

# ES2. Problem 4

## Tearing Mode with Viscosity ~~XXXXXXXXXXXX~~

Equation numbers refer to my lecture notes

We now restore the viscous term in eq. (1).

~~Equation~~

$$\nu \nabla_{\perp}^4 \phi$$

$$\nu \nabla_{\perp}^4 \phi = \nu \left( \frac{\partial^2}{\partial x^2} - k^2 \right)^2 \phi_1 \quad \text{— this is added to the rhs of eq. (4)}$$

This brings in another time scale:

$$\tau_{\text{visc}} = (\nu k^2)^{-1} \quad \text{Assume } \frac{1}{\gamma} \ll \tau_{\text{visc}}$$

and another spatial scale

$$\delta_{\text{visc}} \sim \sqrt{\frac{\nu}{\gamma}} \sim \frac{1}{k} \frac{1}{\sqrt{\gamma \tau_{\text{visc}}}} \ll \frac{1}{k}$$

The outer solution is not affected by viscosity.

For the inner solution, we have ~~to~~ to add

$$\nu \left( \frac{\partial^2}{\partial x^2} - k^2 \right)^2 \phi_1 \approx \nu \frac{\partial^4 \phi_1}{\partial x^4} \quad (\text{NB: } \frac{\partial}{\partial x} \gg k)$$

to the rhs of eq. (9):

$$\gamma \phi_1'' = -k \psi_0''(0) \times \phi_1'' + \nu \phi_1^{IV} \quad (23)$$

$$\text{whence } \phi_1'' = - \frac{\gamma \phi_1''}{k \psi_0''(0)} + \frac{\nu \phi_1^{IV}}{k \psi_0''(0)} =$$

$$= - \frac{\gamma}{k \psi_0''(0)} \left[ \frac{\partial^2}{\partial x^2} \left[ 1 - \frac{\nu}{\gamma} \frac{\partial^2}{\partial x^2} \right] \phi_1 \right]$$

So we simply need to replace  $\phi_1''$  by  $\frac{\partial^2}{\partial x^2} \left( 1 - \frac{\nu}{\gamma} \frac{\partial^2}{\partial x^2} \right) \phi_1$

$\delta x$  defined by (16)

Go again to the  $\bar{X}, \chi(\bar{X})$  variables.

In place of eq. (18), we get

$$\underbrace{\frac{\eta \nu}{\delta^4} [k B_{0y}']^{-2}}_{\uparrow} \left( 1 - \underbrace{\left( \frac{\nu}{\eta \delta^2} \frac{\partial^2}{\partial \bar{X}^2} \right)}_{\downarrow} \right) \frac{\partial^2 \chi}{\partial \bar{X}^2} = \bar{X} (1 + \chi \bar{X}) \quad (24)$$

last time we defined  $\delta$  so that this is 1 [eq. (17)]  
Now we keep  $\delta$  undefined

this is  $\frac{\delta_{visc}^2}{\delta^2}$

If  $\left( \frac{\delta_{visc}^2}{\delta^2} \ll 1 \right)$ , then we simply get another boundary

layer inside the resistive layer. The solution we obtained previously is the outer solution and we have to calculate the inner solution as well.

This will produce a correction to the previous value of  $\chi$ , but the correction is small, so it is not particularly interesting.

What if, however, ~~the viscous boundary~~  $\left( \frac{\delta_{visc}^2}{\delta} \gg 1 \right)$  ?

Then Eq. (24) becomes

$$\frac{\eta \nu}{\delta^6} [k B_{0y}']^{-2} \frac{\partial^4 \chi}{\partial \bar{X}^4} = \bar{X} (1 + \chi \bar{X})$$

define  $\delta$  to set this to 1:  $\delta^6 = \eta \nu [k B_{0y}']^{-2} \quad (25)$

$$\frac{\partial^4 \chi}{\partial \bar{X}^4} = \bar{X} (1 + \chi \bar{X}) \quad (26)$$

Again we have equation with no parameters.

The matching condition (8) now becomes [using (23)]

$$\Delta' = \frac{\gamma^2}{\delta^3} [k_{Boyd}]^{-2} \int_{-\infty}^{+\infty} dX \frac{1}{X} \left( 1 - \underbrace{\frac{\gamma}{\gamma \delta^2} \frac{\partial^2 \phi}{\partial X^2}}_{\text{large}} \right) \frac{\partial^2 \phi}{\partial X^2} =$$

$$= + \frac{\gamma \nu}{\delta^5} [k_{Boyd}]^{-2} \int_{-\infty}^{+\infty} dX (1 + \phi X)$$

↑  
use (26)

again this is just a number,  $I_{visc}$

So, using (25) for  $\delta$ , we get

$$\begin{aligned} \Delta' &= \gamma \nu [k_{Boyd}]^{-2} \eta^{-5/6} \nu^{-5/6} [k_{Boyd}]^{5/3} I_{visc} \\ &= I_{visc} [k_{Boyd}]^{-1/3} \gamma \eta^{-5/6} \nu^{1/6} \end{aligned}$$

$$\boxed{\gamma = \eta^{5/6} \nu^{-1/6} \Delta' [k_{Boyd}]^{1/3} I_{visc} \quad (27)}$$

Using (25)

$$\boxed{\delta = \eta^{1/6} \nu^{1/6} [k_{Boyd}]^{-1/3} \quad (28)}$$

Thus, w/o viscosity,  $\gamma \sim \eta^{3/5}$

with viscosity,  $\gamma \sim \eta^{5/6} \nu^{-1/6}$  different scaling.

and  $\boxed{\delta_{visc} \sim \nu^{7/12} \eta^{-5/12} \Delta'^{-1/2} [k_{Boyd}]^{-1/6}}$

~~is valid for  $\delta \gg \nu^{1/2}$ . Let's check this condition:~~

~~$\delta \gg \nu^{1/2} \Rightarrow \eta^{1/6} \nu^{1/6} [k_{Boyd}]^{-1/3} \gg \nu^{1/2} \Rightarrow [k_{Boyd}]^{-1/3} \gg \nu^{1/3} \Rightarrow [k_{Boyd}]^{2/3} \gg \nu$~~

~~$\eta^{5/6} \nu^{-1/6} [k_{Boyd}]^{1/3} \gg \nu^{1/2} \Rightarrow [k_{Boyd}]^{1/3} \gg \nu^{1/3} \Rightarrow [k_{Boyd}]^{2/3} \gg \nu$~~

Now let's write these results in terms of characteristic times:

$$\gamma \sim \tau_R^{-5/6} k^{-5/3} \tau_{\text{visc}}^{1/6} k^{1/3} \Delta' \tau_A^{-1/3} \left[ \frac{B_{\text{Boy}}'}{B_{\text{Boy}}} \right]^{1/3} =$$

$$= \tau_R^{-5/6} \tau_{\text{visc}}^{1/6} \tau_A^{-1/3} \underbrace{\left( k^{-4/3} \Delta' \left[ \frac{B_{\text{Boy}}'}{B_{\text{Boy}}} \right]^{1/3} \right)}_{\text{order 1}}$$

We need ~~the magnetic Prandtl number~~

$$\gamma \tau_A \sim \tau_A^{2/3} \tau_R^{-5/6} \tau_{\text{visc}}^{1/6} \sim \left( \frac{\tau_A}{\tau_R} \right)^{2/3} \left( \frac{\tau_{\text{visc}}}{\tau_R} \right)^{1/6} \sim S^{-2/3} P_m^{-1/6}$$

where  $P_m = \frac{\nu}{\eta} = \frac{\tau_R}{\tau_{\text{visc}}}$  magnetic Prandtl number  $\downarrow$   
 $P_m \gg S^{-4/3}$

$$\gamma \tau_R \sim \tau_R^{1/6} \tau_{\text{visc}}^{1/6} \tau_A^{-1/3} \sim S^{1/3} P_m^{-1/6} \gg 1$$

$$\gamma \tau_{\text{visc}} \sim \tau_R^{-5/6} \tau_{\text{visc}}^{7/6} \tau_A^{-1/3} \sim S^{1/3} P_m^{-7/6} \gg 1$$

So we need ~~the magnetic Prandtl number~~

$$P_m \gg S^{-4/3}, P_m \gg S^{-4}, P_m \ll S^2, P_m \ll S^{2/7}$$

Assuming  $S \gg 1$ , this gives  $\frac{1}{S^{4/3}} \ll P_m \ll S^{2/7}$

Also  $\delta \sim \tau_R^{-1/6} \tau_{\text{visc}}^{-1/6} \tau_A^{1/3} L \sim \tau_R^{-1/6} \tau_{\text{visc}}^{-1/6} S^{-1/3} P_m^{1/6} L \ll L$

Finally, we required  $\frac{\nu}{\gamma \delta^2} \gg 1$ . Check this:

$$\frac{\nu}{\gamma \delta^2} \sim \tau_{\text{visc}}^{-1} \tau_R^{5/6} \tau_{\text{visc}}^{-1/6} \tau_A^{1/3} \tau_R^{1/3} \tau_{\text{visc}}^{1/3} \tau_A^{-2/3} =$$

$$= \tau_R^{7/6} \tau_{\text{visc}}^{-5/6} \tau_A^{-1/3} \sim S^{1/3} P_m^{5/6} \gg 1$$

overrides  $P_m \gg S^{-2/5}$

Thus, viscosity is important unless it is  $\ll \ll \eta$ !