\phi^I_{,\mu}} \delta\partial_\mu\phi^I = \\ &= \frac{\partial\mathcal{L}}{\partial\phi^I} \epsilon^a T_a^{IJ} \phi^J + \frac{\partial\mathcal{L}}{\partial\phi^I_{,\mu}} \epsilon^a T_a^{IJ} \partial_\mu\phi^J, \end{aligned}$$

where $\phi^I_{,\mu} \equiv \partial_\mu\phi^I$.

Since parameters ϵ^a are arbitrary,

$$\frac{\partial\mathcal{L}}{\partial\phi^I} T_a^{IJ} \phi^J + \frac{\partial\mathcal{L}}{\partial\phi^I_{,\mu}} T_a^{IJ} \partial_\mu\phi^J = 0$$

Now, on shell : $\frac{\partial\mathcal{L}}{\partial\phi^I} = \partial_\mu \frac{\partial\mathcal{L}}{\partial\phi^I_{,\mu}}$

$$\Rightarrow \partial_\mu \left(\frac{\partial\mathcal{L}}{\partial\phi^I_{,\mu}} T_a^{IJ} \phi^J \right) = 0$$

$$\text{i.e. } \partial_\mu J_a^\mu = 0,$$

(25)

$$J_a^\mu \equiv \frac{\partial \mathcal{L}}{\partial \phi_{, \mu}^I} T_a^{IJ} \phi^J$$

Noether current : conserved locally on shell due to global symmetry of the action.

Example: complex scalar field

$$J^\mu = \frac{\partial \mathcal{L}}{\partial (\partial_\mu \psi)} i \psi + \frac{\partial \mathcal{L}}{\partial (\partial_\mu \psi^*)} (-i) \psi^*$$

$$= i (\partial_\mu \psi^* \psi - \partial_\mu \psi \cdot \psi^*)$$

$$\partial_\mu J^\mu = 0 \text{ (on shell)} \Rightarrow \dot{Q} = 0,$$

$$\text{where } Q = \int d^3x J^0.$$

Noether theorem (1st) (general case)

Suppose the action S is invar. under a global continuous symmetry. In general, both fields and coordinates are affected:

$$\left\{ \begin{array}{l} x^\mu \rightarrow x'^\mu = x^\mu + \chi_a^\mu(x) \omega^a (*) \\ \phi^I(x) \rightarrow \phi'^I(x') = \phi^I(x) + \psi_a^I(x) \omega^a (**) \end{array} \right.$$

- global transf. (param. ω^a are indep. of x)
- continuous (transf. belong to a Lie group \Rightarrow enough to consider infinitesimal ones).

To simplify things, we shall suppress all indices. We also introduce new notations.

$$\delta\phi(x) = \phi'(x') - \phi(x) = \psi(x)\omega$$

$$\bar{\delta}\phi(x) \equiv \phi'(x) - \phi(x)$$

We have $\bar{\delta}(\partial\phi) = \partial(\bar{\delta}\phi)$

Indeed, $\bar{\delta}(\partial\phi) = \frac{\partial\phi'(x)}{\partial x} - \frac{\partial\phi(x)}{\partial x} =$

$$= \frac{\partial}{\partial x}(\bar{\delta}\phi).$$

Note that $\partial(\delta\phi) \neq \delta(\partial\phi)$.

$$\frac{\partial}{\partial x} \delta\phi(x) = \frac{\partial}{\partial x} (\phi'(x') - \phi(x)) =$$

$$= \frac{\partial}{\partial x'} \phi'(x') - \frac{\partial\phi(x)}{\partial x} + \frac{\partial\phi'(x')}{\partial x} - \frac{\partial\phi'(x')}{\partial x'}$$

$$= \delta(\partial\phi) + \frac{\partial\phi'(x')}{\partial x} - \frac{\partial\phi'(x')}{\partial x'}$$

Now, $\phi'(x') = \phi(x) + \psi(x)\omega$, so

$$\frac{\partial}{\partial x} \phi'(x') = \frac{\partial\phi(x)}{\partial x} + \frac{\partial\psi}{\partial x} \cdot \omega$$

Also, $\frac{\partial}{\partial x'} = \frac{\partial X}{\partial x'} \frac{\partial}{\partial X} \Rightarrow$

$$\frac{\partial X'}{\partial X} = 1 + \frac{\partial X}{\partial x} \omega \Rightarrow \frac{\partial X}{\partial X'} = 1 - \frac{\partial X}{\partial x} \omega + O(\omega^2)$$

$$\Rightarrow \frac{\partial \phi'(x')}{\partial x'} = \left(1 - \frac{\partial X}{\partial x} \omega\right) \left(\frac{\partial \phi(x)}{\partial x} + \frac{\partial \psi}{\partial x} \omega\right) + O(\omega^2)$$

$$= \frac{\partial \phi(x)}{\partial x} + \frac{\partial \psi}{\partial x} \omega - \frac{\partial X}{\partial x} \frac{\partial \phi(x)}{\partial x} \omega + O(\omega^2)$$

Thus,

$$\frac{\partial}{\partial x} \delta \phi(x) = \delta \left(\frac{\partial \phi}{\partial x}\right) + \frac{\partial \phi(x)}{\partial x} \frac{\partial X}{\partial x} \omega + O(\omega^2)$$

Consider $\delta S = \int_{V'} d\Omega' \mathcal{L}'(x') - \int_V d\Omega \mathcal{L}(x)$

$\delta S = 0$ Here $d\Omega = \sqrt{|g|} d^4x$

$$\mathcal{L}'(x') = \mathcal{L}(\phi'(x'), \frac{\partial \phi'}{\partial x'}(x'))$$

$$\delta S = \int_{V'} d\Omega' \delta \mathcal{L} + \int_{V'} d\Omega' \mathcal{L}(x) - \int_V d\Omega \mathcal{L}(x) \quad (29)$$

where $\delta \mathcal{L} = \mathcal{L}'(x') - \mathcal{L}(x)$.

$$d\Omega' = \left| \frac{\partial x'}{\partial x} \right| d\Omega = \begin{vmatrix} 1 + \frac{\partial \delta x^0}{\partial x^0} & \frac{\partial \delta x^0}{\partial x^1} & \dots \\ \frac{\partial \delta x^1}{\partial x^0} & 1 + \frac{\partial \delta x^1}{\partial x^1} & \dots \\ \vdots & \vdots & \ddots \end{vmatrix}$$

$$= \left(1 + \frac{\partial \delta x^M}{\partial x^M} + O(\omega^2) \right) d\Omega$$

(One can use e.g. $\log \det = \text{tr} \log$)

So,

$$\delta S = \int_V d\Omega \left(\delta \mathcal{L}(x) + \frac{\partial \delta x}{\partial x} \mathcal{L}(x) \right) + O(\omega^2)$$

$$= \int_V d\Omega \left[\delta \mathcal{L}(x) + \frac{\partial \mathcal{L}}{\partial x} \delta x + \frac{\partial \delta x}{\partial x} \mathcal{L} \right] + O(\omega^2)$$

Indeed, for \forall scalar $\psi(x)$: