

## Examples of $\mathcal{H}_I(x)$

- $\mathcal{H}_I(x) = J(x) \phi(x)$  (external source)

$$\begin{array}{c} \xrightarrow{x} \\ \phi(x) \quad J(x) \end{array}$$

- $\mathcal{H}_I(x) = \frac{\lambda}{4!} \phi^4(x)$  ( $\sim \phi J$ ,  $J \sim \phi^3$ )

$$\begin{array}{c} \phi(x) \\ | \\ \bullet \\ | \\ \phi(x) \\ \hline \phi(x) \quad \phi(x) \\ | \\ \phi(x) \end{array}$$

- $\mathcal{H}_I(x) = g \psi^*(x) \psi(x) \phi(x)$

$$(\sim J \phi, J \sim \psi^* \psi)$$

$$\begin{array}{c} \psi^*(x) \\ \nearrow \\ \bullet \\ \searrow \\ \psi(x) \\ \phi(x) \end{array}$$

Example:  $\mathcal{H}_I^{ip} = \frac{\lambda}{4!} \phi_{ip}^4$

Consider  $S_{fi} = \langle f | \hat{S} | i \rangle$ , where

$$|i\rangle = |p\rangle = a_p^\dagger |0\rangle$$

$$|f\rangle = |q\rangle = a_q^\dagger |0\rangle$$

$$S_{fi} \equiv \mathcal{A} = \langle q | T \exp\left[\frac{i}{\hbar} \int d^4x \mathcal{H}_I^{ip}(x)\right] | p \rangle$$

$$S = T \left( \mathbb{1} - \frac{i\lambda}{4!} \int d^4x \phi_{ip}^4(x) - \right.$$

$$\left. - \frac{1}{2!} \left(\frac{\lambda}{4!}\right)^2 \int d^4x_1 d^4x_2 \phi_{ip}^4(x_1) \phi_{ip}^4(x_2) + \dots \right)$$

$$\mathcal{A} = \langle 0 | a_q a_p^\dagger | 0 \rangle - \frac{i\lambda}{4!} \int d^4x \langle 0 | T a_q \phi_{ip}^4(x) a_p^\dagger | 0 \rangle + O(\lambda^2)$$

- $\langle 0 | a_q a_p^\dagger | 0 \rangle = \langle q | p \rangle = (2\pi)^3 2\omega_p \delta^{(3)}(\vec{p} - \vec{q})$
- $\langle 0 | T a_q \phi_{ip}(x) \phi_{ip}(x) \phi_{ip}(x) \phi_{ip}(x) a_p^\dagger | 0 \rangle$

Wick's theorem  $\Rightarrow$  2 types of contractions

$$1) \quad a_q \overbrace{\phi \phi \phi \phi}^{\text{---}} a_p^+$$

$$\overbrace{a_q \phi \phi \phi \phi}^{\text{---}} a_p^+$$

$$\overbrace{a_q \phi \phi \phi \phi}^{\text{---}} a_p^+$$

They contribute  $3 \langle 0 | a_q a_p^+ | 0 \rangle \langle 0 | T \phi \phi | 0 \rangle$   
 $\times \langle 0 | T \phi \phi | 0 \rangle$

$$2) \quad a_q \overbrace{\phi}^{\text{---}} \overbrace{\phi}^{\text{---}} \overbrace{\phi}^{\text{---}} \overbrace{\phi}^{\text{---}} a_p^+ \text{ and similar ones}$$

$$\Rightarrow 12 \langle 0 | a_q \phi_{in}(x) | 0 \rangle \langle 0 | T \phi_{in}(x) \phi_{in}(x) | 0 \rangle$$

$$\times \langle 0 | \phi_{in}(x) a_p^+ | 0 \rangle$$

$$\bullet \langle 0 | \phi_{in}(x) a_p^+ | 0 \rangle =$$

$$= \langle 0 | \int d^3 \tilde{q} (a_q e^{-iqx} + a_q^+ e^{iqx}) a_p^+ | 0 \rangle =$$

$$= \int d^3 \tilde{q} \langle 0 | a_q a_p^+ | 0 \rangle e^{-iqx} = e^{-ipx}$$

Similarly,  $\langle 0 | a_q \phi_{in}(x) | 0 \rangle = e^{iqx}$  (18)

Recall also:

$$D_F(x-y) = \int \frac{d^4 k}{(2\pi)^4} \frac{i}{k^2 - m^2 + i\epsilon} e^{-i k(x-y)}$$

$$\delta^{(4)}(x-y) = \int \frac{d^4 k}{(2\pi)^4} e^{-i k(x-y)}$$

We now have all ingredients to compute part 2):

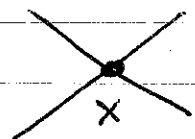
$$\begin{aligned} A_1^{(2)} &= -\frac{i\lambda}{4!} 12 \int d^4 x e^{-i(p-q)x} \int \frac{d^4 k}{(2\pi)^4} \frac{i}{k^2 - m^2 + i\epsilon} \\ &= \frac{\lambda}{2} \delta^{(4)}(p-q) \int \frac{d^4 k}{k^2 - m^2 + i\epsilon} \end{aligned}$$

What about part 1)?

But first, we draw Feynman diagrams for  $A_1^{(2)}$ :

- draw interaction vertices at  $x_1, \dots, x_n$


Here we have only one:

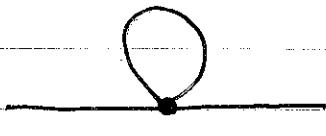


• incoming lines :  $\overline{\phi(x)} a_p^+$

outgoing lines :  $a_q \phi(x)$

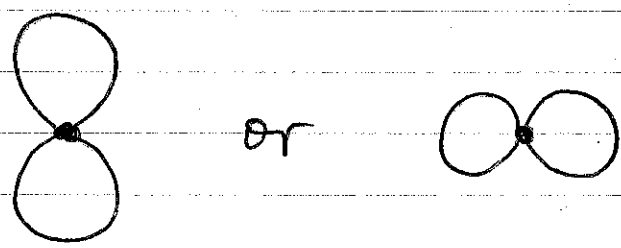
• propagators  $D_F(x-y)$  connect vertices

Here  $y=x$ , so we have 

Together : 

Now, part 1) has a piece not related to  $x$  :  $\langle 0 | a_q a_p^+ | 0 \rangle$  represented by

and also  $x$ -related piece :



In total:

$$A = \text{---} + 12 \text{---} \text{loop} + 3 \text{---} \text{two circles} + O(x^3)$$
  
connected / disconnected  
vacuum diagram  $\langle 0 | S | 0 \rangle$

$$A_2 = -\frac{\lambda^2}{2!(4!)}^2 \langle 0 | \int d^4x d^4y \overline{(a_{iP} - \phi(x) \phi(x) \phi(x) \phi(x) \phi(y) \phi(y) \phi(y) \phi(y) a_{iP}^\dagger) | 0 \rangle$$

This gives 6 diagrams:

